

Quantum Computing with Para-hydrogen

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Joint work with: J.A. Jones (Oxford) and S.B. Duckett (York)







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Entanglement in two qubit mixed states

Entanglement: Pure States

Separable

$$|\psi\rangle_{AB} = |\phi\rangle_A \otimes |\chi\rangle_B$$

$$|0\rangle \otimes |0\rangle = |00\rangle$$

$$|0\rangle \otimes (|0\rangle + |1\rangle) / \sqrt{2} = (|00\rangle + |01\rangle) / \sqrt{2}$$

$$(|0\rangle + |1\rangle) / \sqrt{2} \otimes (|0\rangle + i|1\rangle) / \sqrt{2} = (|00\rangle + i|01\rangle + |10\rangle + i|11\rangle) / 2$$

Entangled

$$|\psi\rangle_{AB} \neq |\phi\rangle_A \otimes |\chi\rangle_B$$

$$|\phi^+\rangle = (|00\rangle + |11\rangle) / \sqrt{2}$$

$$|\phi^-\rangle = (|00\rangle - |11\rangle) / \sqrt{2}$$

$$|\psi^+\rangle = (|01\rangle + |10\rangle) / \sqrt{2}$$

$$|\psi^-\rangle = (|01\rangle - |10\rangle) / \sqrt{2}$$

Bell or EPR States

Entanglement and Density Matrices 1

Density matrix ρ

$$\rho_{AB} = \sum_k \lambda_k |\psi\rangle_{AB} \langle\psi|_{AB}$$

Pure and mixed states

Density matrices of the Bell States

$$(|00\rangle + |11\rangle) / \sqrt{2} \longrightarrow \rho = \begin{pmatrix} 1/2 & 0 & 0 & 1/2 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1/2 & 0 & 0 & 1/2 \end{pmatrix}$$

$$(|01\rangle - |10\rangle) / \sqrt{2} \longrightarrow \rho = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1/2 & -1/2 & 0 \\ 0 & -1/2 & 1/2 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Entanglement and Density Matrices 2

Maximally mixed state

$$\rho = \begin{pmatrix} 1/4 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 0 \\ 0 & 0 & 1/4 & 0 \\ 0 & 0 & 0 & 1/4 \end{pmatrix} = (|00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 10| + |11\rangle\langle 11|) / 4$$

$$= (|0\rangle\langle 0| \otimes |0\rangle\langle 0| + |0\rangle\langle 0| \otimes |1\rangle\langle 1| + |1\rangle\langle 1| \otimes |0\rangle\langle 0| + |1\rangle\langle 1| \otimes |1\rangle\langle 1|) / 4$$

Another mixed state

$$\rho = \begin{pmatrix} 3/4 & 1/4 & 0 & 0 \\ 1/4 & 1/4 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = (|0\rangle\langle 0| \otimes |0\rangle\langle 0| + |0\rangle\langle 0| \otimes (|0\rangle + |1\rangle)(\langle 0| + \langle 1|) / 2) / 2$$

Entanglement and Density Matrices 3

Separable density matrix

$$\rho_{AB} = \sum_k \eta_k (\rho_A \otimes \rho_B)_k$$

Convex sum of direct product states

Entangled density matrix

$$\rho_{AB} \neq \sum_k \eta_k (\rho_A \otimes \rho_B)_k$$

Detecting and Quantifying Entanglement 1

Can we find a separable decomposition of ρ ?

$$\rho = \frac{1}{18} \begin{pmatrix} 5 & 0 & 0 & 1 \\ 0 & 4 & 0 & 0 \\ 0 & 0 & 4 & 0 \\ 1 & 0 & 0 & 5 \end{pmatrix}$$

Ensemble fallacy

$$\rho = \begin{pmatrix} 1/4 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 0 \\ 0 & 0 & 1/4 & 0 \\ 0 & 0 & 0 & 1/4 \end{pmatrix} = (|00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 10| + |11\rangle\langle 11|) / 4$$
$$= (|\varphi^+\rangle\langle\varphi^+| + |\varphi^-\rangle\langle\varphi^-| + |\psi^+\rangle\langle\psi^+| + |\psi^-\rangle\langle\psi^-|) / 4$$

D. T. Pegg, J. Jeffers J. Mod. Opt. **52**, 1835 (2005)

Detecting and Quantifying Entanglement 2

$$\rho = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 1/2 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} = (|01\rangle\langle 01| + |10\rangle\langle 10|) / 2$$
$$= (|\psi^+\rangle\langle\psi^+| + |\psi^-\rangle\langle\psi^-|) / 2$$

Test for Separability

PPT Test $|\psi^+\rangle = 1/\sqrt{2}(|00\rangle + |11\rangle)$

$$\rho = |\psi^+\rangle\langle\psi^+| = \begin{pmatrix} 1/2 & 0 & 0 & 1/2 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 1/2 & 0 & 0 & 1/2 \end{pmatrix}$$

$$\rho^T = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 0 & 1/2 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 1/2 \end{pmatrix}$$

Eigenvalues of ρ^T are $\{1/2, 1/2, 1/2, -1/2\}$.
 ρ is non-separable or entangled.

A. Peres Phys. Rev. Lett. **77**, 1413 (1996)

Separability of the Werner State

Werner or
pseudopure state

$$\rho = (1-\varepsilon)\mathbf{1}/4 + \varepsilon |\psi\rangle\langle\psi|$$

$$= \begin{pmatrix} \frac{1-\varepsilon}{4} & 0 & 0 & 0 \\ 0 & \frac{1+\varepsilon}{4} & -\frac{\varepsilon}{2} & 0 \\ 0 & -\frac{\varepsilon}{2} & \frac{1+\varepsilon}{4} & 0 \\ 0 & 0 & 0 & \frac{1-\varepsilon}{4} \end{pmatrix}$$

$$\rho^T = \begin{pmatrix} \frac{1-\varepsilon}{4} & 0 & 0 & -\frac{\varepsilon}{2} \\ 0 & \frac{1+\varepsilon}{4} & 0 & 0 \\ 0 & 0 & \frac{1+\varepsilon}{4} & 0 \\ -\frac{\varepsilon}{2} & 0 & 0 & \frac{1-\varepsilon}{4} \end{pmatrix}$$

- Eigenvalues of PPT of ρ are $\{1/4(1-3\varepsilon), 1/4(1+\varepsilon), 1/4(1+\varepsilon), 1/4(1+\varepsilon)\}$. ρ is non-separable or entangled only if $\varepsilon > 1/3$.

The problem of initialization

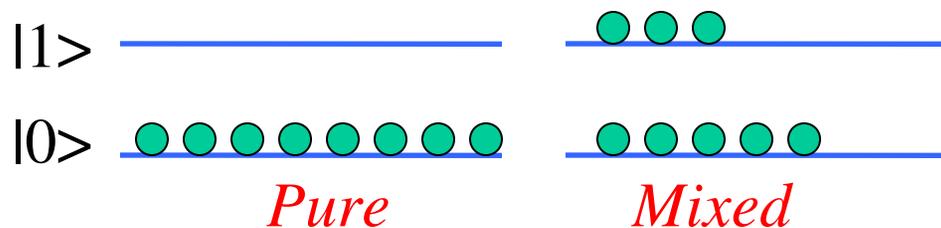
Requirements for Practical QC

- Existence of Qubits
- Initialization
- Universal Quantum Networks
- Read-out or Measure
- Decoherence

D.P. DiVincenzo quant-ph/0002077 (2000)

Initialization 1

NMR states are highly mixed.

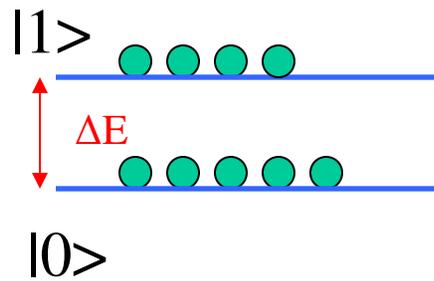


Ideal initial state

$$\rho = |0\rangle\langle 0| \otimes |0\rangle\langle 0| = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Initialization 2

Thermal equilibrium state



$$\rho_{\text{eq}} = \exp(-\beta H) / Z \text{ where } \beta = \hbar / (k_B T)$$

$$\Delta E = \gamma \hbar B_z = \hbar \omega = 2.8 \times 10^{-25} \text{ J}$$

$$k_B T = 4 \times 10^{-21} \text{ J at } 10 \text{ T, } 300 \text{ K}$$

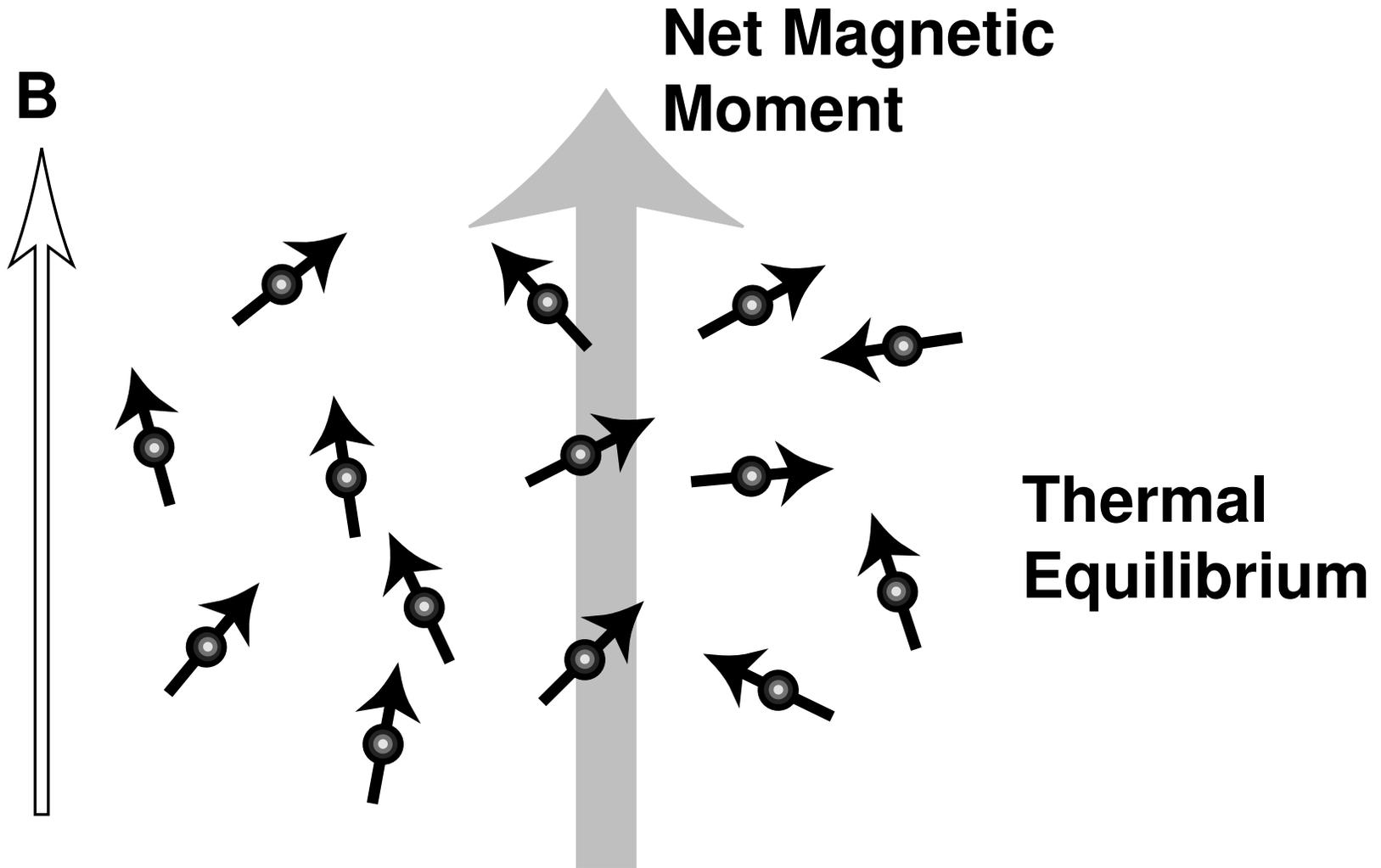
Population difference only $\sim 10^{-4}$

$$k_B T \gg \Delta E$$

$$\rho = \begin{pmatrix} 0.25001 & 0 & 0 & 0 \\ 0 & 0.25 & 0 & 0 \\ 0 & 0 & 0.25 & 0 \\ 0 & 0 & 0 & 0.24999 \end{pmatrix}$$

$$= \{ \{ 1/4 + B, 1/4, 1/4, 1/4 - B \} \} = 1/4 + B/4 (I_z + S_z) \text{ where } B = \hbar \omega / kT.$$

Initialization 3



Pseudopure States

$$\begin{pmatrix} 0.25001 & 0 & 0 & 0 \\ 0 & 0.25 & 0 & 0 \\ 0 & 0 & 0.25 & 0 \\ 0 & 0 & 0 & 0.24999 \end{pmatrix}$$

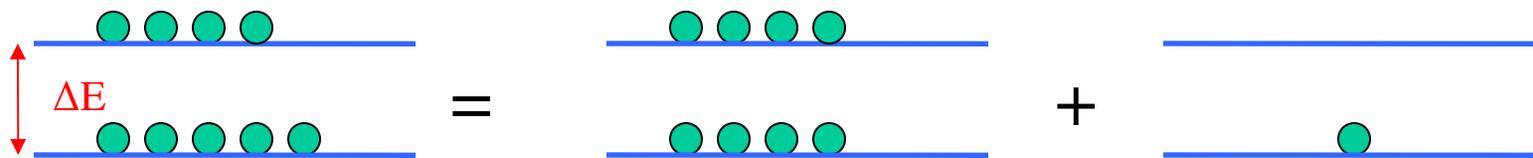
Thermal equilibrium state



$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Ideal initial state

Pseudopure state $\rho_{ps} = (1-\varepsilon)\mathbf{I}/4 + \varepsilon |\psi\rangle\langle\psi|$ ε is polarization.

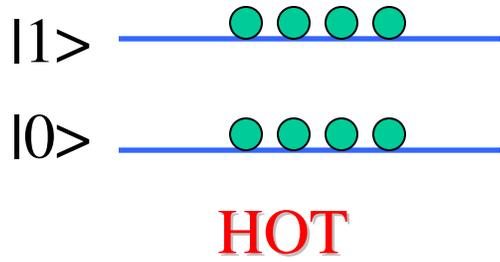
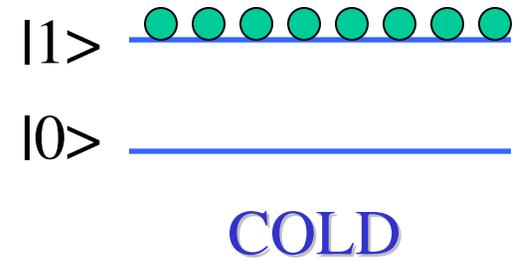
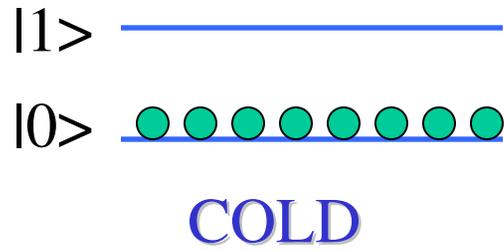


Thermal, mixed

Maximally mixed

Pure

Spin Temperature



Problems with Pseudopure States

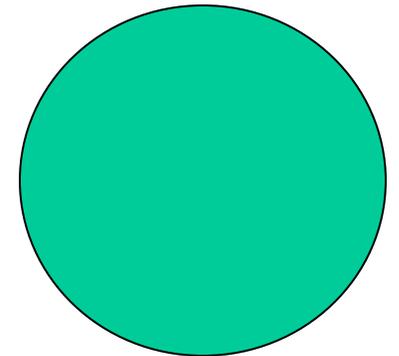
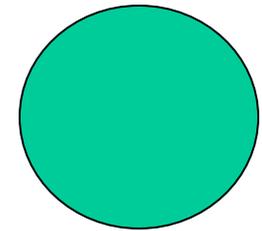
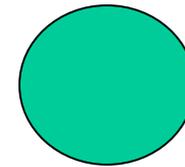
Scalability

Is NMR quantum mechanical at all?

S.L. Braunstein et al. Phys. Rev. Lett. **83**, (1999)

J.A. Jones, Fortsch. der Physik **48**, 909 (2002).

Climbing Mount Scalable

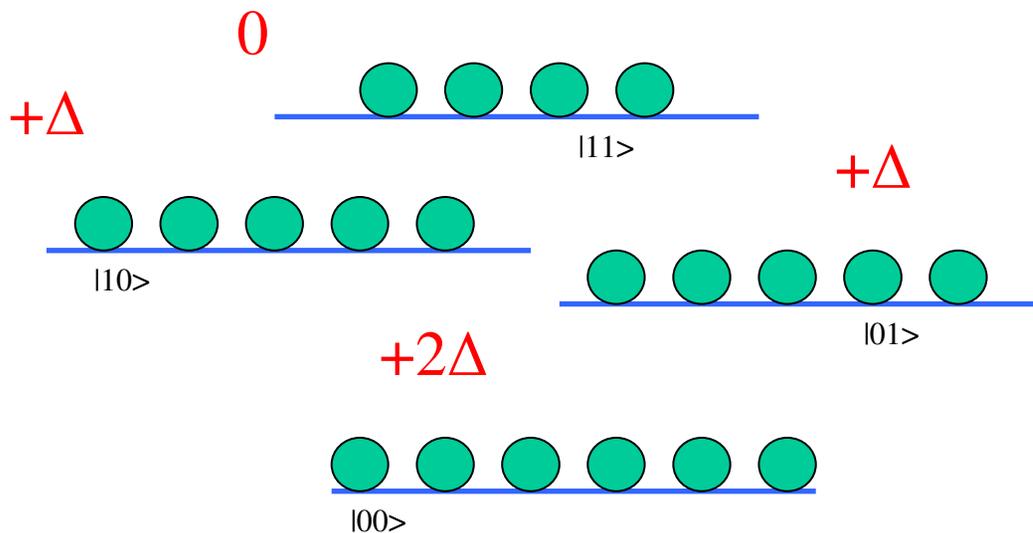


Scalability

Maximum pure state that can be extracted from thermal state

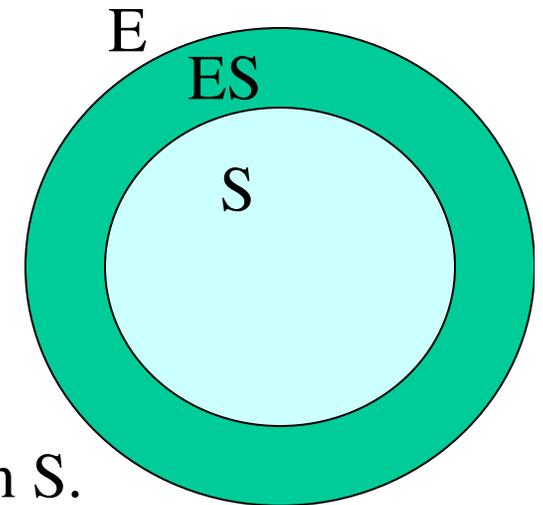
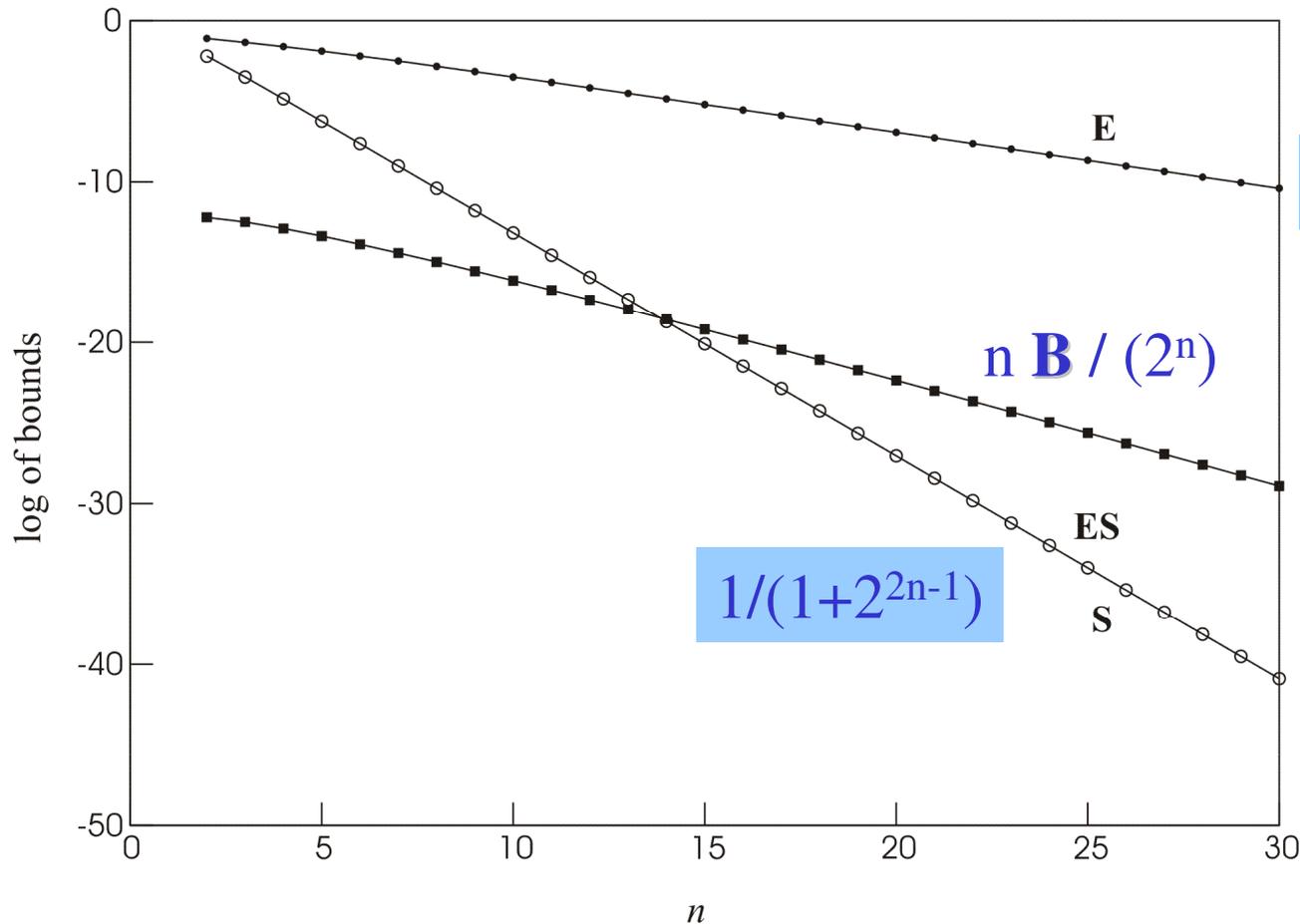
$$\rho_{ps} = (1-\varepsilon) \mathbf{1} / 2^n + \varepsilon |\psi\rangle\langle\psi|$$

Warren proposed a theoretical maximum on ε :
 $\varepsilon \sim nB / (2^n)$ where n is the number of qubits.



W.S. Warren Science **277**, 1688 (1997).

Scalability and Separability



- Present day NMR implementations work in region S.
- We can enter ES if $n > 11$ (room temperature, 10T).
- We may never enter E with thermal states!

Methods for Increasing Polarization

High fields and low temperature $B = \hbar\omega / kT$

Algorithmic concentration of polarization



Polarization transfer from nuclear spins

Polarization transfer from electron spins

Chemically induced
dynamical nuclear polarization (CIDNP)

Is an NMR Device a Quantum Computer?

No entanglement

What makes a quantum computer quantum?

States or Dynamics?

Scaling of the Hilbert space dimensionality?

Superposition?

Other models of better than classical computing.

The para-hydrogen approach

Para-hydrogen Induced Polarization 1

Need for “non-thermal” distributions

Symmetrization postulate

$$|\psi\rangle_t = |\psi\rangle_{tr} |\psi\rangle_{vib} |\psi\rangle_e |\psi\rangle_{ns} |\psi\rangle_r$$

$$|T_1\rangle = |00\rangle$$

$$|\psi\rangle_{ns} \rightarrow |T_0\rangle = 1/2^{1/2} (|01\rangle + |10\rangle)$$

$$|S_0\rangle = 1/2^{1/2} (|01\rangle - |10\rangle)$$

$$|T_{-1}\rangle = |1\rangle$$

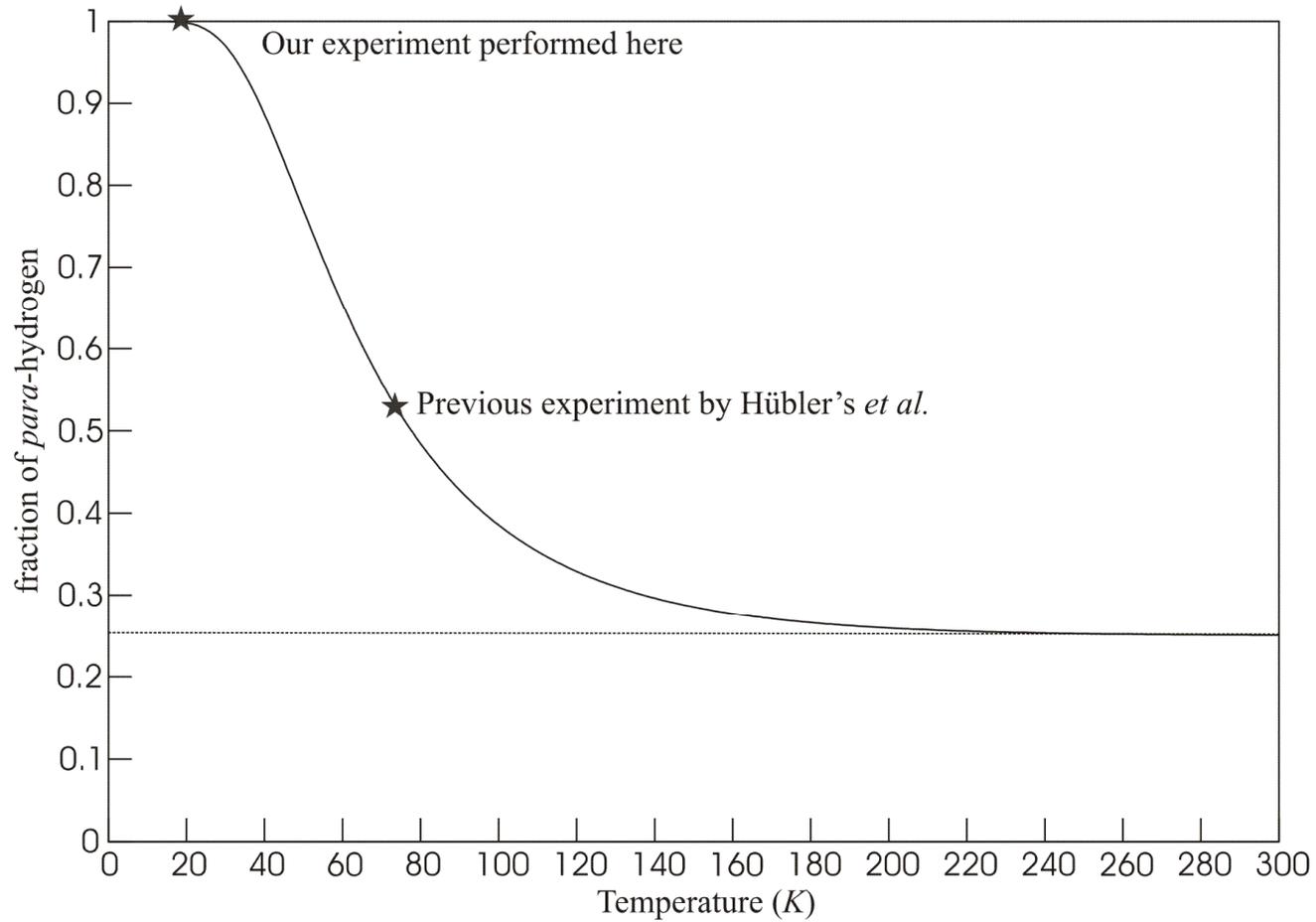
$$|\psi\rangle_r \rightarrow \begin{array}{ll} j=0, 2, 4, \dots & \text{even} \\ j=1, 3, 5, \dots & \text{odd} \end{array}$$

$$(Y_m^j(\pi-\theta, \pi+\varphi) = (-1)^j Y_m^j(\theta, \varphi))$$

Para-hydrogen Induced Polarization 2

j	$ \psi_{\text{rot}}\rangle$	$ \psi_{\text{ns}}\rangle$	o/p
0	s symmetrical	a asymmetrical	p
1	a	s	o
2	s	a	p
3	a	s	o

Making *Para*hydrogen

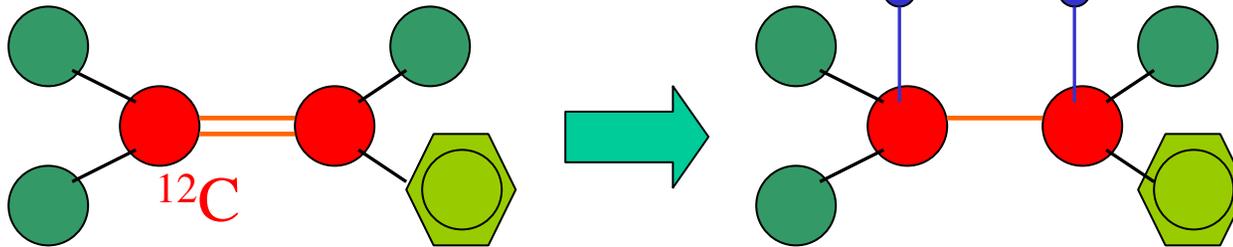


Para-hydrogen once is *para*-hydrogen forever!

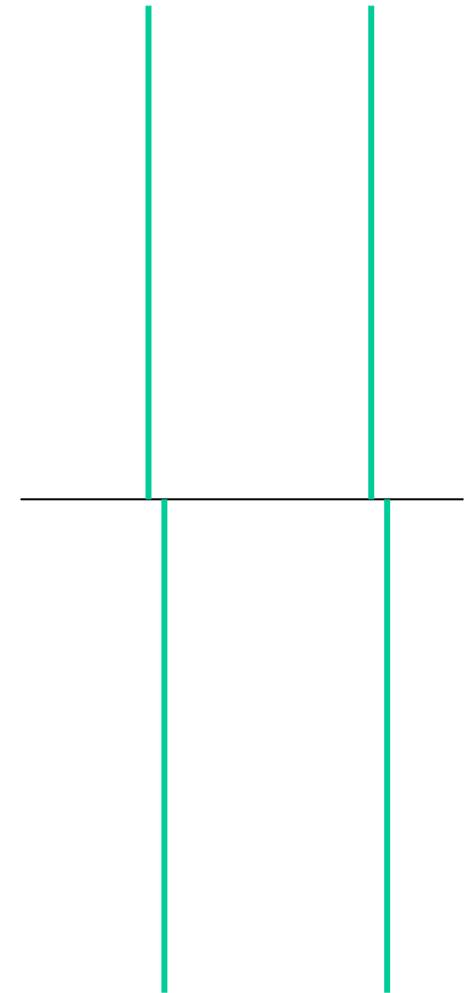
Para-hydrogen Quantum Computer

The PASADENA Experiment

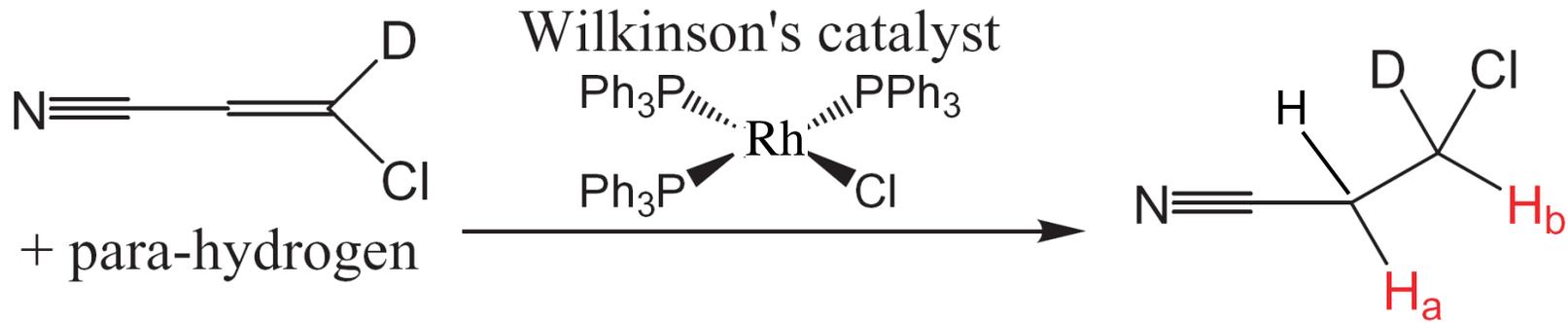
Cl



- pairwise addition
- asymmetric environment
- enhanced polarizations

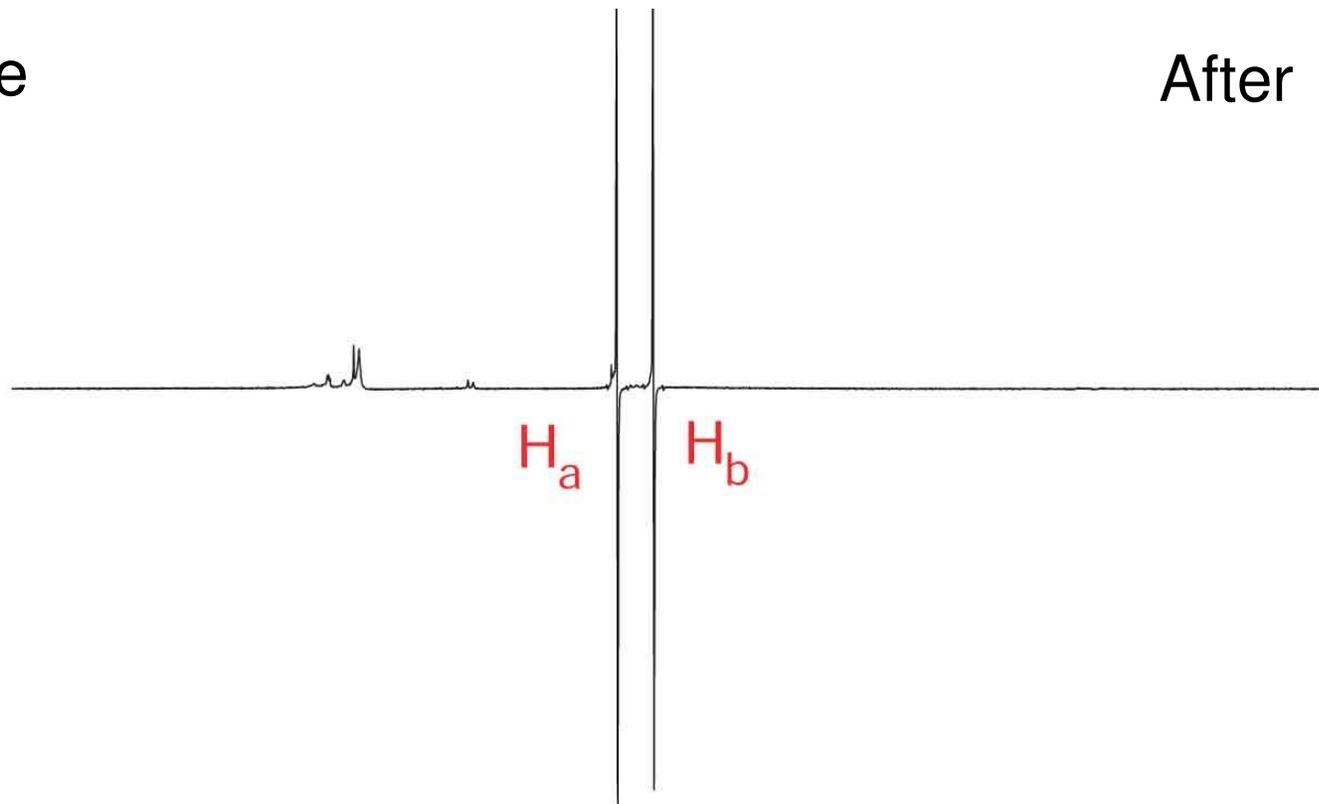


Example PASADENA spectrum

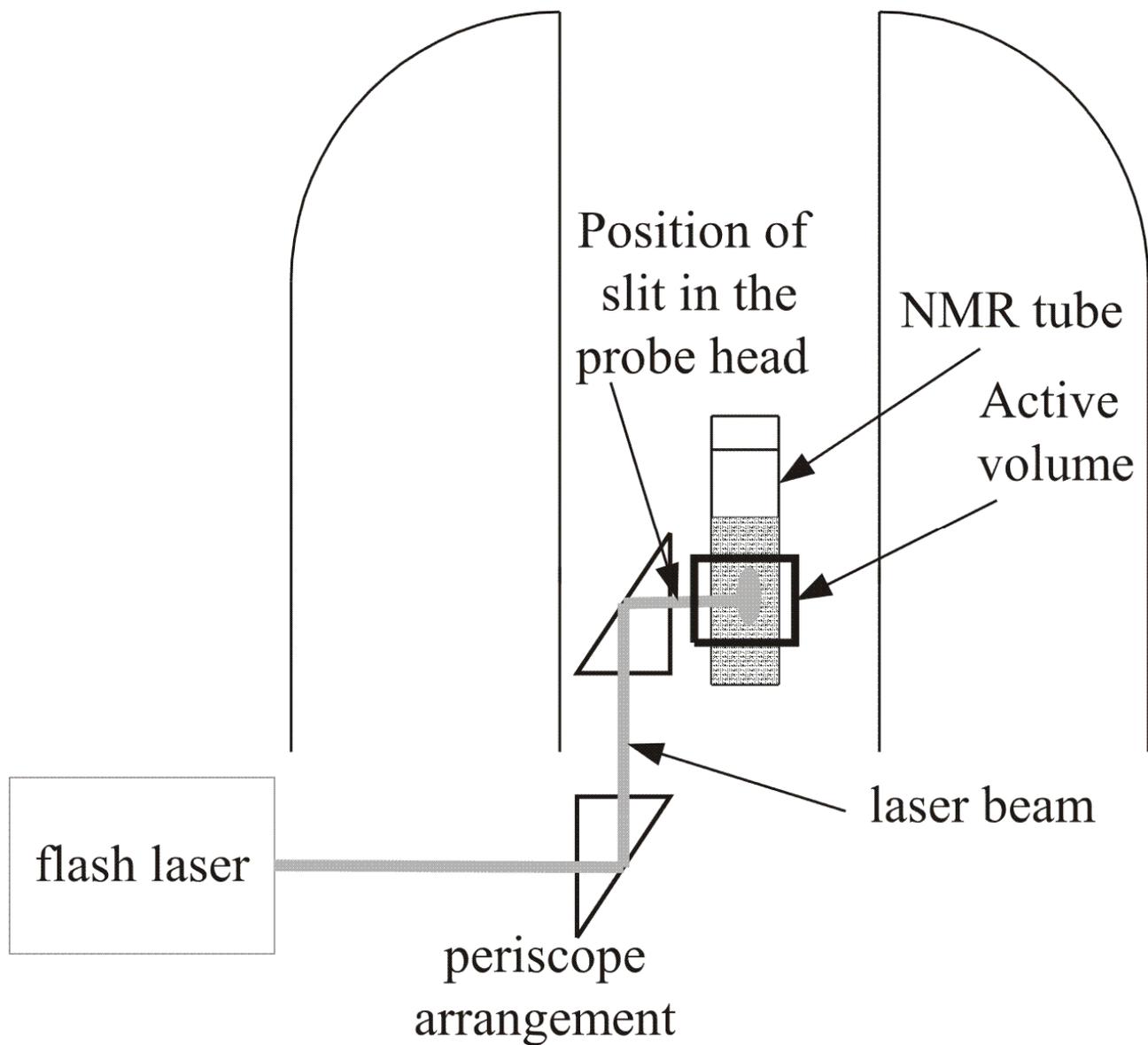


Before

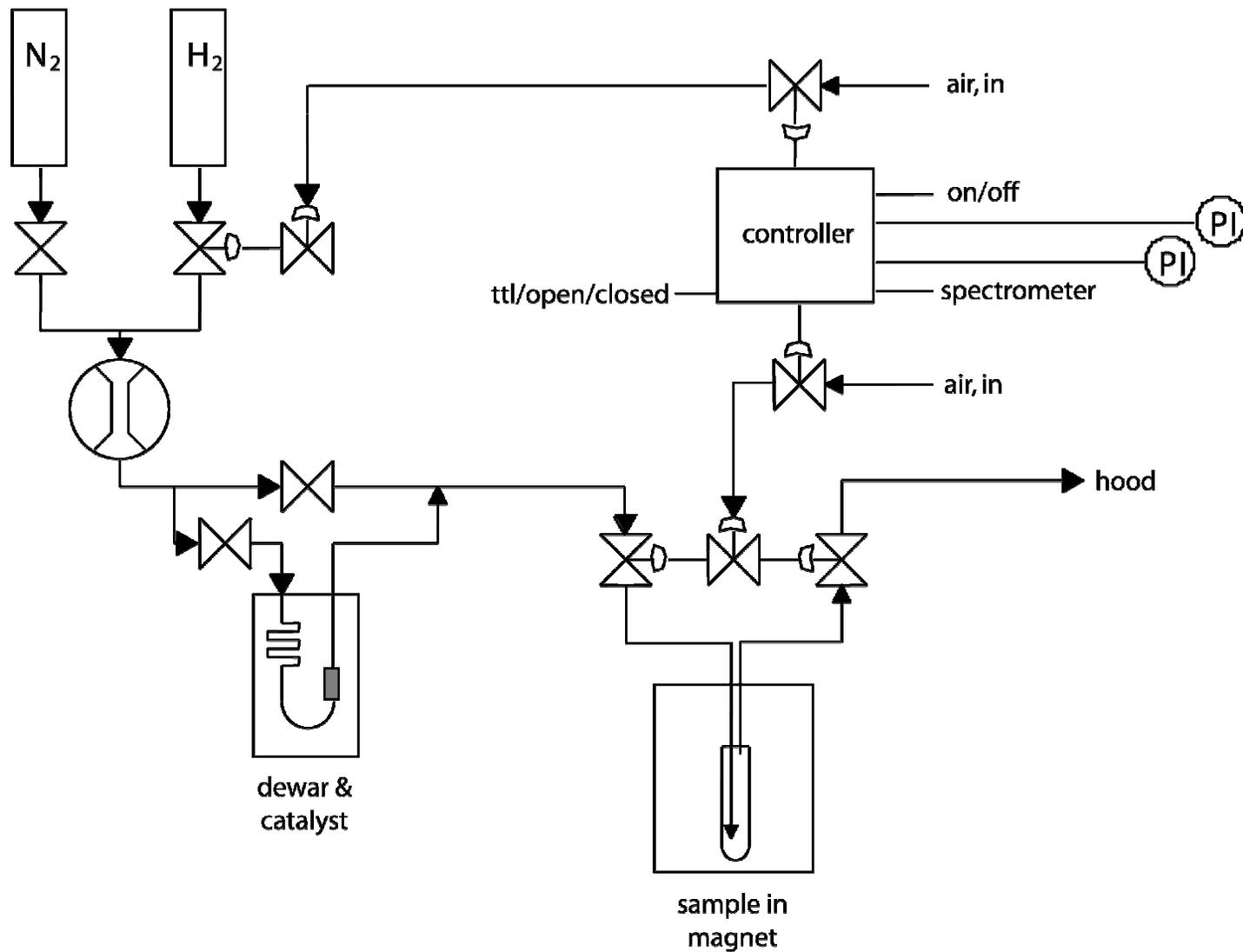
After



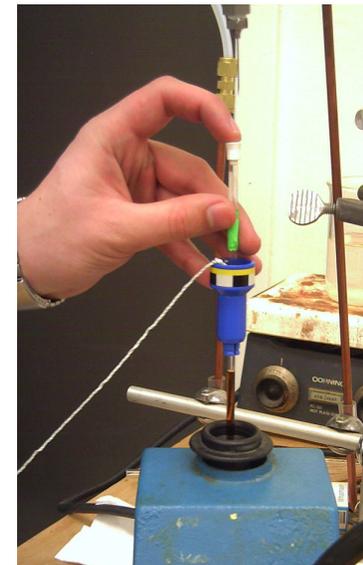
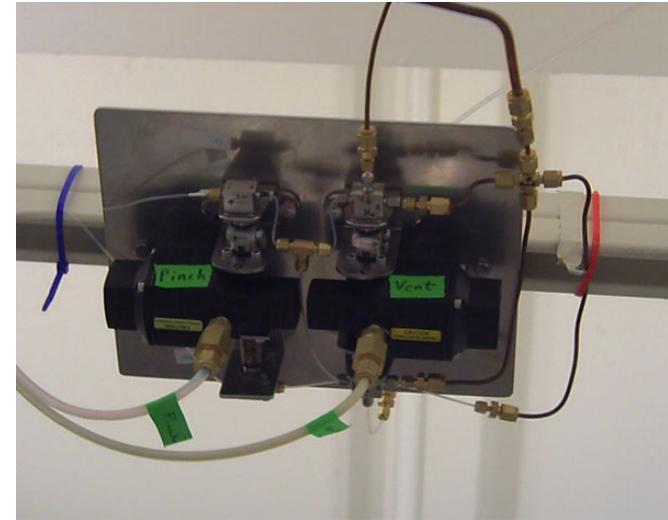
Experimental Setup



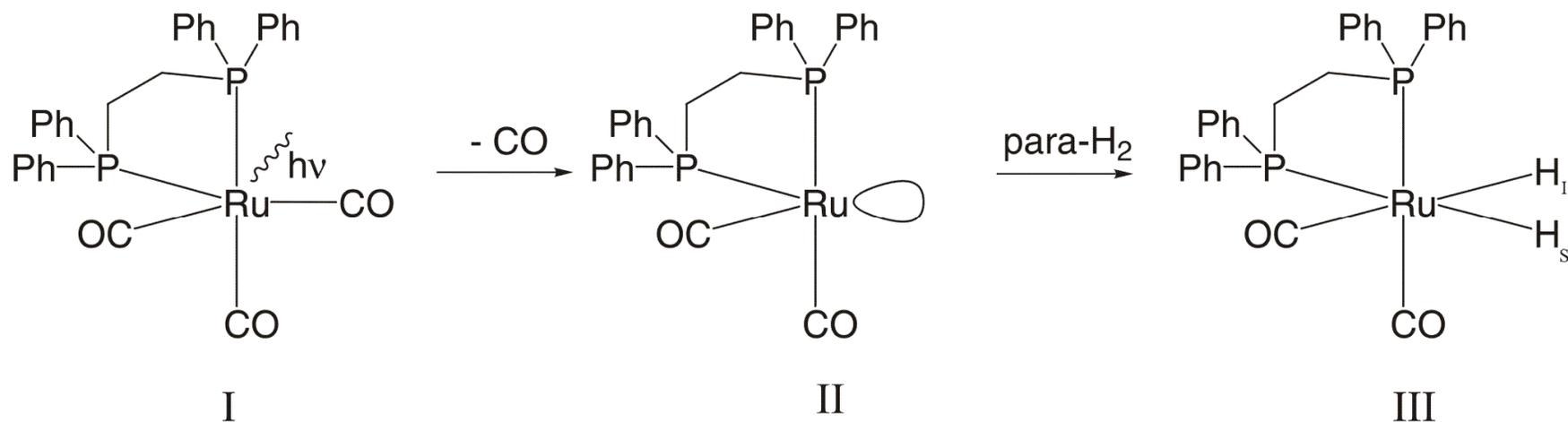
Schematic setup for the PASADENA Experiment



Some photographs depicting various units in the experiment



Flash Photolysis



D. Blazina et al. Magn. Reson. Chem. **43**, 200 (2005)

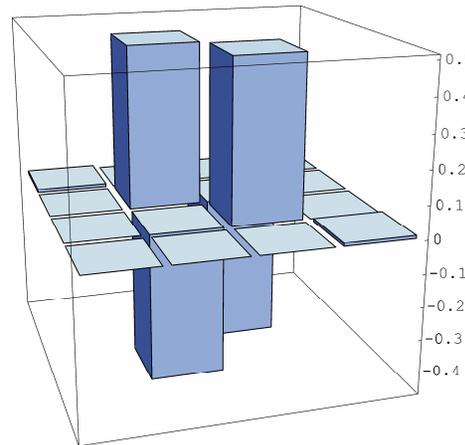
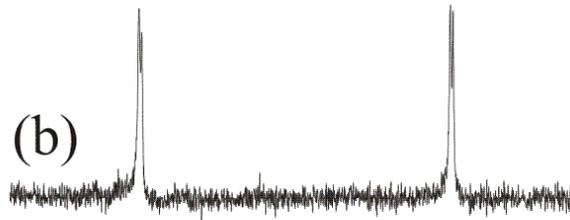
PHIP Spectra

Spin Temperature = 6.4 mK
Effective Field = 0.45 MT
 $\epsilon = 0.916$

(a)



(b)

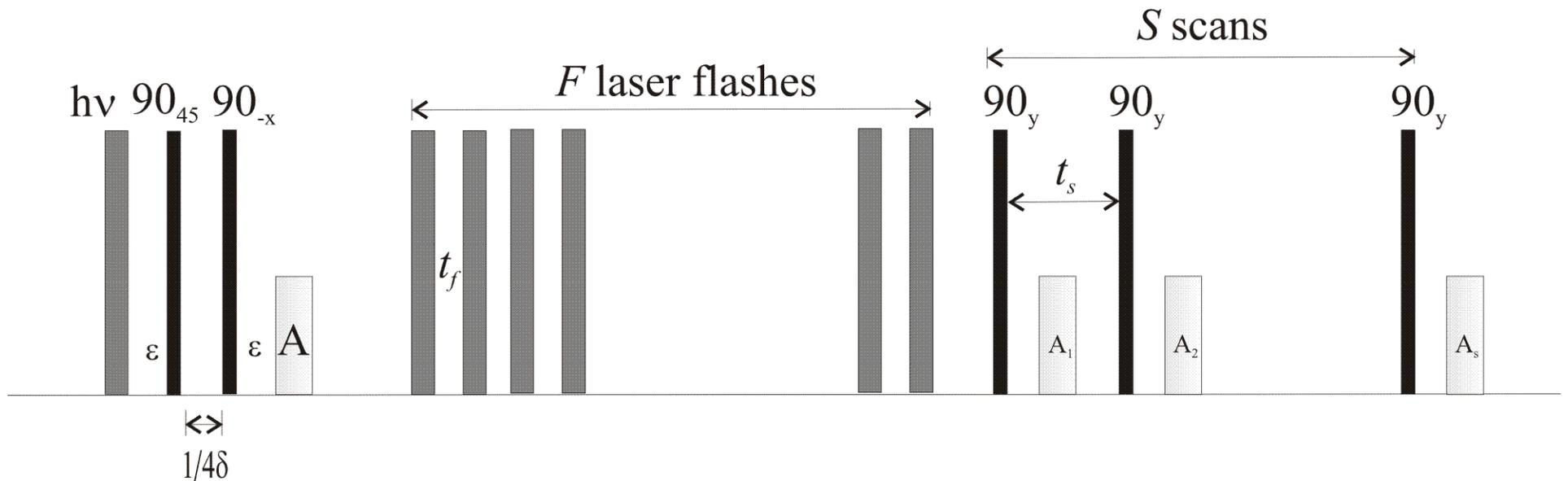


Tomography results

$$\frac{1}{2} \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & -1 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

Ideal singlet

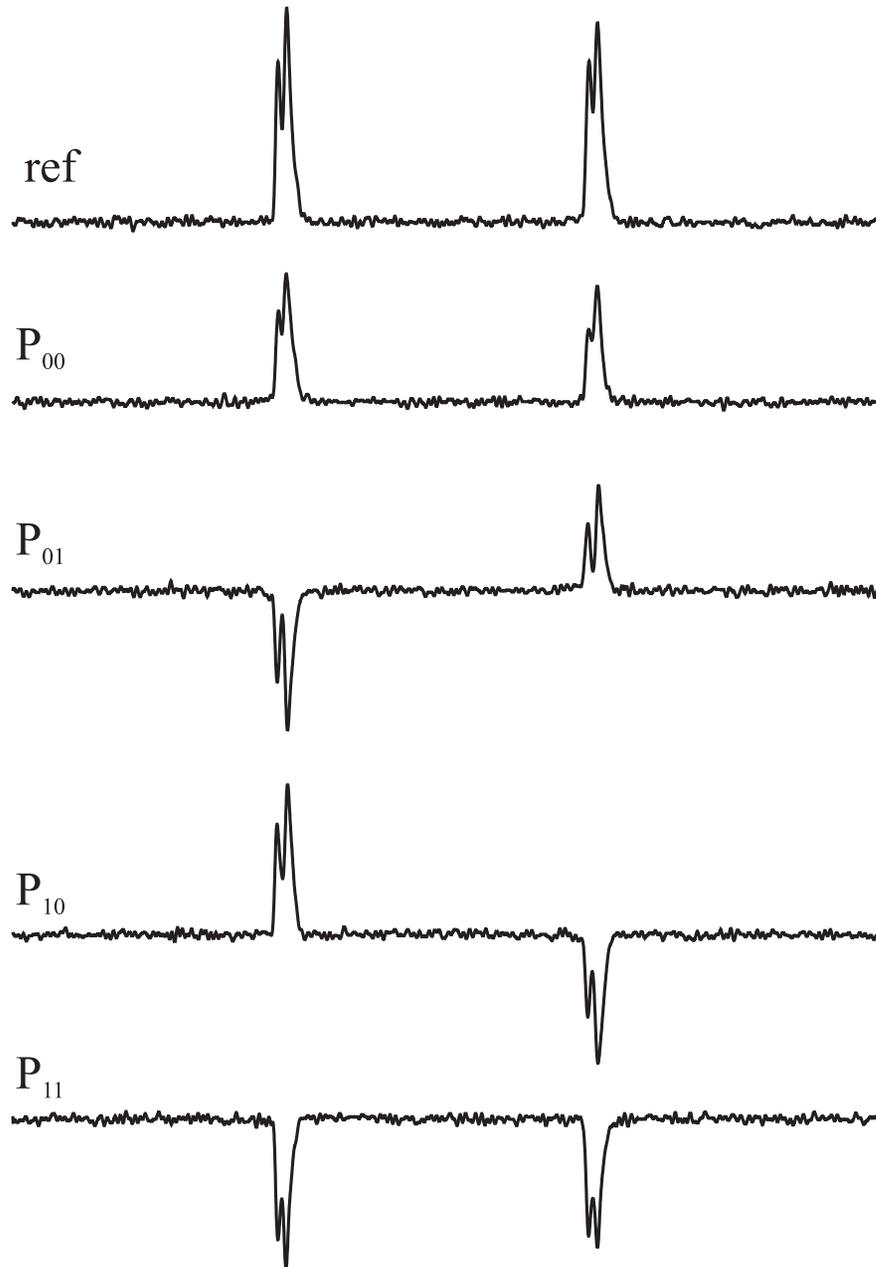
The NMR Experiment



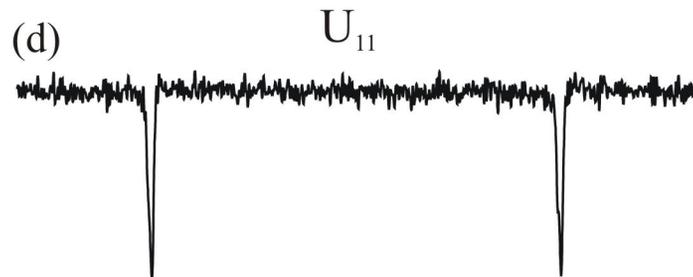
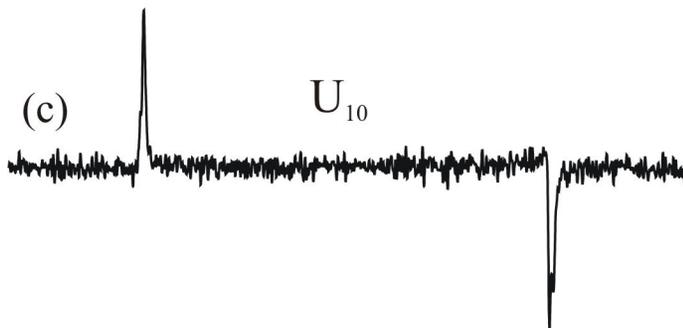
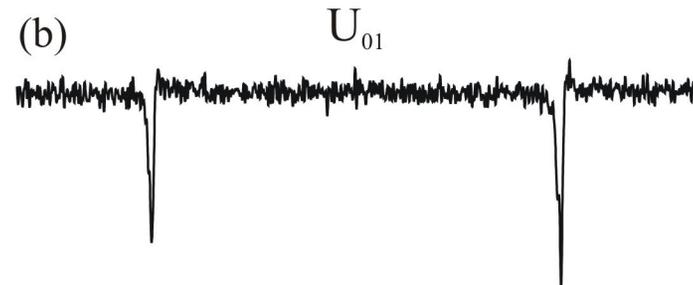
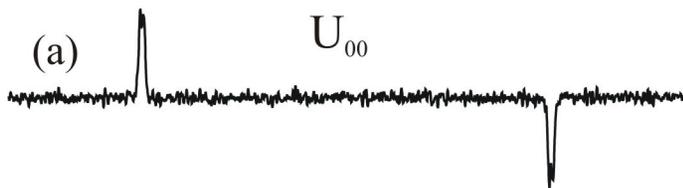
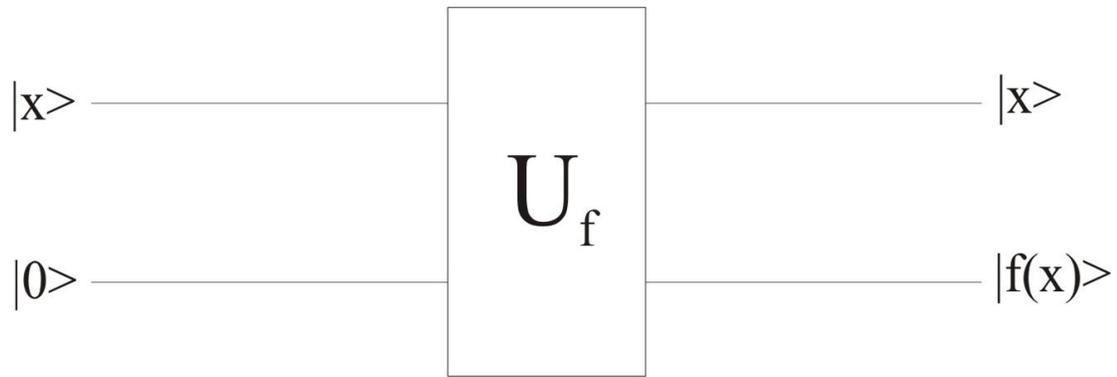
*M.S. Anwar et al. Phys. Rev. Lett. **93**, 040501 (2004)*

*M.S. Anwar et al. Phys. Rev. A **71**, 032327 (2005)*

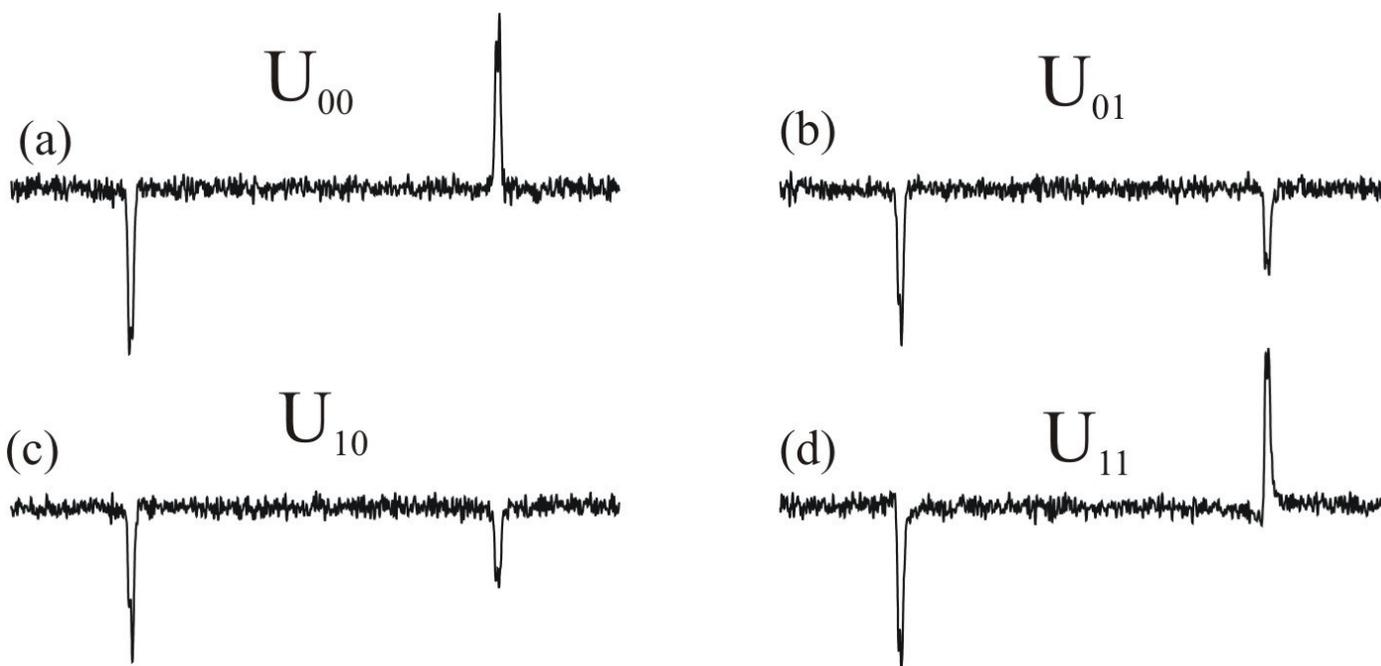
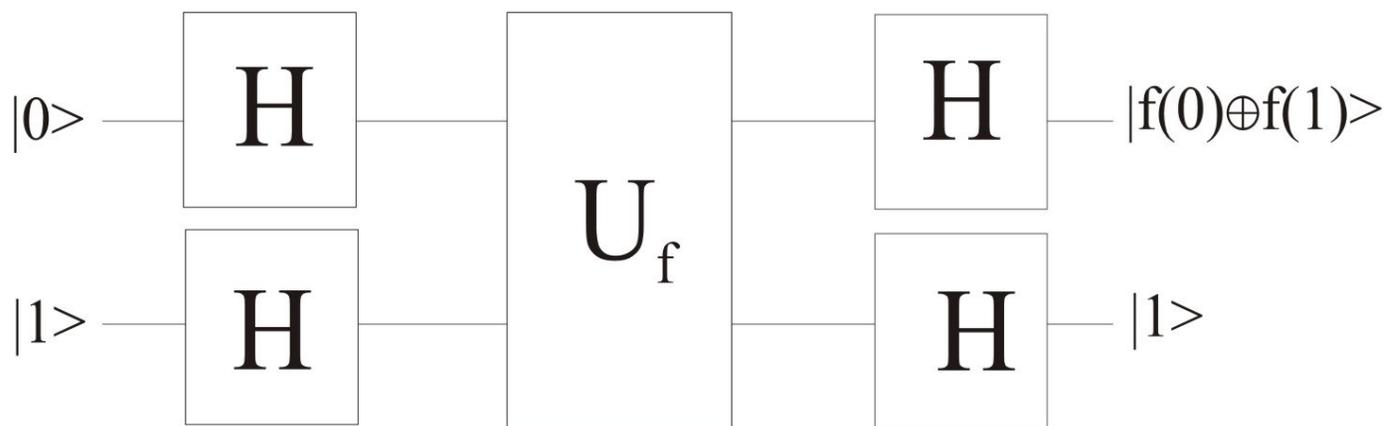
Grover's algorithm



Classical Deutsch's Algorithm: $f(1)$

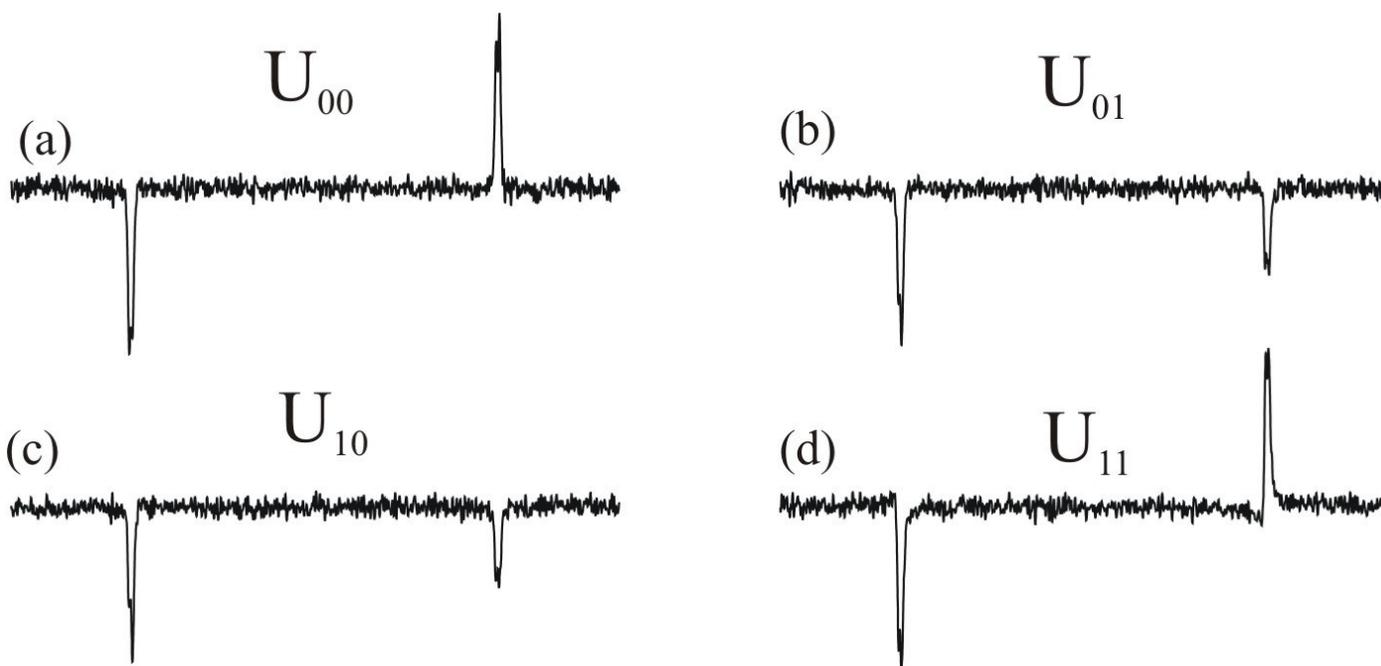
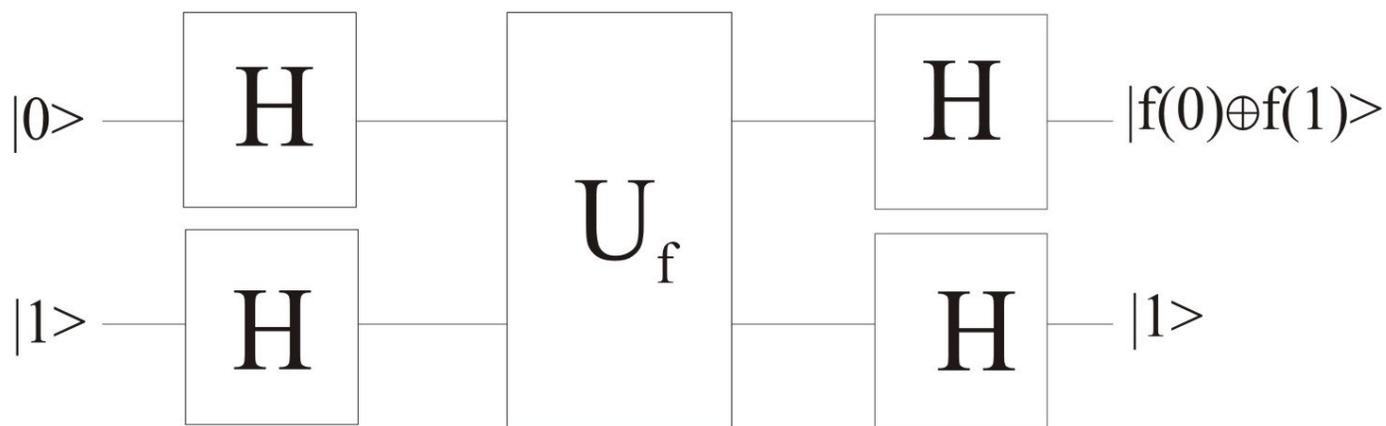


Quantum Deutsch's Algorithm



M.S. Anwar et al. Phys. Rev. A **70**, 032324 (2004)

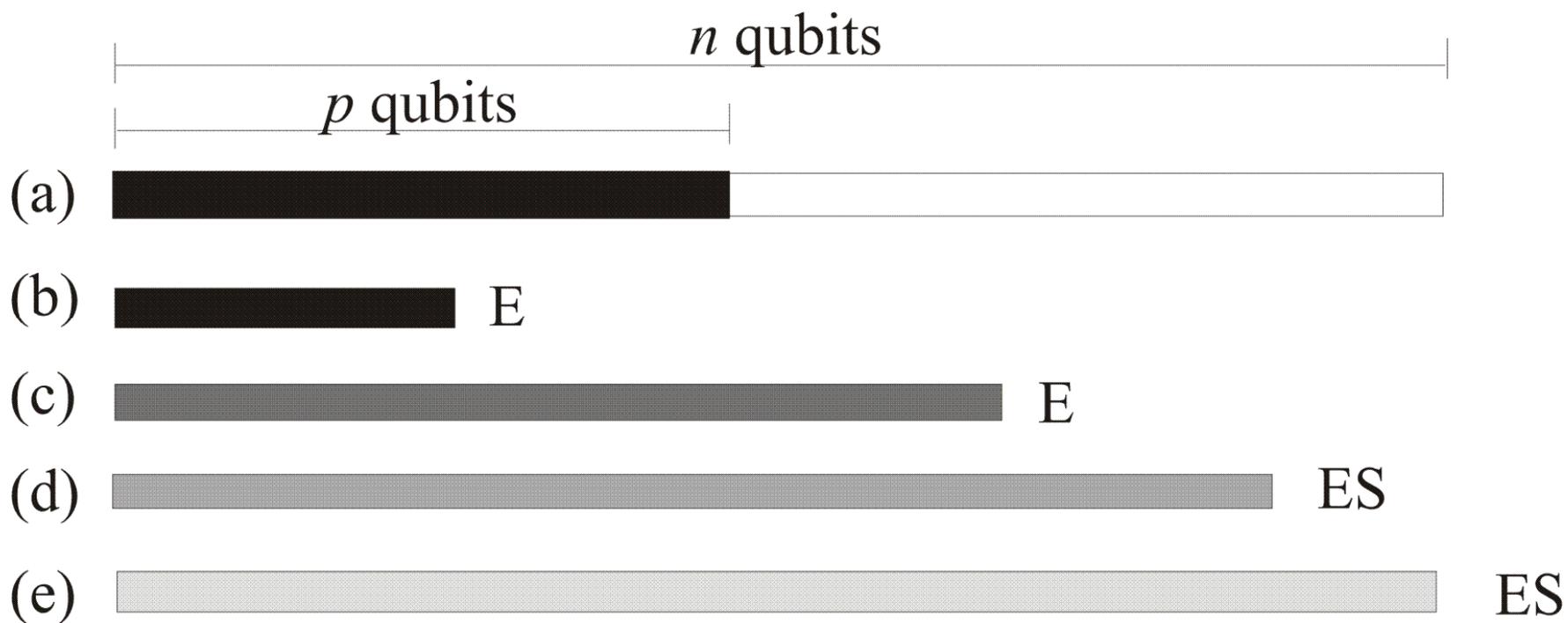
Quantum Deutsch's Algorithm



M.S. Anwar et al. Phys. Rev. A **70**, 032324 (2004)

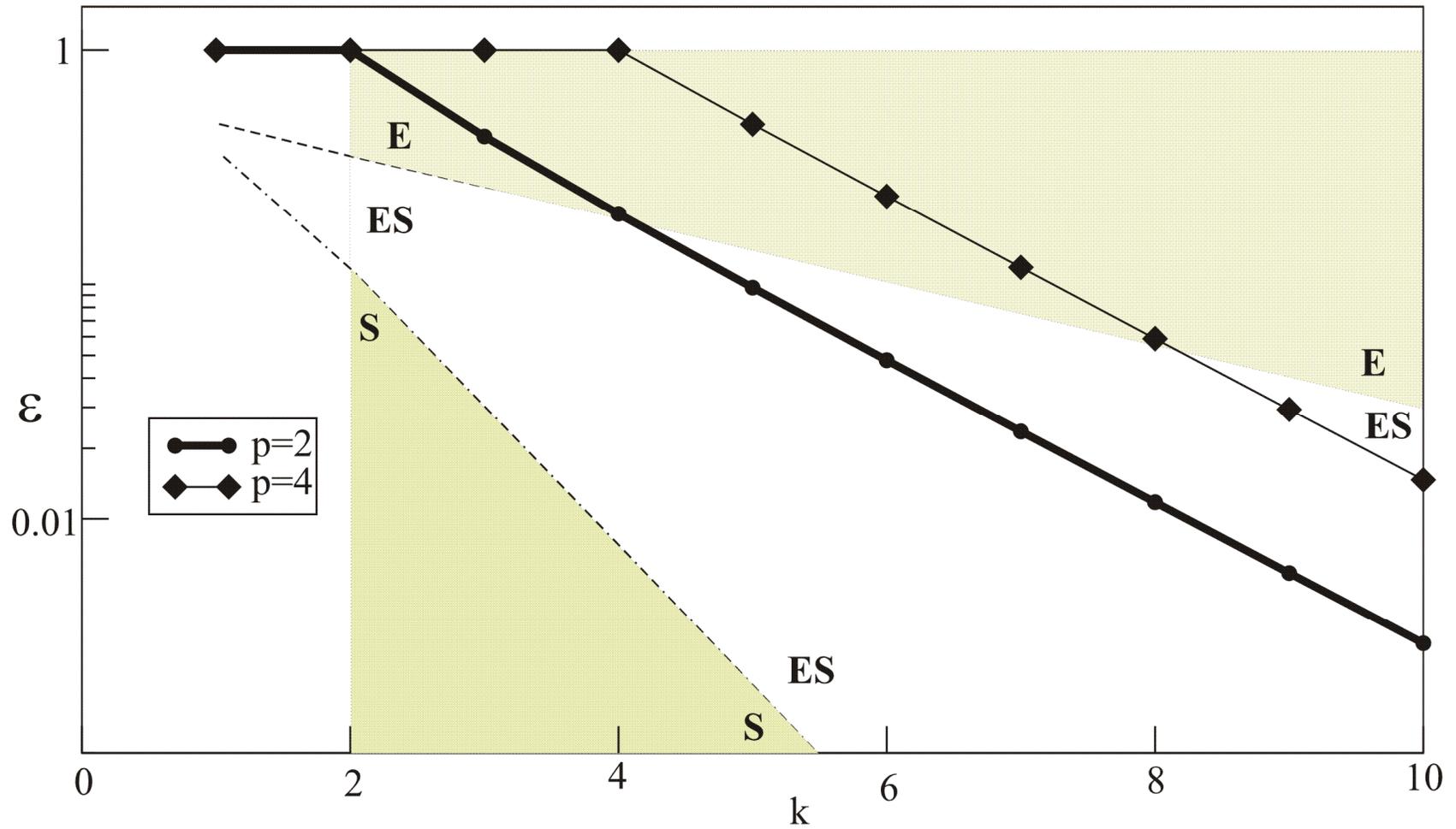
Purity dilution

Purity Sharing 1



M.S. Anwar et al. Phys. Rev. A **73**, 022322 (2006)

Purity Sharing 2



Purity Sharing 3

