A Flavor and Spectral Analysis of the Ultra-High Energy Neutrino Events at IceCube

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C.-Y. Chen, PSBD and A. Soni, Phys. Rev. D 89, 033012 (2014) [arXiv:1309.1764 [hep-ph]]; arXiv:1411.5658 [hep-ph].



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Outline

- UHE Events at IceCube
- Sources and Interactions
- SM Predictions
- Implications for New Physics

- A New Astrophysical Flux
- Conclusion

Neutrinos: Friends across 20 orders of Magnitude



[J. A. Formaggio and G. P. Zeller, Rev. Mod. Phys. 84, 1307 (2012)]

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Neutrino Flux



High-energy Neutrinos: Astrophysical Messengers





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(Ultra) High-energy Neutrino Detectors (Telescopes)

Super-Kamiokande, Baksan, Lake Baikal, ANTARES, AMANDA, IceCube , KM3Net,...



Neutrino Detection at IceCube



- Cherenkov radiation from secondary particles (muons, electrons, hadrons).
- Within the SM, neutrino interacts with matter only via weak (*W* and *Z*) gauge bosons.



CC Muon track (data)





CC electromagnetic/NC hadronic

cascade shower (data)



CC tau 'double bang' (simulation only)

First Observation of UHE Neutrinos





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Follow-Up Analysis



- 26 more events between 20-300 TeV.
- Total 28 events in 662 days of data with 4.1σ excess over expected atmospheric background (10.6^{+5.0}_{-3.6} events).

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21 cascade events and 7 muon tracks.

With 3-year Dataset

[Phys. Rev. Lett. 113, 101101(2014)]



- 9 more events, including one at 2 PeV ("Big Bird").
- Total 37 events in 988 days of data with 5.7σ excess over expected atmospheric background of 6.6^{+5.9}_{-1.6} atmospheric neutrinos and 8.4 ± 4.2 cosmic ray muons.

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28 cascade events and 9 muon tracks.

Understanding the Events

- Two main theoretical aspects:
 - Source (astrophysics): flux and flavor composition
 - Interaction (particle physics): showers and tracks
- Most plausible source: Astrophysical with a power-law flux $\Phi(E_{\nu}) = CE_{\nu}^{-s}$.



Possible Source	$N(1-2~{\rm PeV})$	$\rm N(2-10~PeV)$
Atm. Conv. [45, 46]	0.0004	0.0003
Cosmogenic–Takami [48]	0.01	0.2
Cosmogenic–Ahlers [49]	0.002	0.06
Atm. Prompt [47]	0.02	0.03
Astrophysical E^{-2}	0.2	1
Astrophysical $E^{-2.5}$	0.08	0.3
Astrophysical E^{-3}	0.03	0.06

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[R. Laha, J. F. Beacom, B. Dasgupta, S. Horiuchi and K. Murase, Phys. Rev. D 88, 043009 (2013)]

Flavor Composition

Primary production mechanisms for astrophysical neutrinos:

- $p\gamma$ process: $p\gamma \rightarrow \Delta^+ \rightarrow n\pi^+ \rightarrow ne^+\nu_e \bar{\nu}_\mu \nu_\mu$;
- pp process: $pp \to \pi^{\pm}/K^{\pm} + 2p/n \to \mu\nu_{\mu} + 2p/n \to e\nu_e \bar{\nu}_{\mu}\nu_{\mu} + 2p/n$;
- pn process: $pn \to \pi^{\pm}/K^{\pm} + 2p/n \to \mu\nu_{\mu} + 2p/n \to e\nu_e\bar{\nu}_{\mu}\nu_{\mu} + 2p/n$.

• Predict a flavor ratio of $(\nu_e : \nu_\mu : \nu_\tau) = (1:2:0)$ at source.

• Given a flavor ratio $(f_e^0:f_\mu^0:f_\tau^0)_S$, the corresponding value $(f_e:f_\mu:f_\tau)_E$ on Earth is given by

$$f_{\ell} = \sum_{\ell'=e,\mu,\tau} \sum_{i=1}^{3} |U_{\ell i}|^2 |U_{\ell' i}|^2 f_{\ell'}^0 \equiv \sum_{\ell'} P_{\ell \ell'} f_{\ell'}^0 \ .$$

For the current values of the 3-neutrino oscillation parameters, we get (1:1:1)_E at Earth.

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Possible (New Physics) Interactions

Several exotic phenomena have been invoked to explain the IceCube events, e.g.

- Decaying (PeV-scale) Dark Matter. [B. Feldstein, A. Kusenko, S. Matsumoto and T. T. Yanagida, Phys. Rev. D 88, 015004 (2013); A. Esmaili and P. D. Serpico, JCAP 1311, 054 (2013)]
- Secret neutrino interactions involving a light mediator [K. loka and K. Murase, PTEP 2014, 061E01 (2014); K. C. Y. Ng and J. F. Beacom, Phys. Rev. D 90, 065035 (2014)]
- Resonant production of TeV-scale leptoquarks. [V. Barger and W.-Y. Keung, Phys. Lett. B 727, 190 (2013)]
- Decay of massive neutrinos to lighter ones over cosmological distance scales [P. Baerwald, M. Bustamante and W. Winter, JCAP 1210, 020 (2012); S. Pakvasa, A. Joshipura and S. Mohanty, Phys. Rev. Lett. 110, 171802 (2013)]
- Pseudo-Dirac neutrinos oscillating to sterile ones in a mirror world [A. S. Joshipura, S. Mohanty and S. Pakvasa, Phys. Rev. D 89, 033003 (2014)]
- Superluminal neutrinos and Lorentz invariance violation [F. W. Stecker and S. T. Scully, Phys. Rev. D 90, 043012 (2014); L. A. Anchordoqui, V. Barger, H. Goldberg, J. G. Learned, D. Marfatia, S. Pakvasa, T. C. Paul and T. J. Weiler, Phys. Lett. B 739, 99 (2014)]

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This Talk

- Before embarking on BSM explanations, desirable to know the SM expectation with better accuracy.
- Include known sources of theoretical uncertainty (mainly from PDFs).
- Include realistic detector effects (e.g., effective number of target nucleons, attenuation effects, energy loss).
- Find the event rate for SM interactions, assuming an isotropic astrophysical, power-law flux.

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- Compare the SM predictions with the IceCube data.
- Any statistically significant deviations from the SM prediction might call for BSM!
- In the absence of significant deviations, could use the data to constrain various BSM scenarios.

SM Neutrino-Nucleon Interactions



Differential cross sections: [R. Gandhi, C. Quigg, M. H. Reno and I. Sarcevic, Astropart. Phys. 5, 81 (1996)]

$$\frac{d^2 \sigma_{\nu N}^{CC}}{dx dy} = \frac{2G_F^2 M_N E_{\nu}}{\pi} \left(\frac{M_W^2}{Q^2 + M_W^2} \right)^2 \left[xq(x, Q^2) + x\bar{q}(x, Q^2)(1-y)^2 \right],$$

$$\frac{d^2 \sigma_{\nu N}^{NC}}{dx dy} = \frac{G_F^2 M_N E_{\nu}}{2\pi} \left(\frac{M_Z^2}{Q^2 + M_Z^2} \right)^2 \left[xq^0(x, Q^2) + x\bar{q}^0(x, Q^2)(1-y)^2 \right],$$

where $x = Q^2/(2M_N y E_\nu)$ (Bjorken variable), and $y = (E_\nu - E_\ell)/E_\nu$ (inelasticity).

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Parton Distribution Functions

 q, q
 q
 (q⁰, q
 ⁰) are respectively the quark and anti-quark density distributions in a proton, summed over valence and sea quarks of all flavors relevant for CC (NC) interactions:

$$\begin{array}{lll} q & = & \displaystyle \frac{u+d}{2} + s + b, \\ \bar{q} & = & \displaystyle \frac{\bar{u}+\bar{d}}{2} + c + t, \\ q^0 & = & \displaystyle \frac{u+d}{2}(L_u^2 + L_d^2) + \displaystyle \frac{\bar{u}+\bar{d}}{2}(R_u^2 + R_d^2) + (s+b)(L_d^2 + R_d^2) + (c+t)(L_u^2 + R_u^2), \\ \bar{q}^0 & = & \displaystyle \frac{u+d}{2}(R_u^2 + R_d^2) + \displaystyle \frac{\bar{u}+\bar{d}}{2}(L_u^2 + L_d^2) + (s+b)(L_d^2 + R_d^2) + (c+t)(L_u^2 + R_u^2), \end{array}$$

with $L_u = 1 - (4/3)x_W$, $L_d = -1 + (2/3)x_W$, $R_u = -(4/3)x_W$ and $R_d = (2/3)x_W$ (where $x_W = \sin^2 \theta_W$, and θ_W is the weak mixing angle).

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- Higher E_{ν} means probing smaller *x*-regions (DIS).
- The PDFs must include the lowest possible *x*-grids (up to $\sim 10^{-9}$ extracted so far from HERA data).
- We used NNPDF2.3 [R. D. Ball et al., Nucl. Phys. B 867, 244 (2013)].

Differential Cross Sections



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[C.-Y. Chen, PSBD and A. Soni, Phys. Rev. D 89, 033012 (2014)]

Differential Cross Sections



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Total Cross Sections



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Glashow Resonance

Resonant production of W⁻ in vee⁻ scattering: [S. Glashow, Phys. Rev. 118, 316 (1960)]

 $\bar{\nu}_e + e^- \rightarrow W^- \rightarrow \text{anything}$

$$\frac{d\sigma_{\bar{\nu}_{\theta}e \to \bar{\nu}_{\theta}e}}{dy} = \frac{G_{F}^{2}m_{\theta}E_{\nu}}{2\pi} \left[\frac{R_{\theta}^{2} + L_{\theta}^{2}(1-y)^{2}}{\left(1 + 2m_{\theta}E_{\nu}y/M_{Z}^{2}\right)^{2}} + 4(1-y)^{2} \frac{1 + \frac{L_{\theta}\left(1 - 2m_{\theta}E_{\nu}y/M_{W}^{2}\right)}{1 + 2m_{\theta}E_{\nu}y/M_{Z}^{2}}}{\left(1 - 2m_{\theta}E_{\nu}/M_{W}^{2}\right)^{2} + \Gamma_{W}^{2}/M_{W}^{2}} \right]$$

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where $L_e = 2x_W - 1$ and $R_e = 2x_W$ are the chiral couplings of Z to electron.

• Peak is at energy $E_{\nu} = m_W^2/(2m_e) = 6.3$ PeV.

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where $L_e = 2x_W - 1$ and $R_e = 2x_W$ are the chiral couplings of Z to electron.

- Peak is at energy $E_{\nu} = m_W^2/(2m_e) = 6.3$ PeV.
- Proposed as an explanation of the PeV events. [A. Bhattacharya, R. Gandhi, W. Rodejohann and A. Watanabe, JCAP 1110, 017 (2011); V. Barger, J. Learned and S. Pakvasa, arXiv:1207.4571 [astro-ph.HE]]
- Disfavored by a dedicated IceCube analysis. [IceCube Collaboration, Phys. Rev. Lett. 111, 021103 (2013)]

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 A lighter W' resonance can be similarly ruled out for a range of g_{W'}, which is otherwise inaccessible experimentally. [Chen, PSBD, Soni (work in progress)]

Event Rate

$$N = TN_{A}\Omega \int_{E_{\min}}^{E_{\max}} dE_{dep} \int_{0}^{1} dy \ \Phi(E_{\nu}) V_{eff}(E_{\nu}) S(E_{\nu}) \frac{d\sigma(E_{\nu}, y)}{dy}$$

- T = 988 days for the IceCube data collected between 2010-2013.
- $N_A = 6.022 \times 10^{23} \text{ mol}^{-1} \equiv 6.022 \times 10^{23} \text{ cm}^{-3}$ water equivalent for interactions with nucleons. For interactions with electrons, $N_A \rightarrow (10/18)N_A$.
- $V_{\rm eff}(E_{\nu}) = M_{\rm eff}(E_{\nu})/\rho_{\rm ice}$ is the effective fiducial volume and $\sim 0.4 \text{ km}^3$ at PeV.



Earth Matter Effect

• $\Omega = 4\pi$ sr for an isotropic neutrino flux.

 To take into account Earth Matter effects (for upgoing events), include an attenuation factor [R. Gandhi, C. Quigg, M. H. Reno and I. Sarcevic, Astropart. Phys. 5, 81 (1996)]

$$S(E_{\nu}) = \frac{1}{2} \int_{-1}^{1} d(\cos \theta) \exp \left[-\frac{z(\theta)}{L_{\text{int}}(E_{\nu})}\right]$$

where $L_{int} = 1/(N_A \sigma)$ and $z(\theta)$ is the effective column depth obtained from PREM. [A. Dziewonski and D. L. Anderson, Phys. Earth Planet. Int. **25**, 297 (1981)]



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Makes Earth opaque to UHE neutrinos, thus limiting the upgoing events above ~ 200 TeV.
 For upgoing *τ*-neutrinos, must include regeneration effects. [S. I. Dutta, M. H. Reno and I. Sarcevic, Phys. Rev. D 62, 123001 (2000); J. F. Beacom, P. Crotty and E. W. Kolb, Phys. Rev. D 66, 021302 (2002)]

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Astrophysical Neutrino Flux

Parametrize by a single-component unbroken power-law:

$$\Phi(E_{
u}) = \Phi_0 \left(rac{E_{
u}}{E_0}
ight)^{-\gamma}$$

where Φ_0 is the total $\nu + \bar{\nu}$ flux for all flavors at $E_0 = 100$ TeV in units of $\text{GeV}^{-1}\text{cm}^{-2}\text{sr}^{-1}\text{s}^{-1}$.

• The exact value of γ depends on the source evolution model.

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- Expected to be between 2 and 2.5 for standard astrophysical sources (such as GRBs, AGNs).
- Upper bound on diffuse neutrino flux: [E. Waxman and J. N. Bahcall, Phys. Rev. D 59, 023002 (1999)]

$$[E_{\nu}^{2}\Phi_{\nu}]_{\rm WB} \approx 2.3 \times 10^{-8} \epsilon_{\pi} \xi_{Z} \, {\rm GeV cm^{-2} s^{-1} sr^{-1}}$$

Use the standard flavor composition of (1:1:1)_E corresponding to (1:2:0)_S.

Deposited Energy

 Deposited em-equivalent energy is *always* less than the incoming neutrino energy by a factor which depends on the interaction channel:

$$E_{\mathrm{em},\mathrm{e}} = (1 - y)E_{\nu}, \qquad E_{\mathrm{em},\mathrm{had}} = F_X y E_{\nu}$$

[F_X = 1 − (E_X/E₀)^{-m}(1 − f₀), with E₀ = 0.399 GeV, m = 0.130 and f₀ = 0.467 from simulations of hadronic vertex cascade [M. P. Kowalski, Ph.D. thesis, Humboldt-Universität zu Berlin (2004)]
 Contained vertex search to veto atmospheric background].











[C.-Y. Chen, PSBD and A. Soni arXiv:1411.5658 [hep-ph]]

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Zenith Angle Distribution



[C.-Y. Chen, PSBD and A. Soni, Phys. Rev. D 89, 033012 (2014)]

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χ^2 -Analysis



Two Potential Problems

- SM predictions with (1:1:1)_E flavor composition seem to be consistent with current IceCube data.
- Salient Features:
 - An unbroken power-law flux with $\gamma \simeq$ 2.5.
 - Less upgoing events due to Earth attenuation effect.
 - Most of the UHE (PeV) events are expected to be downgoing showers.
 - A possible cut-off beyond 10 PeV to explain the absence of more UHE events.

So far, no need for any exotic explanation!

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 - Less upgoing events due to Earth attenuation effect.
 - Most of the UHE (PeV) events are expected to be downgoing showers.
 - A possible cut-off beyond 10 PeV to explain the absence of more UHE events.
- So far, no need for any exotic explanation!
- However, a closer look seems to suggest two potential problems (though not statistically significant).
 - An apparent 'energy gap' between 400 TeV 1 PeV.
 - A potential 'muon deficit problem' in the high-energy bins (60 TeV $< E_{dep}$).

Two Potential Problems



Deposited EM-Equivalent Energy in Detector (TeV)

	Atm. Bkg.	(1:1:1) _E best-fit	IceCube
Total	2.8+ < 5.3	19.9	20
Up	1.5 + < 3.7	7.7	5
Down	1.2+ < 1.6	12.2	15
Track	$\sim 2.1 + < 1.0$	6.1	4
Shower	\sim 0.7 $+$ < 4.2	13.8	16
<i>p</i> -value		0.95	

Muon Deficit Problem



- A dedicated statistical analysis disfavors the (1:1:1)_E solution at 81% CL. [O. Mena, S. Palomares-Ruiz and A. C. Vincent, Phys. Rev. Lett. 113, 091103 (2014)]
- Their best-fit solution is (1:0:0)_E.
- Cannot be attained from *any* flavor ratio at an astrophysical source within the standard neutrino oscillation framework.

A Possible BSM Solution

- Invoke exotic lepton flavor violating interactions, e.g. mediated by an MeV-scale Z'.
- Could also explain the longstanding muon (g 2) anomaly.
- However, the parameter space for this to happen is very limited. [W. Altmannshofer, S. Gori, M. Pospelov and I. Yavin, Phys. Rev. Lett. 113, 091801 (2014)]



A Possible BSM Solution

 Absorption by relic neutrinos could explain the gap between 300 TeV - 1 PeV [T. Araki, F. Kaneko, Y. Konishi, T. Ota, J. Sato and T. Shimomura, arXiv:1409.4180 [hep-ph]]



- However, requires non-trivial (asymmetric) flavor structure for $Z' \bar{\ell}_{\alpha} \ell_{\beta}$ couplings, which is hard to motivate in a realistic model.
- Moreover, if a similar coupling to quarks is allowed, then ruled out by the IceCube data. [Chen, PSBD, Soni (work in progress)]

A New Solution (within the SM Framework)

- Coexistence of another astrophysical source with (1:0:0)_S flavor composition.
- Several well-motivated sources, e.g.
 - Nuclear beta decay of relativistic neutrons.
 - UHECRs interacting with relativistic electrons.
 - *e*⁺*e*⁻ scattering in a dense astrophysical system.
- Predicts a flavor ratio of (2:1:1)_E at Earth.
- Solves the muon deficit problem without invoking BSM interactions.
- Once the (2:1:1)_E flux is recognized, it is rather natural to consider a two-component flux consisting of both (1:1:1)_E and (2:1:1)_E.

Offers a simple explanation of the apparent energy gap.

χ^2 -Analysis



Comparison of the Number of Events

$$\Phi(E_{\nu}) = \Phi_1 \left(\frac{E_{\nu}}{E_0}\right)^{-\gamma_1} e^{-E_{\nu}/E_1} + \Phi_2 \left(\frac{E_{\nu}}{E_0}\right)^{-\gamma_2}$$

	Background	(1:1:1) _E	(2:1:1) _E	Two-comp	IceCube
Total	2.8+ < 5.3	19.9	19.7	19.4	20
Up	1.5+ < 3.7	7.7	7.5	7.3	5
Down	1.2+ < 1.6	12.2	12.2	12.2	15
Track	$\sim 2.1 + < 1.0$	6.1	4.1	4.3	4
Shower	$\sim 0.7 + < 4.2$	13.8	15.6	15.1	16
<i>p</i> -value		0.95	0.95	0.75	

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[C.-Y. Chen, PSBD and A. Soni, arXiv:1411.5658 [hep-ph]]

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Conclusion and Outlook

- Understanding all aspects of the UHE neutrino events at IceCube is very important for both Astrophysics and Particle Physics ramifications.
- From astrophysics point of view,
 - Need to pin down the source(s) of UHE neutrinos and their flavor composition.
 - Golden era of Neutrino Astrophysics.
- From particle physics point of view,
 - Current data seems to be consistent with the SM interactions.
 - Any significant deviations might call for BSM interpretations.
 - With more statistics, can be used to constrain (otherwise inaccessible) BSM scenarios, such as light Z'.

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 If the 'muon deficit' and/or the energy gap become statistically significant, our two-component flux can offer a natural solution within the SM framework.

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Declination



Neutrino Portal Dark Matter



[J. F. Cherry, A. Friedland and I. M. Shoemaker, arXiv:1411.1071 [hep-ph]]

Astrophysical Neutrino Flux

- Three primary mechanisms:
 - Proton collisions with energetic photons (Photo-meson production)
 - Proton-gas collision
 - Decay of UHE neutrons
- Upper bound on diffuse neutrino flux: [E. Waxman and J. N. Bahcall, Phys. Rev. D 59, 023002 (1999)]



$$[E_{\nu}^{2}\Phi_{\nu}]_{\rm WB} \approx 2.3 \times 10^{-8} \epsilon_{\pi} \xi_{Z} \, {\rm GeV cm}^{-2} {\rm s}^{-1} {\rm sr}^{-1}$$

Upper Limit on Diffuse Flux



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Different Power Law Spectra



[R. Laha, J. F. Beacom, B. Dasgupta, S. Horiuchi and K. Murase, Phys. Rev. D 88, 043009 (2013)]

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Decaying DM

- DM annihilation saturating the unitarity limit $\sigma_{ann} \leq 4\pi/(m_{DM}^2 v^2)$ cannot explain the PeV events: $\Gamma_{events} \sim V_{eff} L_{halo} n_N \sigma_N \left(\frac{\rho_{DM}}{m_{DM}}\right)^2 \langle \sigma_{ann} v \rangle \lesssim 1$ per few hundred years
- Decaying PeV-scale DM with lifetime $\tau_{DM} \simeq 1.9 N_{\nu} \times 10^{28}$ s can explain the IceCube PeV events. [B. Feldstein, A. Kusenko, S. Matsumoto and T. T. Yanagida, arXiv:1303.7320 [hep-ph]]



Leptoquarks



Resonant production at threshold energy $E_{\nu} = M_{LQ}^{2}/(2M_{N})$. [V. Barger and W. -Y. Keung, arXiv:1305.6907] $_{\odot,\odot}$