Implications of Recent Measurements in Neutrino Sector & Future Directions

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The Standard Model: Massless Neutrinos

The Standard Model is a gauge theory & it unifies strong, weak & electromagnetic forces!

$$SU(3)_C \times SU(2)_L \times U(1)_Y \Rightarrow SU(3)_C \times U(1)_{EM}$$

$(1, 2)_{-\frac{1}{2}}$ $(3, 2)_{\frac{1}{6}}$	$(1,1)_{-1}$	$(3,1)_{\frac{2}{3}}$	$(3,1)_{-\frac{1}{3}}$
$\left(\begin{array}{c} oldsymbol{ u_e} \\ e \end{array}\right)_L \left(\begin{array}{c} u^i \\ d^i \end{array}\right)_L$	e_R	u_R^i	d_R^i
$\begin{pmatrix} \nu_{\mu} \\ \mu \end{pmatrix}_L \begin{pmatrix} c^i \\ s^i \end{pmatrix}_L$	μ_R	c_R^i	s_R^i
$\begin{pmatrix} \nu_{\tau} \\ \tau \end{pmatrix}_{L} \begin{pmatrix} t^{i} \\ b^{i} \end{pmatrix}_{L}$	$ au_R$	t_R^i	b_R^i

3-fold repetition of the same representation!

- 3 active neutrinos: v_e , v_μ , v_τ
- Neutral elementary particles of Spin ½
- Only couple to *weak force* (& gravity)
- Only *left handed* neutrinos
- There are no right-handed neutrinos
- *No* Dirac Mass term: $m(\bar{\psi}_L\psi_R + \bar{\psi}_R\psi_L)$

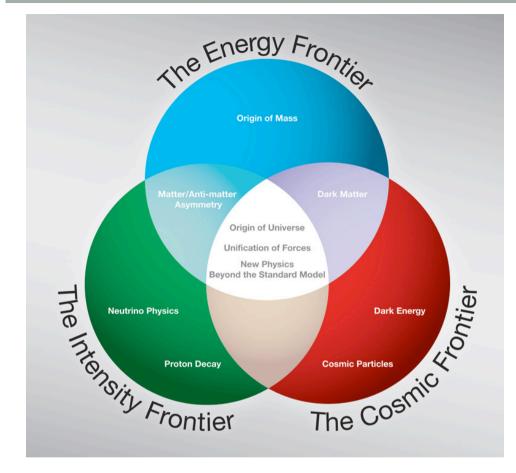
Neutrinos are massless in the Basic SM

- Over the past decade, marvelous data from world class neutrino experiments firmly established that they change flavor after propagating a finite distance
- ☐ Neutrino flavor change (oscillation) demands non-zero mass and mixing

Non-zero v mass: first experimental proof for physics beyond the Standard Model

!! An extension of the Standard Model is necessary !!

Big News in Neutrino Sector: Discovery of θ_{13}



We are going through an exciting phase

Exciting results from all the three frontiers

The Energy Frontier:
Discovery of Higgs at LHC

The Intensity Frontier: Discovery of θ_{13}

The Cosmic Frontier:
High Precision Planck measurements

BICEP2 detected B-mode polarization Smoking gun evidence for Inflation

Intensity Frontier: Neutrino properties: A window to our Universe and New Physics

Discovery of moderately large value of θ_{13} has crucial consequences for future theoretical and experimental efforts

Non-zero θ_{13} is the gateway to discover leptonic CP violation & to measure δ_{CP}

Neutrino Physics: An Exercise in Patience

Three most fundamental questions were being asked in the past century...

- 1. How tiny is the neutrino mass? (Pauli, Fermi, '30s)
- Planck + BAO + WMAP polarization data: upper limit of 0.23 eV for the sum of v masses!

 Planck Collaboration, arXiv:1303.5076 [astro-ph.CO]
- 2. Can a neutrino turn into its own antiparticle? (Majorana, '30s)
 Hunt for v-less Double- β decay (Z,A \rightarrow Z+2, A) is still on, demands lepton number violation!

 Nice Review by Avignone, Elliott, Engel, Rev.Mod.Phys. 80 (2008) 481-516
- 3. Do different v flavors 'oscillate' into one another? (Pontecorvo, Maki-Nakagawa-Sakata, '60s)

 B. Pontecorvo, Sov. Phys. JETP 26, 984 (1968) [Zh. Eksp. Teor. Fiz. 53, 1717 (1967)]

Last question positively answered only in recent years. Now an established fact that **neutrinos are massive** and leptonic flavors are not **symmetries of Nature**!

Recent measurement of θ_{13} , a clear first order picture of the 3-flavor lepton mixing matrix has emerged, signifies a major breakthrough in v physics!

This year marks the 60th anniversary since v detector of Reines & Cowan was turned on

Neutrino Oscillations in 3 Flavors

It happens because flavor (weak) eigenstates do not coincide with mass eigenstates

 $c_{ij} = \cos \theta_{ij}$ and $s_{ij} = \sin \theta_{ij}$

$$\begin{pmatrix} \nu_e \\ \nu_{\mu} \\ \nu_{\tau} \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix}$$

$$\theta_{23}$$
: $P(\nu_{\mu} \rightarrow \nu_{\mu})$ by θ_{13} : $P(\nu_{e} \rightarrow \nu_{e})$ by Reactor ν θ_{12} : $P(\nu_{e} \rightarrow \nu_{e})$ by Atoms. ν and ν beam θ_{13} & δ : $P(\nu_{\mu} \rightarrow \nu_{e})$ by ν beam Reactor and solar

$$\left(\begin{array}{ccc} c_{12} & s_{12} & 0\\ -s_{12} & c_{12} & 0\\ 0 & 0 & 1 \end{array}\right) \left(\begin{array}{c} \nu_1\\ \nu_2\\ \nu_3 \end{array}\right)$$

$$\theta_{12}$$
: P($\nu_e \rightarrow \nu_e$) by
Reactor and solar ν

$$\left(\theta_{23},\theta_{13},\theta_{12}\right)$$

Three mixing angles: $(\theta_{23}, \theta_{13}, \theta_{12})$ and one CP violating (Dirac) phase (δ_{CP})



$$\tan^2 \theta_{12} \equiv \frac{|U_{e2}|^2}{|U_{e1}|^2}; \quad \tan^2 \theta_{23} \equiv \frac{|U_{\mu 3}|^2}{|U_{\tau 3}|^2}; \quad U_{e3} \equiv \sin \theta_{13} e^{-i\delta}$$

3 mixing angles simply related to flavor components of 3 mass eigenstates

Over a distance L, changes in the relative phases of the mass states may induce flavor change!

$$P(\nu_{\alpha} \to \nu_{\beta}) = \delta_{\alpha\beta} - 4 \sum_{i>j} \text{Re}[U_{\alpha i}^* U_{\alpha j} U_{\beta i} U_{\beta j}^*] \sin^2 \Delta_{ij} - 2 \sum_{i>j} \text{Im}[U_{\alpha i}^* U_{\alpha j} U_{\beta i} U_{\beta j}^*] \sin 2\Delta_{ij}.$$

$$\Delta_{ij} = \Delta m_{ij}^2 L/4E_{\nu}$$

$$\Delta m_{ij}^2 = m_i^2 - m_i^2$$

$$\Delta_{ij} = \Delta m_{ij}^2 L/4E_{\nu}$$

$$\Delta m^2 = m^2 - m^2$$

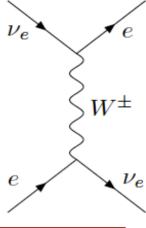
2 independent mass splittings (Δm_{21}^2) and (Δm_{32}^2) , for anti-neutrinos replace δ_{CP} by $-\delta_{CP}$

Neutrino Oscillations in Matter

Neutrino propagation through matter modify the oscillations significantly

Coherent forward elastic scattering of neutrinos with matter particles

Charged current interaction of v_e with electrons creates an extra potential for v_e



Wolfenstein matter term:

$$A = \pm 2\sqrt{2}G_F N_e E \quad 0$$

or

$$A(eV^2) = 0.76 \times 10^{-4} \rho \text{ (g/cc)} E(GeV)$$

 N_e = electron number density, + (-) for neutrinos (anti-neutrinos), ρ = matter density in Earth

Matter term changes sign when we switch from neutrino mode to anti-neutrino mode

$$P(\nu_{\alpha} \to \nu_{\beta}) - P(\bar{\nu}_{\alpha} \to \bar{\nu}_{\beta}) \neq 0$$

even if $\delta_{CP} = 0$, causes fake CP asymmetry

Matter term modifies oscillation probability differently depending on the sign of Δm^2

$$\Delta m^2 \simeq A \quad \Leftrightarrow \quad E_{\rm res}^{\rm Earth} = 6 - 8 \, {\rm GeV}$$



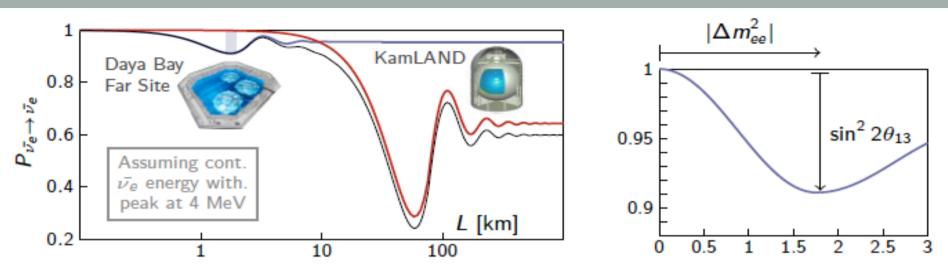
Resonant conversion – Matter effect

2002	ν	$ar{ u}$
$\Delta m^2 > 0$	MSW	· - ·
$\Delta m^2 < 0$	_	MSW



Resonance occurs for neutrinos (anti-neutrinos) if Δm^2 is positive (negative)

Short Baseline Reactor Neutrino Oscillation



 θ_{13} measured by seeing the deficit of reactor anti-neutrinos at ~ 2 km

 θ_{13} governs overall size of electron anti-neutrino deficit

Effective mass-squared difference $|\Delta m_{ee}^2|$ determines deficit dependence on L/E

$$P_{\bar{\nu_e} \to \bar{\nu_e}} = 1 - \frac{\sin^2 2\theta_{13} \sin^2 \left(\Delta m_{ee}^2 \frac{L}{4E}\right)}{\text{Short Baseline}} - \frac{\sin^2 2\theta_{12} \cos^4 2\theta_{13} \sin^2 \left(\Delta m_{21}^2 \frac{L}{4E}\right)}{\text{Long Baseline}}$$

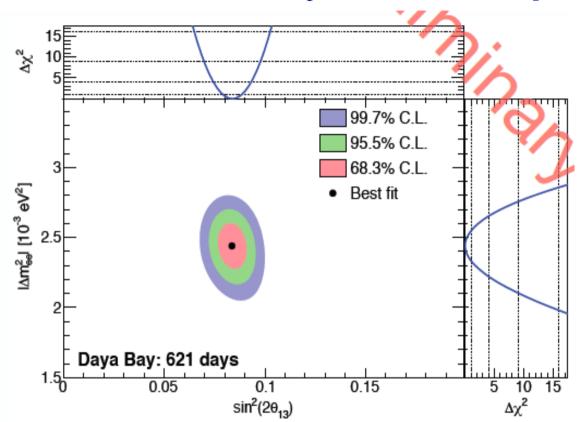
$$+ \sin^2 \theta_{12} \sin^2 \left(\Delta m_{31}^2 \frac{L}{4E}\right)$$

$$\left|\Delta m^2_{ee}\right| \simeq \left|\Delta m^2_{32}\right| \pm 5.21 \times 10^{-5} {\rm eV}^2$$
 +: Normal Hierarchy -: Inverted Hierarchy

Hierarchy discrimination requires $\sim 2\%$ precision on both Δm_{ee}^2 and $\Delta m_{\mu\mu}^2$

Latest Oscillation Results from Daya Bay

Rate + Shape Oscillation Results [Announced in Neutrino 2014]



$$\sin^2 2\theta_{13} = 0.084^{+0.005}_{-0.005}$$

 $|\Delta m^2_{ee}| = 2.44^{+0.10}_{-0.11} \times 10^{-3} \text{eV}^2$
 $\chi^2/NDF = 134.7/146$

6% precision in $\sin^2 2\theta_{13}$ achieved

Strong confirmation of oscillation-interpretation of observed $\bar{\nu_e}$ deficit

	Normal MH Δm_{32}^2 [10 ⁻³ eV ²]	Inverted MH Δm_{32}^2 [10 ⁻³ eV ²]
From Daya Bay Δm_{ee}^2	$2.39^{+0.10}_{-0.11}$	$-2.49^{+0.10}_{-0.11}$
From MINOS $\Delta m_{\mu\mu}^2$	$2.37^{+0.09}_{-0.09}$	$-2.41^{+0.11}_{-0.09}$

Present Understanding of the 2-3 Mixing Angle

Information on θ_{23} comes from: a) atmospheric neutrinos and b) accelerator neutrinos

In two-flavor scenario:
$$P_{\mu\mu} = 1 - \sin^2 2\theta_{\text{eff}} \sin^2 \left(\frac{\Delta m_{\text{eff}}^2 L}{4E}\right)$$

For accelerator neutrinos: relate effective 2-flavor parameters with 3-flavor parameters:

$$\Delta m_{\text{eff}}^2 = \Delta m_{31}^2 - \Delta m_{21}^2 (\cos^2 \theta_{12} - \cos \delta_{\text{CP}} \sin \theta_{13} \sin 2\theta_{12} \tan \theta_{23})$$

$$\sin^2 2\theta_{\text{eff}} = 4\cos^2 \theta_{13}\sin^2 \theta_{23} \left(1 - \cos^2 \theta_{13}\sin^2 \theta_{23}\right)$$
 where $\frac{|U_{\mu 3}|^2}{|U_{\tau 3}|^2} = \tan^2 \theta_{23}$

Nunokawa etal, hep-ph/0503283; A. de Gouvea etal, hep-ph/0503079

Combining beam and atmospheric data in MINOS, we have:

MINOS Collaboration: arXiv:1304.6335v2 [hep-ex]

$$\sin^2 2\theta_{\rm eff} = 0.95^{+0.035}_{-0.036} \ (10.71 \times 10^{21} \ \rm p.o.t) \\ \sin^2 2\bar{\theta}_{\rm eff} = 0.97^{+0.03}_{-0.08} \ (3.36 \times 10^{21} \ \rm p.o.t)$$

Atmospheric data, dominated by Super-Kamiokande, still prefers maximal value of $\sin^2 2\theta_{eff} = 1 \ (\geq 0.94 \ (90\% \ C.L.))$

Talk by Y. Itow in Neutrino 2012 conference, Kyoto, Japan

Bounds on θ_{23} from the global fits

In ν_{μ} survival probability, the dominant term mainly sensitive to $sin^22\theta_{23}$

If $\sin^2 2\theta_{23}$ differs from 1 (as indicated by recent data), we get two solutions for θ_{23} :

one in lower octant (LO: θ_{23} < 45 degree), other in higher octant (HO: θ_{23} > 45 degree)

In other words, if $(0.5 - \sin^2\theta_{23})$ is +ve (-ve) then θ_{23} belongs to LO (HO)

This is known as the octant ambiguity of θ_{23}

Fogli and Lisi, hep-ph/9604415

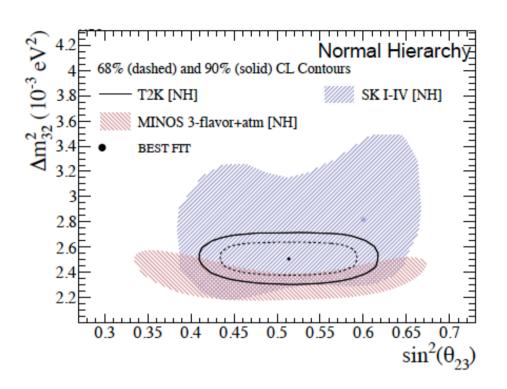
Conferences	After Neutrino 2012	After NeuTel 2013	After TAUP 2013
$\sin^2 \theta_{23}$	$0.41^{+0.037}_{-0.025} \oplus 0.59^{+0.021}_{-0.022}$	$0.437^{+0.061}_{-0.031}$	$0.446^{+0.007}_{-0.007} \oplus 0.587^{+0.032}_{-0.037}$
3σ range	$0.34 \to 0.67$	$0.357 \rightarrow 0.654$	$0.366 \to 0.663$
1σ precision (relative)	13.4%	11.3%	11.1%

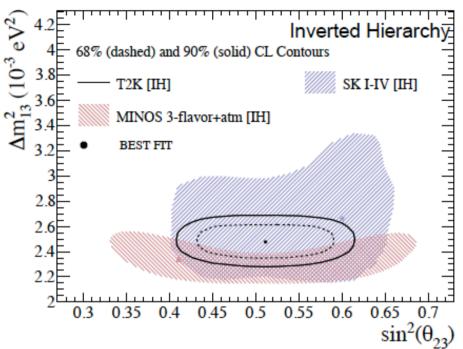
Based on Gonzalez-Garcia, Maltoni, Salvado, Schwetz, http://www.nu-fit.org

Global fit disfavors maximal 2-3 mixing at 1.4 σ confidence level (mostly driven by MINOS) v_{μ} to v_{e} oscillation data can break this degeneracy

The preferred value would depend on the choice of the neutrino mass hierarchy

New Measurements of Atmospheric Parameters





		Best-fit ± FC 68% CL (Δm² units 10 ⁻³ eV²/c⁴)
NILL	$sin^2\theta_{23}$	0.514 ^{+0.055} _{-0.056}
NH	Δm_{32}^2	2.51 ± 0.10
	$sin^2\theta_{23}$	0.511 ± 0.055
IH	Δm_{13}^2	2.48 ± 0.10

Already mixing angle is better constrained by T2K in comparison to SK and MINOS

Talk by C. Walter in Neutrino 2014

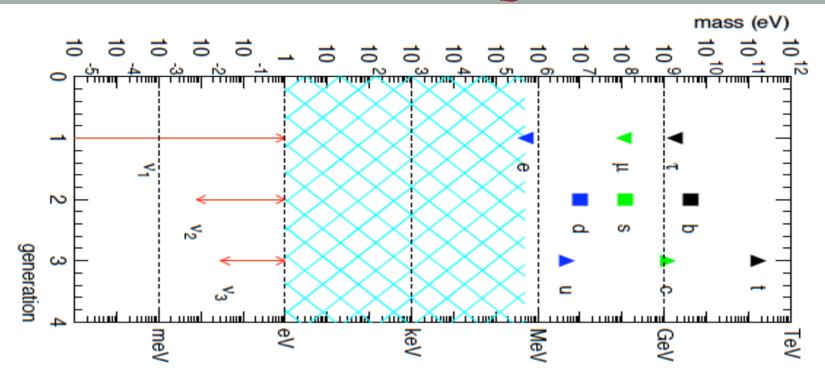
Present Status of Neutrino Oscillation Parameters

	bfp $\pm 1\sigma$	3σ range	Relative 1σ Precision
$\sin^2 \theta_{12}$	$0.306^{+0.012}_{-0.012}$	$0.271 \rightarrow 0.346$	10 1 lecision
$\theta_{12}/^{\circ}$	$33.57^{+0.77}_{-0.75}$	$31.38 \rightarrow 36.01$	4%
$\sin^2 \theta_{23}$ θ_{23} /° Non-maximal	$0.446^{+0.007}_{-0.007} \oplus 0.587^{+0.032}_{-0.037}$	$0.366 \rightarrow 0.663$	11%
$\theta_{23}/^{\circ} \frac{\text{Non}}{71.46}$	$41.9^{+0.4}_{-0.4} \oplus 50.0^{+1.9}_{-2.2}$	$37.2 \rightarrow 54.5$	11 70
$\sin^2 \theta_{13}$ Non-zero	$0.0229^{+0.0020}_{-0.0019}$	$0.0170 \rightarrow 0.0288$	8.7%
//	$8.71^{+0.37}_{-0.38}$	$7.50 \to 9.78$	0.770
$\delta_{\rm CP}/{}^{\circ}\sin\delta_{\rm CP}^{\rm P}$ C.L.	265 ⁺⁵⁶ ₋₆₁ See also the F. Capon D.V. For	zzi etal $0 \rightarrow 360$	(Not Known)
$\frac{\Delta m_{21}^2}{10^{-5} \text{ eV}^2}$	$7.45^{+0.19}_{-0.16}$	$6.98 \rightarrow 8.05$	2.4%
$\frac{\Delta m_{31}^2}{10^{-3} \text{ eV}^2} \text{ (N)}$	$+2.417^{+0.013}_{-0.013}$	$+2.247 \rightarrow +2.623$	2.5%
$\frac{\Delta m_{32}^2}{10^{-3} \text{ eV}^2}$ (I)	$-2.410^{+0.062}_{-0.062}$	$-2.602 \to -2.226$	2. 3 / 0

Based on the data available after TAUP 2013 conference

Gonzalez-Garcia, Maltoni, Salvado, Schwetz, http://www.nu-fit.org

The Two Fundamental Questions



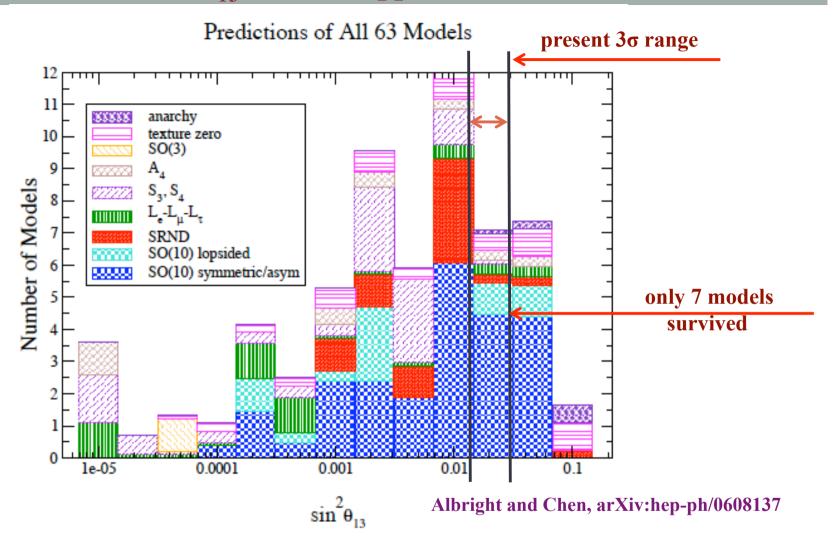
Why are neutrinos so light? The origin of Neutrino Mass!

	Neutrinos (PMNS)	Quarks (CKM)
θ_{12}	35°	13°
θ_{32}	43°	2°
θ_{13}	9°	0.2°
δ	unknown	68°

Why are lepton mixings so different from quark mixings?

The Flavor Puzzle!

Latest Results on θ_{13} : What happened to Mass models?



Survey of 63 v mass models in June 2006 by Carl H. Albright and Mu-Chun Chen

Future high precision measurements of mixing angles, new information on neutrino mass ordering and CP phase will severely constrain these presently allowed models

Implications of Recent Measurement of θ_{13}

Simplest models that are ruled out!

■ Bimaximal mixing: [Vissani (97), Barger, Pakvasa, Weiler, Whisnant (98)]

It predicts: $\theta_{12} = 45^{\circ}$, $\theta_{23} = 45^{\circ}$, and $\theta_{13} = 0^{\circ}$

predicted in flavor symmetry models with symmetry groups like A4, S4, A5

■ Tri-bimaximal mixing: [Vissani (97), Harrison, Perkins, Scot (02)]

$$U_{\text{TBM}} = R_{32} \left(\theta_{32} = \frac{\pi}{4} \right) R_{13} (\theta_{13} = 0) R \left(\theta_{21} = \tan^{-1} \left(\frac{1}{\sqrt{2}} \right) \right) = \frac{1}{\sqrt{6}} \begin{pmatrix} 2 & \sqrt{2} & 0 \\ -1 & \sqrt{2} & \sqrt{3} \\ 1 & -\sqrt{2} & \sqrt{3} \end{pmatrix}$$

• Golden ratio: [Datta, Ling, Ramond (03), Kajiyama, Raidal, Strumia (07)]

It predicts: $\theta_{12} = 31.7^{\circ}$, $\theta_{23} = 45^{\circ}$, and $\theta_{13} = 0^{\circ}$

Simplest models that are still alive!

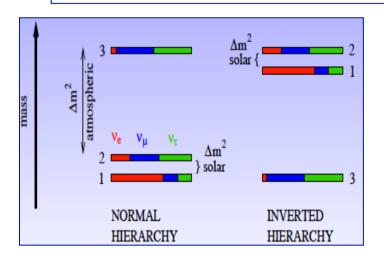
- Anarchy (v mass matrix completely random): [Hal, Murayama, Weiner (99), de Gouvea, Murayama (03, 12)] It predicts: large θ_{13} , okay with observed value of θ_{13}
- Quark-Lepton Complementarity: [Minakata, Smirnov (94), Raidal (04)]

Based on observation: θ_{12} (PMNS) + θ_{12} (CKM) = 45°

It predicts: $\sin \theta_{13} \approx \sin \theta_{C} / \sqrt{2} \approx 0.16$ (close to the observed value, other relations needs to be tested!)

Fundamental Unknowns in Neutrino Oscillation

1. What is the hierarchy of the neutrino mass spectrum, normal or inverted?



- The sign of $\Delta m_{31}^2 = m_3^2 m_1^2$ is not known!
- Currently do not know which neutrino is the heaviest?
- Only have a lower bound on the mass of the heaviest v!

$$\sqrt{2.5 \cdot 10^{-3} \text{eV}^2} \sim 0.05 \text{ eV}$$

2. What is the octant of the 2-3 mixing angle, lower ($\theta_{23} < 45^{\circ}$) or higher ($\theta_{23} > 45^{\circ}$)?

Measure θ_{23} precisely, Establish deviation from maximality at higher C.L. Then look for Octant

2. Is there CP violation in the leptonic sector, as in the quark sector?

Mixing can cause CP violation in the leptonic sector (if δ_{CP} differs from 0° and 180°) Need to measure the CP-odd asymmetries: $\Delta P_{\alpha\beta} \equiv P(\nu_{\alpha} \to \nu_{\beta}; L) - P(\bar{\nu}_{\alpha} \to \bar{\nu}_{\beta}; L)$ ($\alpha \neq \beta$)

With current knowledge of θ_{13} , resolving these unknowns fall within our reach Sub-leading 3 flavor effects are extremely crucial in current & future oscillation expts

JHEP04 (2014) 047

Analytical Understanding of Neutrino Oscillation Probability



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Analytical approximation of the neutrino oscillation matter effects at large θ_{13}

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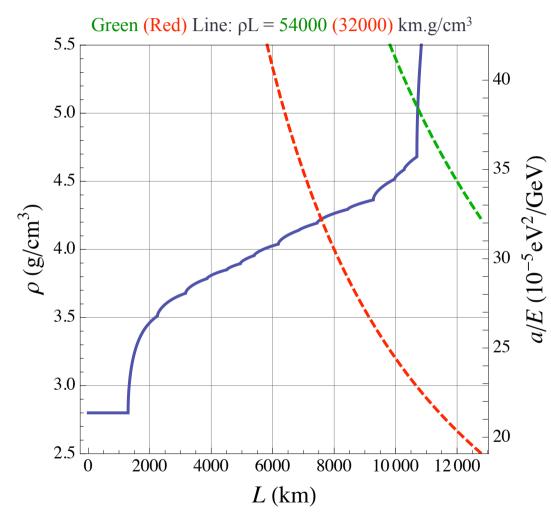
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ABSTRACT: We argue that the neutrino oscillation probabilities in matter are best understood by allowing the mixing angles and mass-squared differences in the standard parametrization to 'run' with the matter effect parameter $a = 2\sqrt{2}G_FN_eE$, where N_e is the electron density in matter and E is the neutrino energy. We present simple analytical approximations to these 'running' parameters. We show that for the moderately large

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Matter Effect Parameter a

$$a = 2\sqrt{2}G_F N_e E = 7.63 \times 10^{-5} (\text{eV}^2) \left(\frac{\rho}{\text{g/cm}^3}\right) \left(\frac{E}{\text{GeV}}\right)$$



Agarwalla, Kao, Takeuchi, JHEP 1404, 047 (2014)

- Matter effects play an important role
- Mixing angles and and mass-squared differences run with the matter effect parameter 'a'
- We present simple analytical approximations to these running parameters using the Jacobi method
- We show that for large θ_{13} , the running of θ_{23} and δ_{CP} can be neglected, simplifying the probability expression
- We need to rotate only θ_{12} and θ_{13}

Our Approach

Use the expressions for the vacuum oscillation probabilities as it is, but make the following replacements:

$$\theta_{12} \rightarrow \theta'_{12}, \quad \theta_{13} \rightarrow \theta'_{13}, \quad \delta m^2_{jk} \rightarrow \lambda_j - \lambda_k$$

where

$$\tan 2\theta_{12}' = \frac{(\delta m_{21}^2 / c_{13}^2)\sin 2\theta_{12}}{(\delta m_{21}^2 / c_{13}^2)\cos 2\theta_{12} - a}, \qquad \tan 2\theta_{13}' = \frac{(\delta m_{31}^2 - \delta m_{21}^2 s_{12}^2)\sin 2\theta_{13}}{(\delta m_{31}^2 - \delta m_{21}^2 s_{12}^2)\cos 2\theta_{13} - a},$$

$$\lambda_{1} = \lambda'_{-} \qquad \lambda'_{\pm} = \frac{(\delta m_{21}^{2} + ac_{13}^{2}) \pm \sqrt{(\delta m_{21}^{2} - ac_{13}^{2})^{2} + 4ac_{13}^{2}s_{12}^{2}\delta m_{21}^{2}}}{2}$$

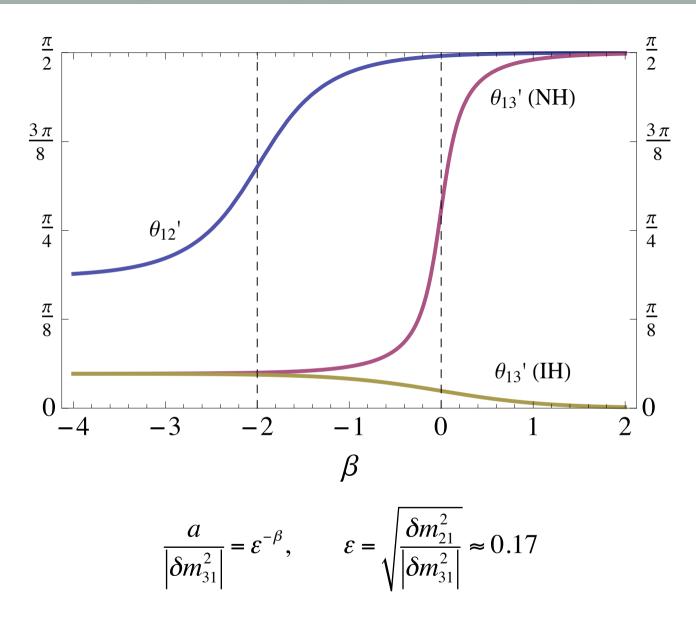
$$\lambda_{2} = \lambda''_{\pm} \qquad \lambda''_{\pm} = \frac{\left[\lambda'_{+} + (\delta m_{31}^{2} + as_{13}^{2})\right] \pm \sqrt{\left[\lambda'_{+} - (\delta m_{31}^{2} + as_{13}^{2})\right]^{2} + 4a^{2}s_{12}^{2}c_{13}^{2}s_{13}^{2}}}{2}$$

$$\lambda''_{\pm} = \frac{\left[\lambda'_{+} + (\delta m_{31}^{2} + as_{13}^{2})\right] \pm \sqrt{\left[\lambda'_{+} - (\delta m_{31}^{2} + as_{13}^{2})\right]^{2} + 4a^{2}s_{12}^{2}c_{13}^{2}s_{13}^{2}}}}{2}$$

upper (lower) sign is for the normal (inverted) hierarchy

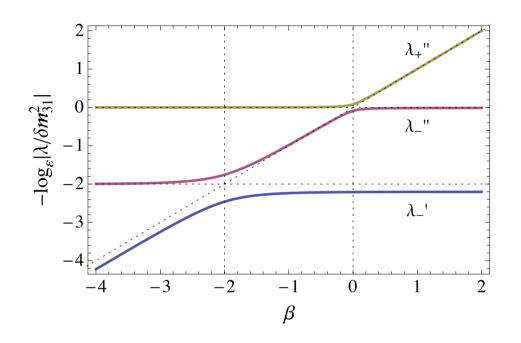
Agarwalla, Kao, Takeuchi, JHEP 1404, 047 (2014)

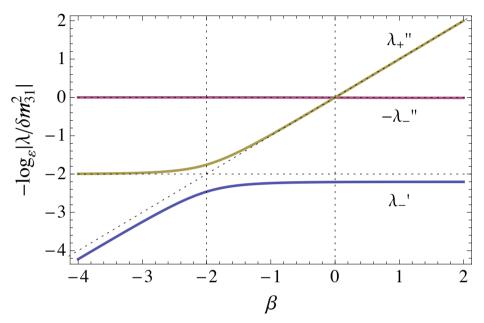
a-dependence of effective mixing angles



Agarwalla, Kao, Takeuchi, JHEP 1404, 047 (2014)

a-dependence of effective mass-squared differences



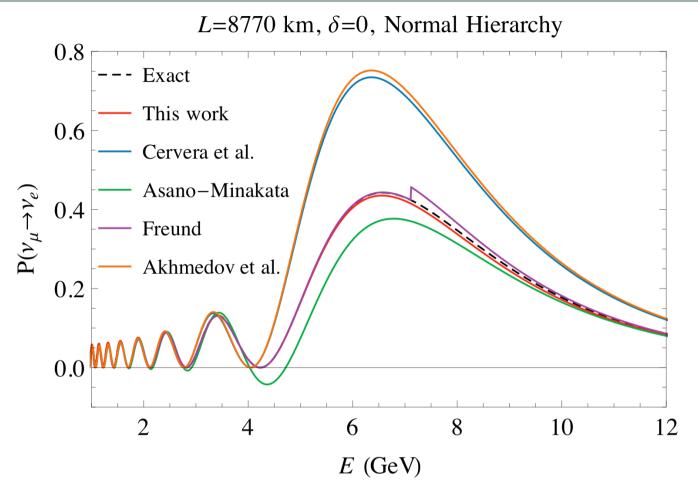


Normal Hierarchy

Inverted Hierarchy

Agarwalla, Kao, Takeuchi, JHEP 1404, 047 (2014)

Accuracy of Our Method and Comparison with Existing Literature

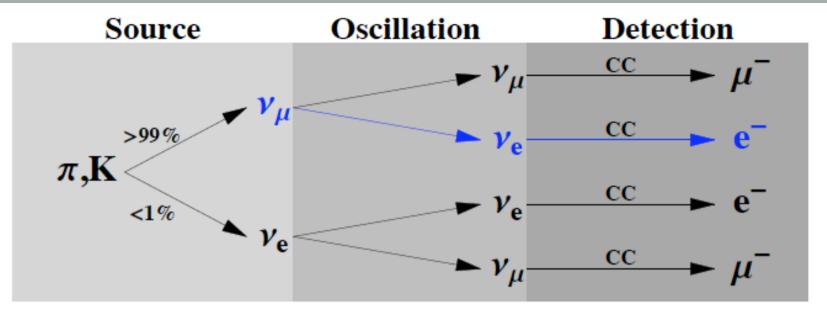


Agarwalla, Kao, Takeuchi, JHEP 1404, 047 (2014)

Other analytical expressions suffer in accuracy due to their reliance on expansion in θ_{13} , or in simplicity when higher order terms in θ_{13} included

Our method gives accurate probability for all channels, baselines and energies

Superbeams



Traditional approach: Neutrino beam from pion decay

Current Generation Experiments:

Tokai to Kamioka (T2K): 295 km (2.5° off-axis, 1st Osc. Max = 0.6 GeV)

J-PARC Beam: 0.75 MW, Total 7.8 × 10²¹ protons on target, 5 years v run

Detector: Super-Kamiokande (22.5 kton fiducial volume)

FNAL to Ash River (NOvA): 810 km (0.8° off-axis, 1st Osc. Max = 1.7 GeV) NuMI Beam: 0.7 MW, Total 3.6×10^{21} protons on target, 3 yrs v + 3 yrs anti-v Detector: 14 kton Totally Active Scintillator Detector (TASD)

Three Flavor Effects in $v_{\mu} \rightarrow v_{e}$ oscillation probability

The appearance probability $(\nu_{\mu} \rightarrow \nu_{e})$ in matter, upto second order in the small parameters $\alpha \equiv \Delta m_{21}^2/\Delta m_{31}^2$ and $\sin 2\theta_{13}$, $\frac{\sin^2 2\theta_{13} \sin^2 \theta_{23}}{(1-\hat{A})^2} \frac{\sin^2 [(1-\hat{A})\Delta]}{(1-\hat{A})^2} \rightarrow \theta_{13} \text{ Driven}$ $\alpha \frac{1}{\alpha} \frac{\sin 2\theta_{13} \xi \sin \delta_{CP} \sin(\Delta)}{\hat{A}} \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1-\hat{A})\Delta]}{(1-\hat{A})} \implies \text{CP odd}$ Resolves octant + $\alpha \sin 2\theta_{13} \xi \cos \delta_{CP} \cos(\Delta) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin[(1-\hat{A})\Delta]}{(1-\hat{A})} \implies \text{CP even}$ $+\underbrace{\alpha^2\cos^2\theta_{23}\sin^22\theta_{12}}_{0.0000};$ \Longrightarrow Solar Term where $\Delta \equiv \Delta m_{31}^2 L/(4E)$, $\xi \equiv \cos \theta_{13} \sin 2\theta_{21} \sin 2\theta_{23}$, and $\hat{A} \equiv \pm (2\sqrt{2}G_F n_e E)/\Delta m_{31}^2$ Cervera etal., hep-ph/0002108 Freund etal., hep-ph/0105071 changes sign with sgn(Δm_{31}^2) changes sign with polarity See also, Agarwalla etal., arXiv:1302.6773 [hep-ph] key to resolve hierarchy! causes fake CP asymmetry!

This channel suffers from: (Hierarchy – δ_{CP}) & (Octant – δ_{CP}) degeneracy! How can we break them?



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Potential of optimized NO ν A for large θ_{13} & combined performance with a LArTPC & T2K

Sanjib Kumar Agarwalla, a Suprabh Prakash, b Sushant K. Raut b,c and S. Uma Sankar b,d

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- $^b Department\ of\ Physics,\ Indian\ Institute\ of\ Technology\ Bombay,$
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ABSTRACT: NO ν A experiment has reoptimized its event selection criteria in light of the recently measured moderately large value of θ_{13} . We study the improvement in the sensitivity to the neutrino mass hierarchy and to leptonic CP violation due to these new features. For favourable values of $\delta_{\rm CP}$, NO ν A sensitivity to mass hierarchy and leptonic CP violation is



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Resolving the octant of θ_{23} with T2K and NO ν A

Sanjib Kumar Agarwalla, a,b Suprabh Prakash and S. Uma Sankar

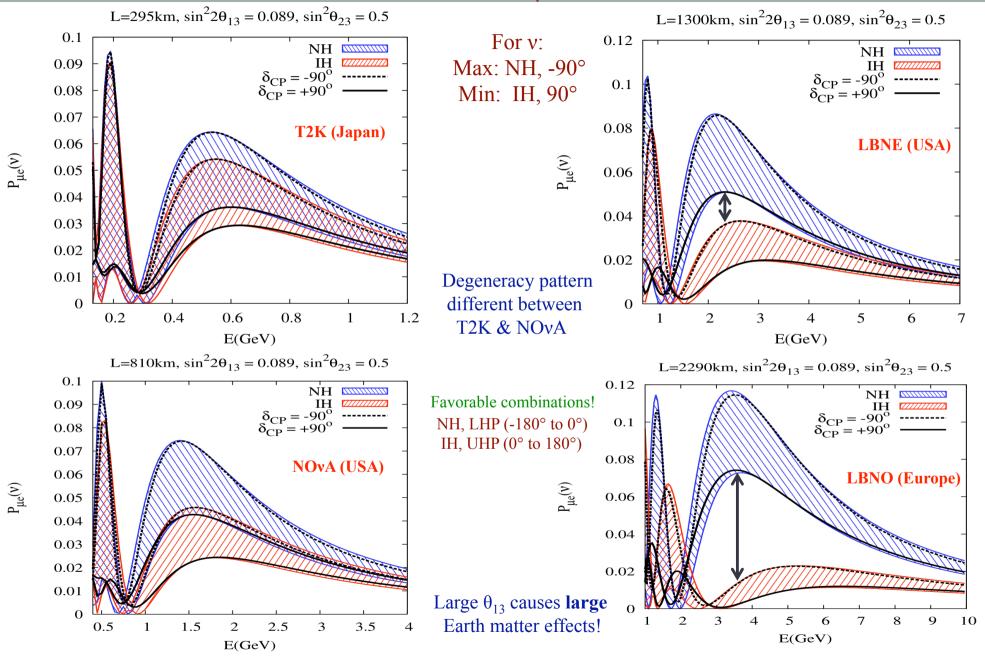
- ^a Institute of Physics, Sachivalaya Marg, Sainik School Post, Bhubaneswar 751005, India
- b Instituto de Física Corpuscular, CSIC-Universitat de València, Apartado de Correos 22085, E-46071 Valencia, Spain
- ^cDepartment of Physics, Indian Institute of Technology Bombay, Mumbai 400076, India

E-mail: sanjib@iopb.res.in, suprabh@phy.iitb.ac.in, uma@phy.iitb.ac.in

ABSTRACT: Preliminary results of MINOS experiment indicate that θ_{23} is not maximal. Global fits to world neutrino data suggest two nearly degenerate solutions for θ_{23} : one in the lower octant (LO: $\theta_{23} < 45^{\circ}$) and the other in the higher octant (HO: $\theta_{23} > 45^{\circ}$). $\nu_{\mu} \rightarrow \nu_{e}$ oscillations in superbeam experiments are sensitive to the octant and are capable of resolving this degeneracy. We study the prospects of this resolution by the current T2K and upcoming NO ν A experiments. Because of the hierarchy- $\delta_{\rm CP}$ degeneracy and the octant- $\delta_{\rm CP}$ degeneracy, the impact of hierarchy on octant resolution has to be taken into account.

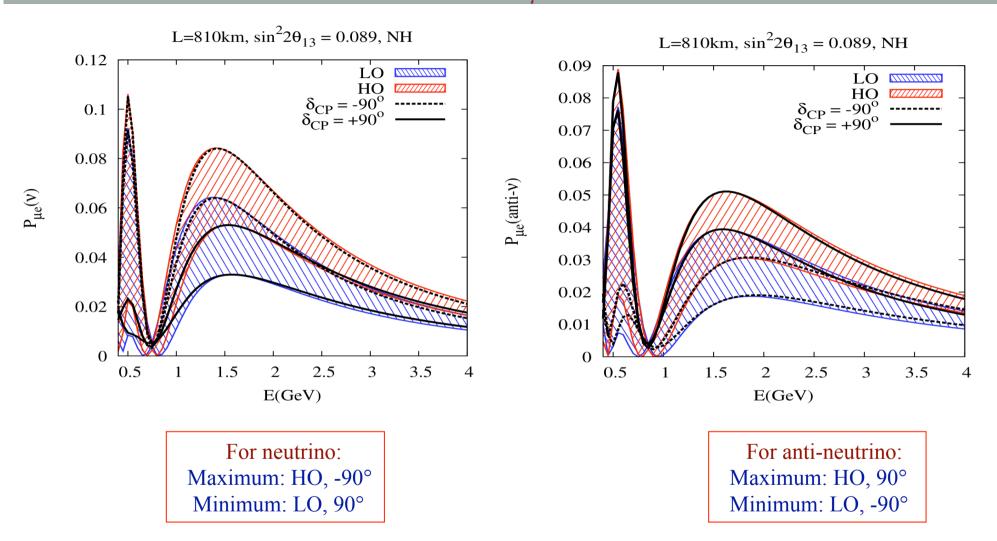
Combined data from T2K and NOvA are expected to provide the first hint for neutrino mass hierarchy, leptonic CP violation, and octant of θ_{23}

Hierarchy – δ_{CP} degeneracy in $v_{\mu} \rightarrow v_{e}$ oscillation channel



Agarwalla, Prakash, Raut, Sankar, 2012-2013

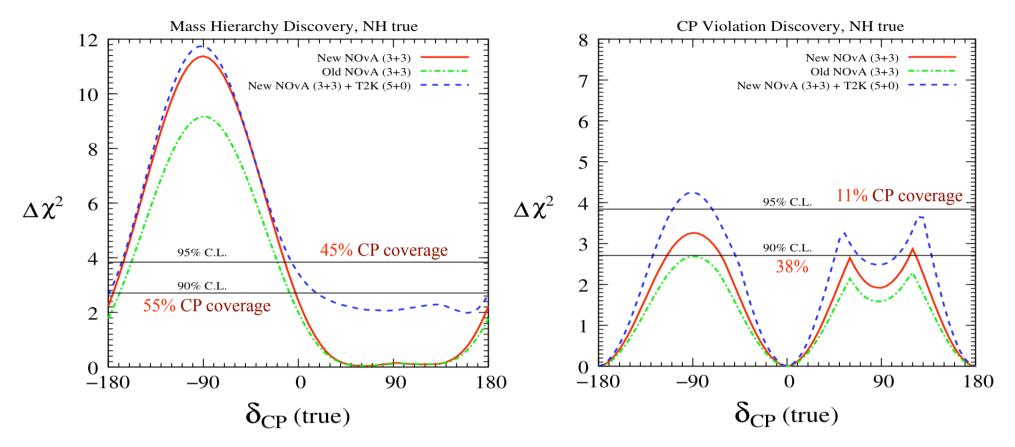
$Octant - \delta_{CP}$ degeneracy in $v_{\mu} \rightarrow v_{e}$ oscillation channel



Unfavorable CP values for neutrino are favorable for anti-neutrino & vice-versa

Agarwalla, Prakash, Sankar, arXiv: 1301.2574

Mass Hierarchy & CP Violation Discovery with T2K and NOvA

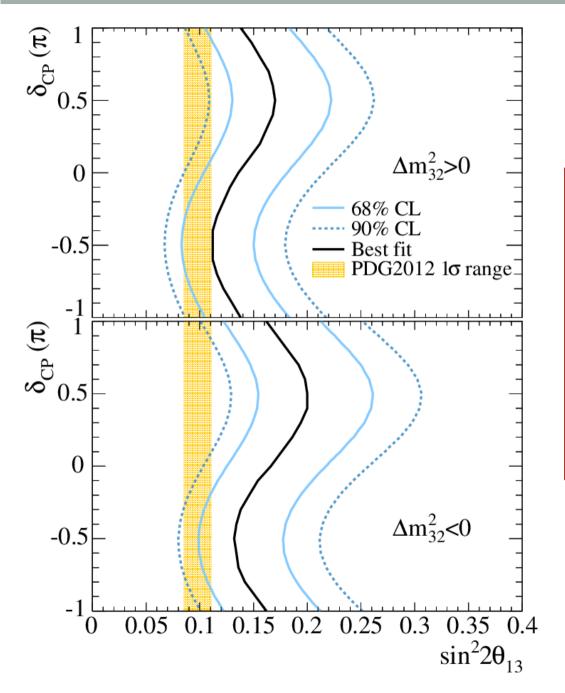


Agarwalla, Prakash, Raut, Sankar, arXiv: 1208.3644 See also, Huber, Lindner, Schwetz, Winter, arXiv: 0907.1896; Machado, Minakata, Nunokawa, Funchal, arXiv: 1307.3248; Ghosh, Ghosal, Goswami, Raut, arXiv: 1401.7243

Adding data from T2K and NOvA is useful to kill the intrinsic degeneracies

CP asymmetry ∞ 1/sin2 θ_{13} , large θ_{13} increases statistics but reduces asymmetry, Systematics are important

Important Synergy between Reactor and Accelerator data



First hint of δ_{CP} combining Reactor and Accelerator data

Best overlap is for Normal hierarchy & $\delta_{CP} = -\pi/2$

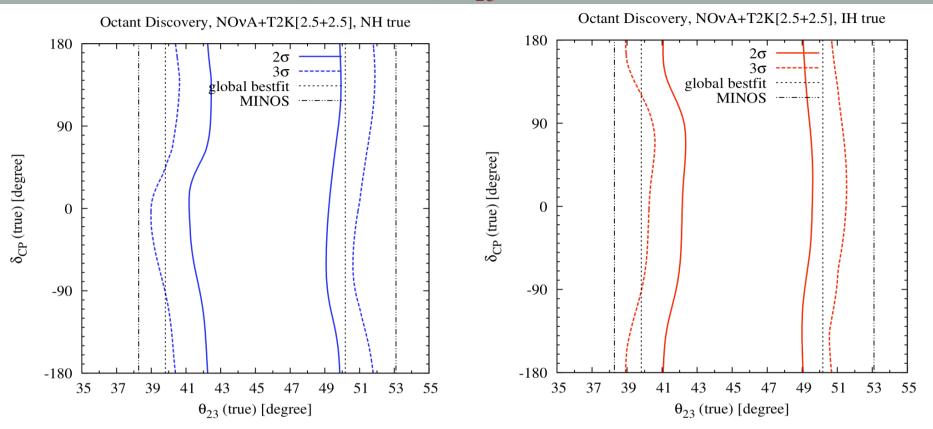
- **★** Is Nature very kind to us?
- **★** Are we very lucky?
- **★ Is CP violated maximally?**

Strong motivation for anti-neutrino run in T2K

In these plots, atmospheric parameters are marginalized over

Courtesy C. Walter (T2K Collaboration)
Talk at Neutrino 2014

Resolving Octant of θ_{23} with T2K and NOvA



Agarwalla, Prakash, Sankar, arXiv:1301.2574 [hep-ph] See also, Chatterjee, Ghoshal, Goswami, Raut, arXiv:1302.1370 [hep-ph]

If θ_{23} < 41° or θ_{23} > 50°, we can resolve the octant issue at 2 σ irrespective of δ_{CP} If θ_{23} < 39° or θ_{23} > 52°, we can resolve the octant issue at 3 σ irrespective of δ_{CP}

Important message: T2K must run in anti-neutrino mode in future

Future Facilities for Long Baseline Neutrino Experiments

LBNE: FNAL to Homestake : 1300 km (1st Osc. Max = 2.52 GeV)

Beam: 120 GeV, 0.7 MW, 6×10^{20} POT/yr, 5 yrs v + 5 yrs anti-v

Detector: 10 kton LArTPC (Phase1), 35 kton LArTPC (Phase2)

LBNO: CERN to Phyasalmi : 2300 km $(1^{st} Osc. Max = 4.54 GeV)$

Beam: 400 GeV, 0.77 MW, 1.5×10^{20} POT/yr, 5 yrs v + 5 yrs anti-v

Detector: 20 kton LArTPC (Phase1), 70 kton LArTPC (Phase2)



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Exploring the three flavor effects with future superbeams using liquid argon detectors

Sanjib Kumar Agarwalla, Suprabh Prakash and S. Uma Sankar

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ABSTRACT: Recent measurement of a moderately large value of θ_{13} signifies an important breakthrough in establishing the standard three flavor oscillation picture of neutrinos. It has provided an opportunity to explore the sub-dominant three flavor effects in present and future long-baseline experiments. In this paper, we perform a comparative study of the physics reach of two future superbeam facilities, LBNE and LBNO in their first phases of run, to resolve the issues of neutrino mass hierarchy, octant of θ_{23} , and leptonic CP violation. We also find that the sensitivity of these future facilities can be improved significantly by adding the projected data from T2K and NO ν A. Stand-alone LBNO setup

These powerful future facilities are capable enough to settle the remaining fundamental unknowns in neutrino oscillation at high C.L., needed to claim the discovery

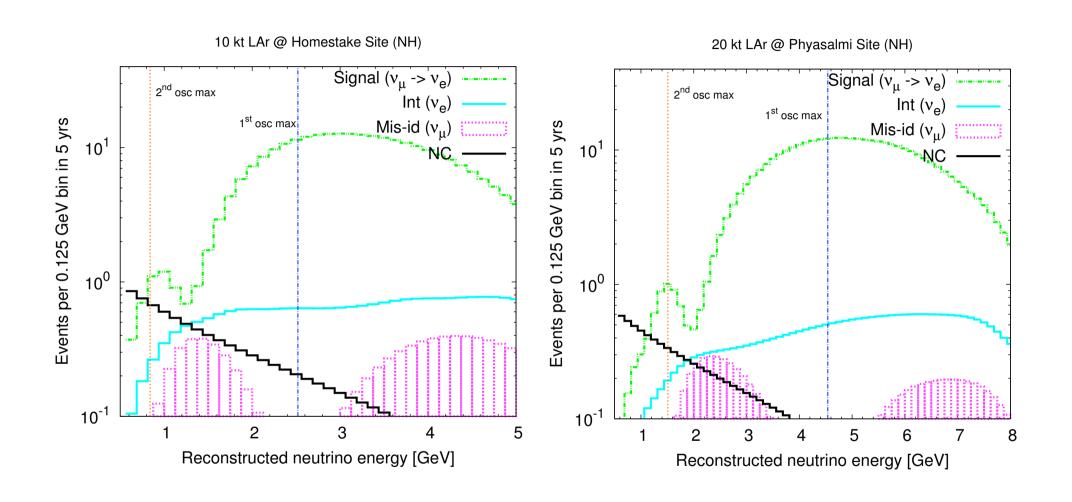
Projected data from T2K and NOvA will play a crucial role in the first phases of LBNE and LBNO

HEP03 (2014) 08

^a Institute of Physics, Sachivalaya Marg, Sainik School Post, Bhubaneswar 751005, India

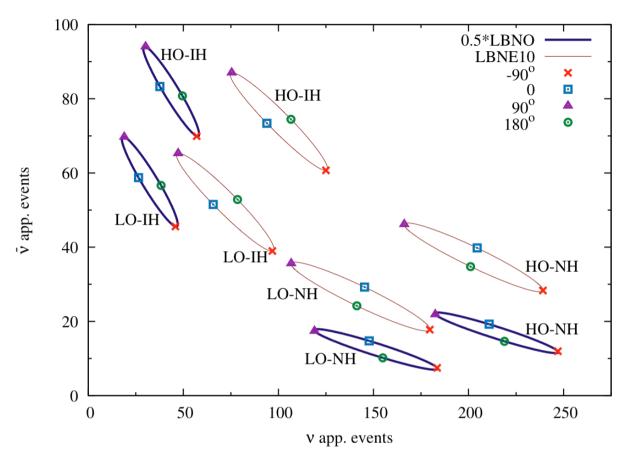
^b Department of Physics, Indian Institute of Technology Bombay, Mumbai 400076, India

Event Spectrum in LBNE and LBNO Experiments



Review on long-baseline phenomenology, S.K. Agarwalla, arXiv: 1401.4705

Future Superbeam Expts with LAr Detector: LBNE & LBNO



LBNO with 10 kt LArTPC

(LO/HO)-IH ellipses well separated from (LO/HO)-NH ellipses

Excellent hierarchy discrimination capability with just neutrino data

For octant, balanced v & anti-v data must

LBNE with 10 kt LArTPC

For LO, hierarchy discovery is limited

Octant discovery: similar to 0.5*LBNO

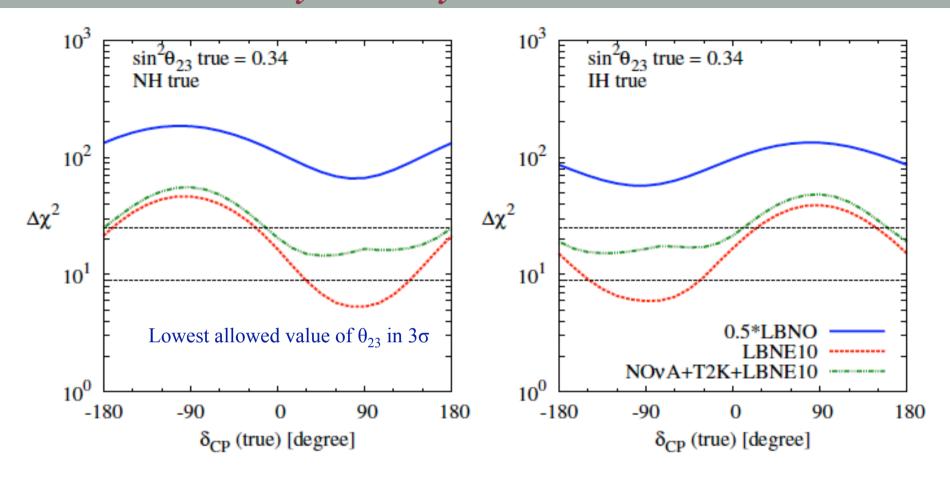
Agarwalla, Prakash, Sankar, arXiv:1304.3251 [hep-ph]

Wide Band Beam → Higher statistics → Cover several L/E values → Kill clone solutions

LAr Detector → Excellent Detection efficiency at 1st & 2nd Osc. maxima, good background rejection

High L → High E → High cross-section → Less uncertainties in cross-section at high E

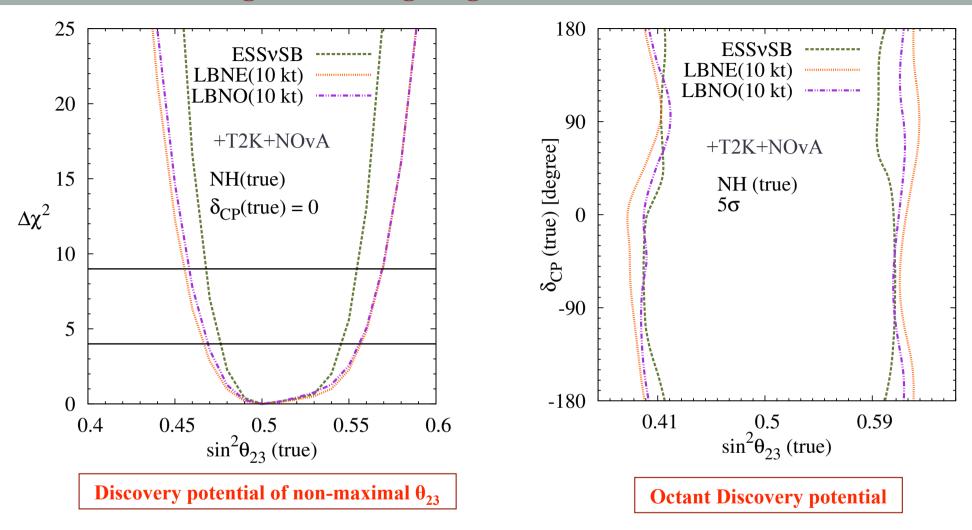
Median Hierarchy Discovery Potential with LBNE and LBNO



Agarwalla, Prakash, Sankar, arXiv:1304.3251 [hep-ph]
See also, arXiv:1312.6520 [hep-ph] from LAGUNA-LBNO Collaboration

LBNO w/ 10 kt $> 7\sigma$ median hierarchy discovery irrespective of the choice of θ_{23} - δ_{CP} -hierarchy LBNE w/ 10 kt $+ T2K + NOvA > 3\sigma$ median hierarchy discovery for any parameter choice

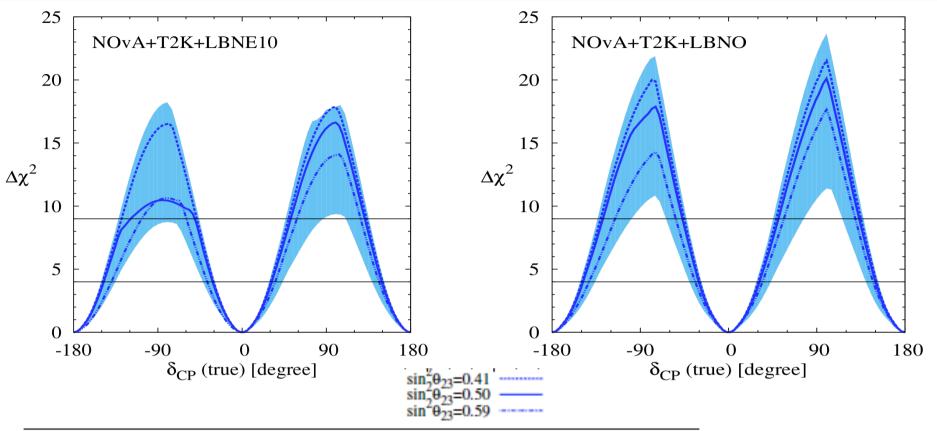
Probing 2-3 Mixing Angle with LBNE and LBNO



Agarwalla, Choubey, Prakash, arXiv:1406.2219 [hep-ph]

If $\sin^2\theta_{23} \le 0.39$ or $\sin^2\theta_{23} \ge 0.62$, octant can be resolved at 5σ irrespective of δ_{CP}

CP violation Discovery with LBNE and LBNO



Sotups	Fraction of $\delta_{\rm CP}({\rm true})$	
Setups	2σ confidence level	3σ confidence level
LBNE10 (10 kt)	0.51	0.03
$LBNE10 + T2K + NO\nu A$	0.63	0.43
LBNO (20 kt)	0.51	0.23
$LBNO + T2K + NO\nu A$	0.69	0.46

Assuming maximal mixing as true choice

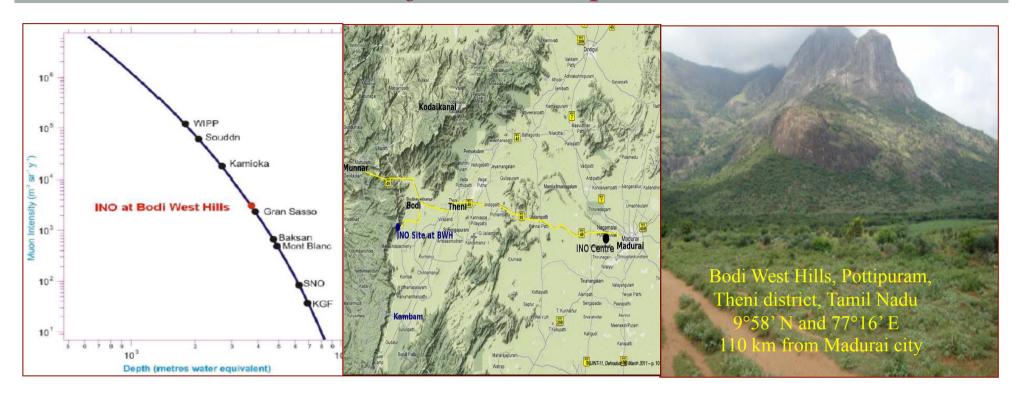
Agarwalla, Prakash, Sankar, arXiv:1304.3251 [hep-ph]

India-Based Neutrino Observatory

- A multi-institutional attempt to build a world-class underground facility to study fundamental issues in science with special emphasis on neutrinos
- With ~1 km all-round rock cover accessed through a 2 km long tunnel.

 A large and several smaller caverns to pursue many experimental programs
- Complementary to ongoing efforts worldwide to explore neutrino properties
- A mega-science project (~250 M\$) in India, jointly funded (50:50) by the Department of Atomic Energy and the Department of Science and Technology
- INO project was discussed and approved by the Atomic Energy Commission on 17th August, 2013 at New Delhi
- Regarding Final approval: Clearance from the Cabinet expected soon
- International Community is welcome to participate in ICAL@INO as well as the INO facility is available to the entire community for setting up experiments like Neutrino-less Double Beta Decay, Direct Dark Matter searches

Location of INO & Unique Features



- > Transport:
- Geotechnical
 Issues:
- > Environmental Issues:
- Weather:

Flat terrain with good access from major roads

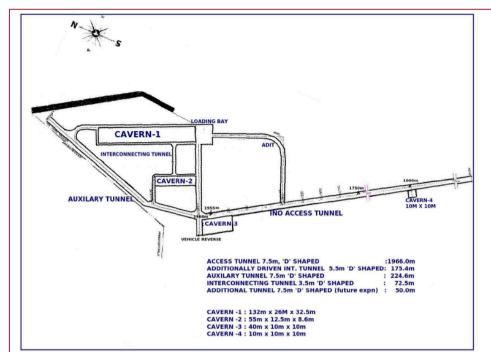
Good rock quality, Cavern set in massive Charnockite rock under the 1589 m peak, Vertical cover approx. 1289 m, Tunnel length 1.91 km

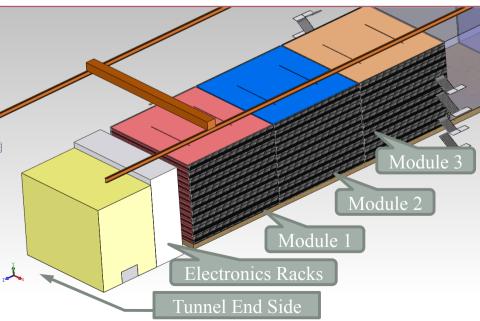
Portal set outside the Reserved Forest boundary, no disturbance. Surface facilities not on Forest Land. No clearing of forest

Warm, low rainfall area, low humidity throughout the year

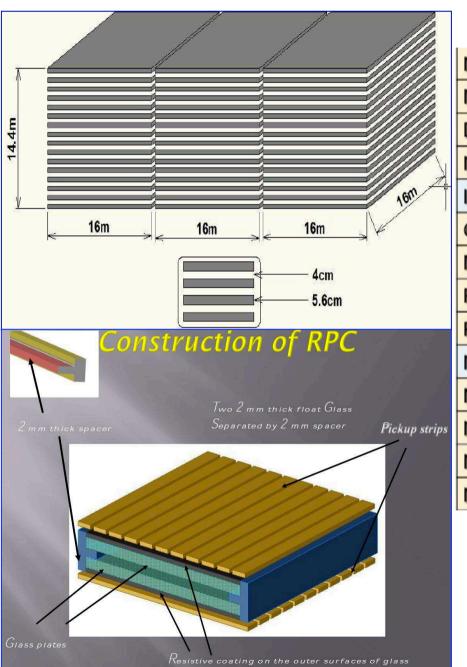
Approved projects under INO

- Come up with an underground lab & surface facilities near Pottipuram village in Theni district of Tamil Nadu
- Build massive 50 kt magnetized Iron calorimeter (ICAL) detector to study properties of atmospheric neutrinos
- Construction of INO centre at Madurai: Inter-Institutional Centre for High Energy Physics (IICHEP)
- Human Resource Development (INO Graduate Training Program)
- Completely in-house Detector R&D with substantial INO-Industry interface
- Time Frame for 1st module: 2019





Specifications of the ICAL Detector



No. of modules	3
Module dimensions	16m×16m×14.5m
Detector dimensions	48.4m × 16m × 14.5m
No. of layers	150
Iron plate thickness	56mm
Gap for RPC trays	40mm
Magnetic field	1.3Tesla
RPC dimensions	1,950mm×1,840mm×24mm
Readout strip pitch	3 omm
No. of RPCs/Road/Layer	8
No. of Roads/Layer/Module	8
No. of RPC units/Layer	192
No. of RPC units	28,800 (97,505m²)
No. of readout strips	3,686,400

Rapid progress in all fronts
2011-2014: A productive phase for INO
Several milestones achieved

Physics Issues with ICAL-INO in Phase 1

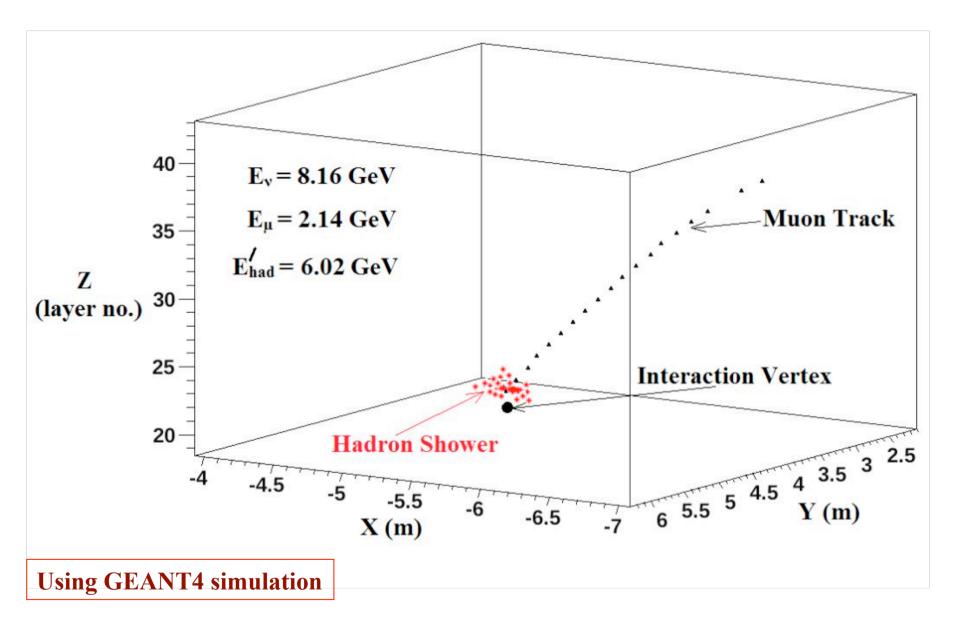
Study Atmospheric neutrinos w/ a wide range of Baselines & Energies

Recent discovery of large θ_{13} : A great news for ICAL-INO

What do we want to achieve?

- * Reconfirm neutrino oscillations using neutrinos and anti-neutrinos separately
- ❖ Improved precision of atmospheric oscillation parameters
- ❖ Determine neutrino mass hierarchy using matter effects via charge discrimination
- * Measure the deviation of 2-3 mixing angle from its maximal value and its octant
- * Test bed for various new physics like NSI, CPT violation, long range forces
- ❖ Detect Ultra High Energy Neutrinos, Cosmic Muons, Indirect searches of DM

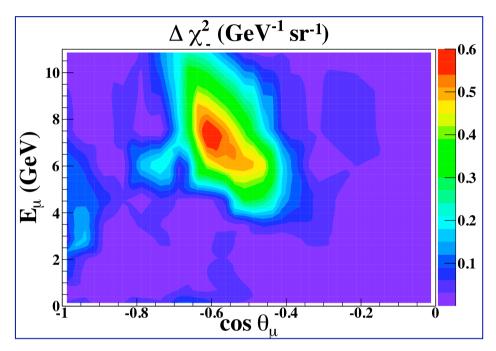
Event Display Inside the ICAL Detector



Devi, Thakore, Agarwalla, Dighe, arXiv:1406.3689 [hep-ph] (INO Collaboration)

Neutrino Mass Hierarchy Discrimination

Distribution of $\Delta \chi^2$ [χ^2 (IH) - χ^2 (NH)] for mass hierarchy discrimination considering μ events



 $\Delta \chi^{2}_{-} \text{ (GeV}^{-1} \text{ sr}^{-1})$ 0.6 0.5 0.3 0.2 0.1

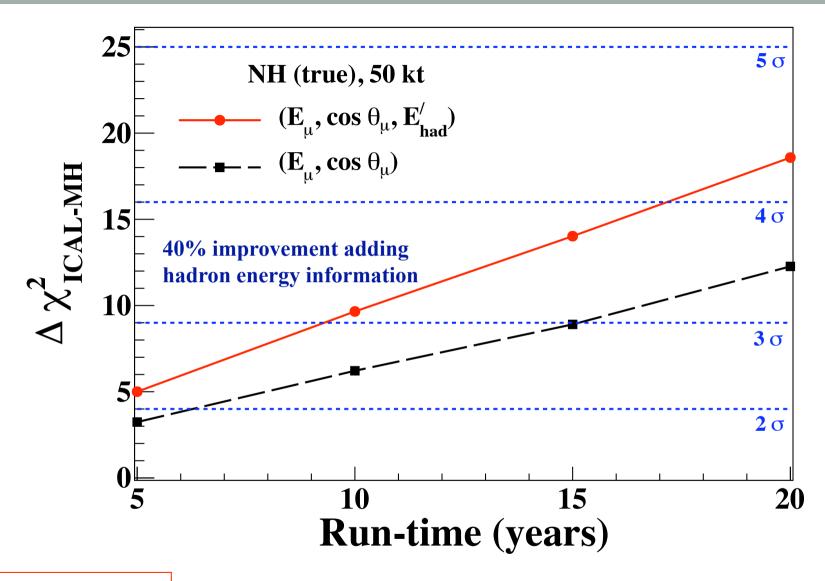
Hadron energy information not used

Observable Bin width Total bins Range [1, 4)0.56 E_{μ} (GeV) 3 [4, 7)10 [7, 11)12 [-1.0, -0.4)0.05 $\cos \theta_{\mu}$ [-0.4, 0.0)0.1 21 [0.0, 1.0]0.25 2 [0, 2)1 E'_{had} (GeV) [2, 4) 2 [4, 15)11

Hadron energy information used

- Further subdivide the events into four hadron energy bins
- Hadron energy carries crucial information
- Correlation between hadron energy and muon momentum is very important

Identifying Neutrino Mass Hierarchy with ICAL

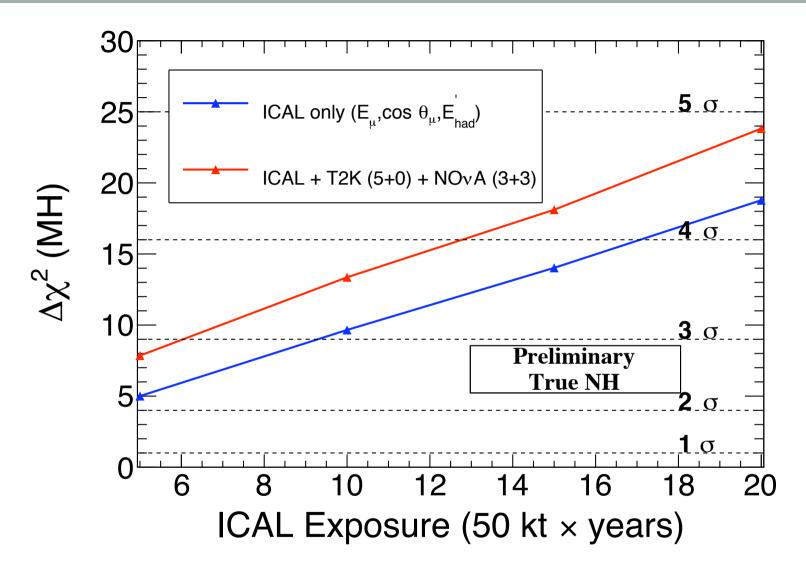


Median Sensitivity

Devi, Thakore, Agarwalla, Dighe, arXiv:1406.3689 [hep-ph] (INO Collaboration)

50 kt ICAL can rule out the wrong hierarchy with $\Delta \chi^2 \approx 9.5$ in 10 years

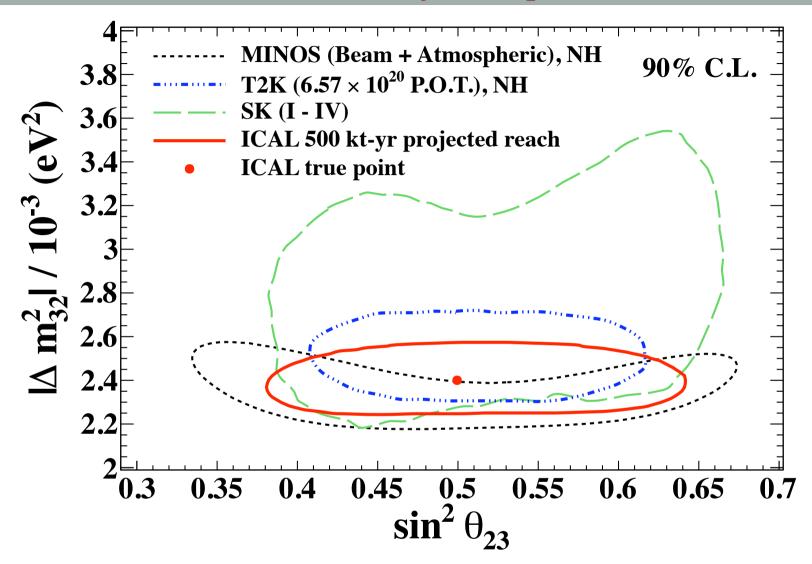
MH Discovery with ICAL+T2K+NOvA



Devi, Thakore, Agarwalla, Dighe, work in progress (INO Collaboration)

3σ median sensitivity can be achieved in 6 years

Precision Measurement of Atmospheric Parameters



Devi, Thakore, Agarwalla, Dighe, arXiv:1406.3689 [hep-ph] (INO Collaboration)

ICAL expected precision on atmospheric mass splitting is far superior than SK

Concluding Remarks

Recent discovery of θ_{13} signifies an important breakthrough in establishing the standard three flavor oscillation picture of neutrinos

It has opened up exciting possibilities for current & future oscillation experiments

At present, we have:

$$|U|_{\text{LEP}(3\sigma)} = \begin{pmatrix} 0.799 \to 0.844 & 0.515 \to 0.581 & 0.129 \to 0.173 \\ 0.212 \to 0.527 & 0.426 \to 0.707 & 0.598 \to 0.805 \\ 0.233 \to 0.538 & 0.450 \to 0.722 & 0.573 \to 0.787 \end{pmatrix}$$

Satisfactory progress in last 15 years but still very far from the 'dream' precision:

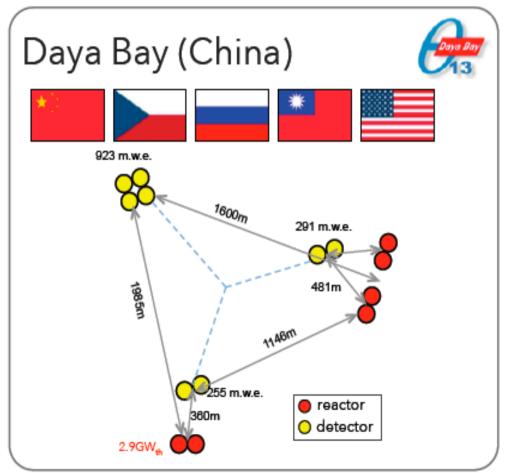
$$|V|_{\text{CKM}} = \begin{pmatrix} 0.97427 \pm 0.00015 & 0.22534 \pm 0.0065 & (3.51 \pm 0.15) \times 10^{-3} \\ 0.2252 \pm 0.00065 & 0.97344 \pm 0.00016 & (41.2^{+1.1}_{-5}) \times 10^{-3} \\ (8.67^{+0.29}_{-0.31}) \times 10^{-3} & (40.4^{+1.1}_{-0.5}) \times 10^{-3} & 0.999146^{+0.000021}_{-0.000046} \end{pmatrix}$$

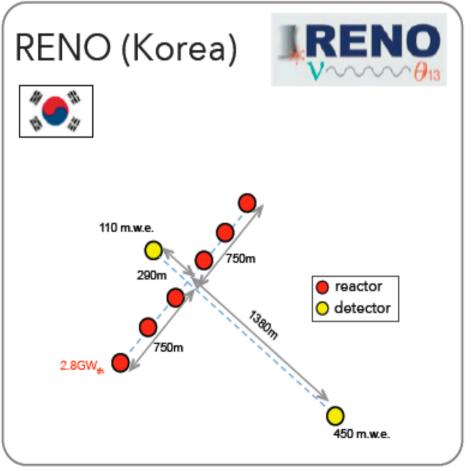
!! Let us work together and achieve it !!

Thank you!

Backup Slides: Currently Running Reactor θ_{13} Experiments





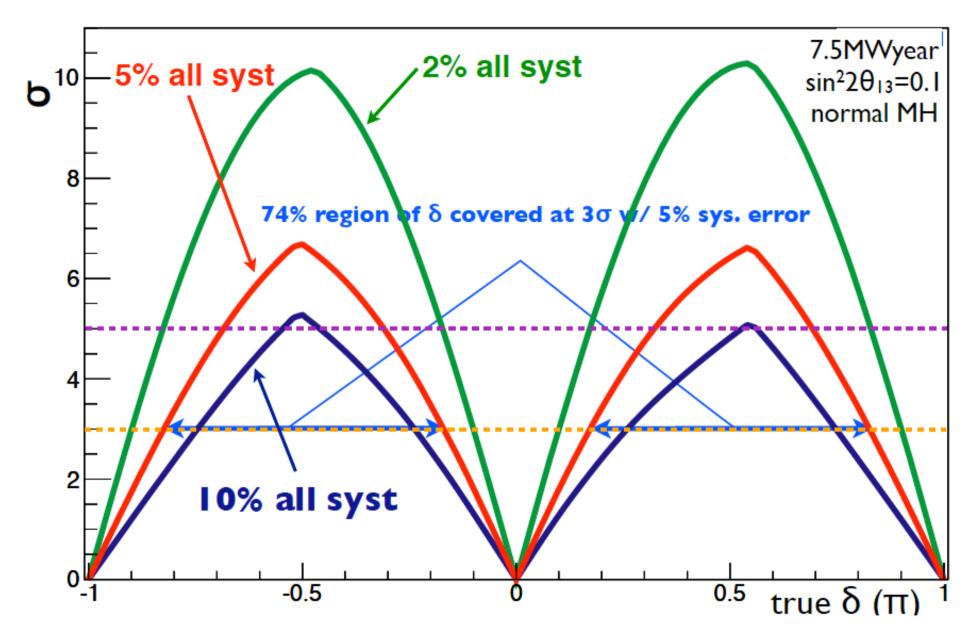


Backup Slides: Key Features of three Reactor Experiments

Experiment	Double Chooz	Daya Bay	RENO	
# of reactors (total power)	2 (9.4 GW)	3 (17.4 GW)	6 (16.8 GW)	
Reactor configuration	2	3	6 inline	
Detector configuration	1 near + 1 far	2 near + 1 far	1 near + 1 far	
Baseline [m]	(400, 1050)	(364, 480, 1912)	(290, 1380)	
Overburden [m.w.e.]	(120, 300)	(280, 300, 880)	(120, 450)	
Target mass [ton]	(8.3, 8.3)	(40, 40, 80)	(16, 16)	
Detector geometry	Cylindrical detector (Gd-LS, γ-catcher, buffer)			
Outer shield	0.5m of LS & 0.15 m of steel	2.5m water	1.5m of water	
Muon veto system	LS & Scinti-Strip	Water Cerenkov & RPC	Water Cerenkov	
Designed sensitivity (90% C.L.)	~0.03	~0.01	~0.02	

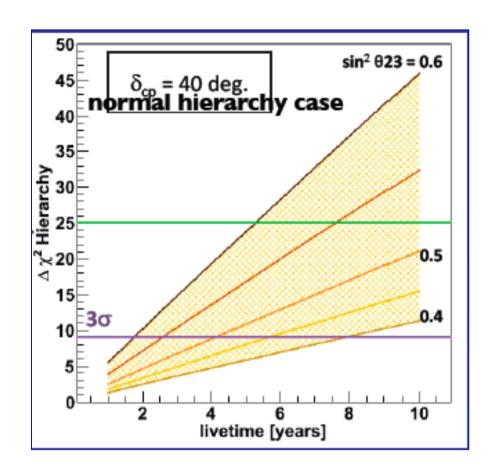
Daya Bay Strategy: Go strong, big and deep!

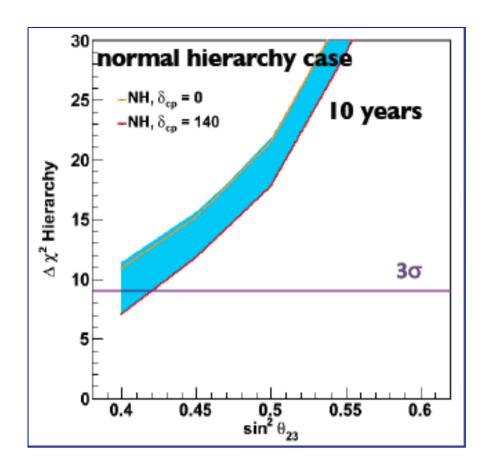
Backup Slides: CPV Discovery in T2HK Setup (w/ MH known)



Hyper-Kamiokande, Letter of Intent, arXiv:1109.3262 [hep-ex]

Backup Slides: MH Discovery in T2HK combining Atmospheric v





 3σ hierarchy discrimination for $\sin^2\theta_{23} > 0.42$ in case of normal hierarchy

Hyper-Kamiokande, Letter of Intent, arXiv:1109.3262 [hep-ex]

Backup Slides

The New Minimal Standard Model

- Minimal Extensions to give Mass to the Neutrino:
 - * Introduce ν_R AND impose L conservation \Rightarrow Dirac $\nu \neq \nu^c$:

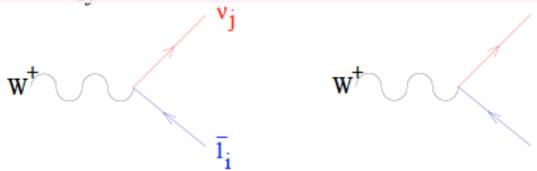
$$\mathcal{L} = \mathcal{L}_{SM} - M_{\nu} \overline{\nu_L} \nu_R + h.c.$$

* NOT impose L conservation \Rightarrow Majorana $\nu = \nu^c$

$$\mathcal{L} = \mathcal{L}_{SM} - \frac{1}{2} M_{\nu} \overline{\nu_L} \nu_L^C + h.c.$$

The charged current interactions of leptons are not diagonal (same as quarks)

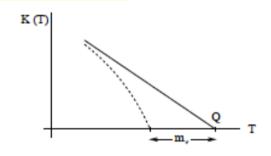
$$\frac{g}{\sqrt{2}} W_{\mu}^{+} \sum_{ij} \left(U_{\rm LEP}^{ij} \, \overline{\ell^{i}} \, \gamma^{\mu} \, L \, \nu^{j} \, + \, U_{\rm CKM}^{ij} \, \overline{U^{i}} \, \gamma^{\mu} \, L \, D^{j} \right) \, + \, h.c.$$



Backup Slides

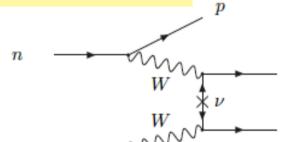
Neutrino Mass Scale

Single β decay: Dirac or Majorana ν mass modify spectrum endpoint



$$m_{\nu_e}^2 = \sum \! m_j^2 |U_{ej}|^2 = c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2$$

 ν -less Double- β decay: \Leftrightarrow Majorana $\nu's$ sensitive to Majorana phases



If
$$m_{
u}$$
 only source of $\Delta L \ (T_{1/2}^{0
u})^{-1} \propto (m_{ee})^2$

$$m_{ee} = |\sum U_{ej}^2 m_j|$$

$$\begin{aligned} & _{e^{-}} \ m_{ee} = |\sum U_{ej}^{2} m_{j}| \\ & _{e^{-}} = \left| c_{13}^{2} c_{12}^{2} m_{1} \, e^{i \eta_{1}} + c_{13}^{2} s_{12}^{2} m_{2} \, e^{i \eta_{2}} + s_{13}^{2} m_{3} \, e^{-i \delta_{CP}} \right| \end{aligned}$$

COSMO Neutrino mass (Dirac or Majorana) modify the growth of structures



Backup Slides (Neutrinoless double beta decay)

Experimental Limits

Isotope	0vββ half life	Experiment	<m> eV</m>
⁴⁸ Ca	> 1.4*10 ²² (90%CL)	ELEGANT-VI	< 7 - 44
⁷⁶ Ge	> 1.9*10 ²⁵ (90%CL)	Heidelberg-Moscow	< 0.35
⁷⁶ Ge	2230+440 ₋₃₁₀ (90%CL)	Subset of HM coll.	0.32 +/- 0.03
⁷⁶ Ge	> 2.1*10 ²⁵ (90%CL)	GERDA†	< 0.2 – 0.4
⁸² Se	> 2.1*10 ²³ (90%CL)	NEMO-3	<1.2 – 3.2
¹⁰⁰ Mo	> 5.8*10 ²³ (90%CL)	NEMO-3	< 0.6 – 2.7
¹¹⁶ Cd	> 1.7*10 ²³ (90%CL)	Solotvino	< 1.7
¹³⁰ Te	> 2.8*10 ²⁴ (90%CL)	Cuoricino	< 0.41 - 0.98
¹³⁶ Xe	> 1.9*10 ²⁵ (90%CL)	KamLAND-Zen††	< 0.12 - 0.25
¹³⁶ Xe	> 1.6×10 ²⁵ (90%CL)	EXO-200†††	< 0.14 - 0.38
¹⁵⁰ Nd	> 1.8*10 ²² (90%CL)	NEMO-3	

Courtesy to Liang Yang

[F. Avignone, S. Elliot, J. Engel, arXiv:0708: 1033v2 (2007)]

† [GERDA Collaboration, arXiv:1307.4720 (2013]

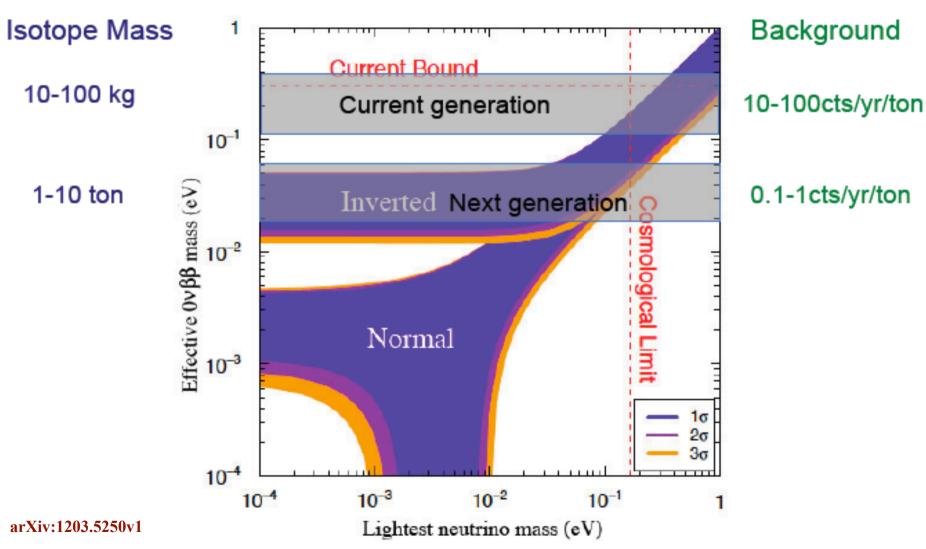
†† [KamLAND-Zen Collaboration, Phys. Rev. Lett. 110, 062502(2013)]

††† [EXO Collaboration, Phys. Rev. Lett.109, 0322505 (2012)]



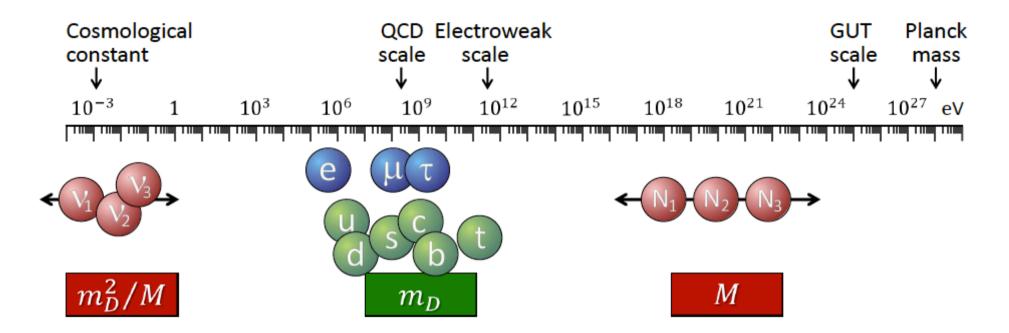
Backup Slides (Neutrino Mass)

Experimental Sensitivity to Neutrino Mass



Courtesy to Liang Yang

Backup Slides (See-Saw & Neutrino Mass)



Mass matrix for one family of ordinary and heavy r.h. neutrinos

$$(\overline{\nu}_L, \overline{N}_R)\begin{pmatrix} \mathbf{0} & m_D \\ m_D & \mathbf{M} \end{pmatrix}\begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

Diagonalization

$$(\overline{\nu}_L, \overline{N}_R)\begin{pmatrix} m_D^2/M & 0 \\ 0 & M \end{pmatrix}\begin{pmatrix} \nu_L \\ N_R \end{pmatrix}$$

One light and one heavy Majorana neutrino

