

The Schwarzschild solution

Bharat Kumar^{1,2*}

¹*Institute of Physics, Bhubaneswar-751 005, India. and*
²*Homi Bhabha National Institute, Training School Complex,*
Anushakti Nagar, Mumbai 400085, India

Derivation of TOV equation:

Consider the general static, spherically symmetric metric [1]

$$ds^2 = -e^{2\alpha(r)} dt^2 + e^{2\beta(r)} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2 \quad (1)$$

Let's now take this metric and use Einstein's equation to solve the function $\alpha(r)$ and $\beta(r)$. We are looking for nonvacuum solutions, so we turn to the full Einstein equation,

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G T_{\mu\nu} \quad (2)$$

where $R_{\mu\nu}$ and R are the Ricci tensor and Ricci scalar, respectively. $g_{\mu\nu}$ is the metric

$$g_{\mu\nu} = \begin{bmatrix} g_{tt} & g_{tr} & g_{t\theta} & g_{t\phi} \\ g_{rt} & g_{rr} & g_{r\theta} & g_{r\phi} \\ g_{\theta t} & g_{\theta r} & g_{\theta\theta} & g_{\theta\phi} \\ g_{\phi t} & g_{\phi r} & g_{\phi\theta} & g_{\phi\phi} \end{bmatrix} = \begin{bmatrix} -e^{2\alpha(r)} & 0 & 0 & 0 \\ 0 & e^{2\beta(r)} & 0 & 0 \\ 0 & 0 & r^2 & 0 \\ 0 & 0 & 0 & r^2 \sin^2 \theta \end{bmatrix}$$

We begin by evaluating the christoffel (or affine connection or Levi-Civita or Riemann connection) symbols. If we use labels (t, r, θ, ϕ) .

$$\Gamma_{\mu\nu}^\lambda = \frac{1}{2} g^{\lambda\sigma} (\partial_\mu g_{\sigma\nu} + \partial_\nu g_{\mu\sigma} - \partial_\sigma g_{\mu\nu}) \quad (3)$$

$$\begin{aligned} \Gamma_{tr}^t &= \frac{1}{2} g^{t\sigma} (\partial_t g_{\sigma r} + \partial_r g_{t\sigma} - \partial_\sigma g_{tr}) \\ &= \frac{1}{2} g^{t\sigma} (0 + \partial_r g_{t\sigma} - 0) \\ &= \frac{1}{2} g^{tt} (\partial_r g_{tt}) \\ &= \frac{1}{2} (-e^{-2\alpha(r)}) (\partial_r (-e^{2\alpha(r)})) \\ &= \frac{1}{2} (-e^{-2\alpha(r)}) (-2e^{2\alpha(r)} \partial_r \alpha) \\ &= \partial_r \alpha \end{aligned} \quad (4)$$

*Electronic address: bharat@iopb.res.in

$$\begin{aligned}
\Gamma_{tt}^r &= \frac{1}{2}g^{r\sigma}(\partial_t g_{\sigma t} + \partial_t g_{t\sigma} - \partial_\sigma g_{tt}) \\
&= \frac{1}{2}g^{rr}(\partial_t g_{rt} + \partial_t g_{tr} - \partial_r g_{tt}) \\
&= \frac{1}{2}(e^{-2\beta(r)})(0 + 0 - \partial_r(-e^{2\alpha(r)})) \\
&= \frac{1}{2}(e^{-2\beta(r)})(2e^{2\alpha(r)})\partial_r\alpha \\
\Gamma_{tt}^r &= e^{2(\alpha(r)-\beta(r))}\partial_r\alpha
\end{aligned} \tag{5}$$

$$\begin{aligned}
\Gamma_{rr}^r &= \frac{1}{2}g^{r\sigma}(\partial_r g_{\sigma r} + \partial_r g_{r\sigma} - \partial_\sigma g_{rr}) \\
&= \frac{1}{2}g^{rr}(\partial_r g_{rr} + \partial_r g_{rr} - \partial_r g_{rr}) \\
&= \frac{1}{2}g^{rr}(\partial_r g_{rr}) \\
&= \frac{1}{2}(e^{-2\beta(r)})\partial_r(e^{2\beta(r)}) \\
&= \frac{1}{2}(e^{-2\beta(r)})(2)e^{2\beta(r)}\partial_r\beta \\
\Gamma_{rr}^r &= \partial_r\beta
\end{aligned} \tag{6}$$

$$\begin{aligned}
\Gamma_{r\theta}^\theta &= \frac{1}{2}g^{\theta\sigma}(\partial_r g_{\sigma\theta} + \partial_\theta g_{r\sigma} - \partial_\sigma g_{r\theta}) \\
&= \frac{1}{2}g^{\theta\theta}(\partial_r g_{\theta\theta} + \partial_\theta g_{r\theta} - \partial_\theta g_{r\theta}) \\
&= \frac{1}{2}(r^{-2})\partial_r(r^2) = \frac{1}{2}(r^{-2})(2r) \\
\Gamma_{r\theta}^\theta &= \frac{1}{r}
\end{aligned} \tag{7}$$

$$\begin{aligned}
\Gamma_{\theta\theta}^r &= \frac{1}{2}g^{r\sigma}(\partial_\theta g_{\sigma\theta} + \partial_\theta g_{\theta\sigma} - \partial_\sigma g_{\theta\theta}) \\
&= \frac{1}{2}g^{rr}(\partial_\theta g_{r\theta} + \partial_\theta g_{\theta r} - \partial_r g_{\theta\theta}) \\
&= \frac{1}{2}g^{rr}(0 + 0 - \partial_r g_{\theta\theta}) = \frac{1}{2}g^{rr}(-\partial_r g_{\theta\theta}) \\
&= \frac{1}{2}(e^{-2\beta(r)})(-\partial_r r^2) = -\frac{1}{2}(e^{-2\beta(r)})(2r) \\
\Gamma_{\theta\theta}^r &= -re^{-2\beta(r)}
\end{aligned} \tag{8}$$

$$\begin{aligned}
\Gamma_{r\phi}^\phi &= \frac{1}{2}g^{\phi\sigma}(\partial_r g_{\sigma\phi} + \partial_\phi g_{r\sigma} - \partial_\sigma g_{r\phi}) \\
&= \frac{1}{2}g^{\phi\phi}(\partial_r g_{\phi\phi} + \partial_\phi g_{r\phi} - \partial_\phi g_{r\phi}) \\
&= \frac{1}{2}g^{\phi\phi}(\partial_r g_{\phi\phi} + 0 - 0) = \frac{1}{2}g^{\phi\phi}(\partial_r g_{\phi\phi}) \\
&= \frac{1}{2}(r^2 \sin^2 \theta)^{-1} \partial_r (r^2 \sin^2 \theta) \\
&= \frac{1}{2}(r^2 \sin^2 \theta)^{-1} (2r \sin^2 \theta) = \frac{1}{r}
\end{aligned} \tag{9}$$

$$\begin{aligned}
\Gamma_{\phi\phi}^r &= \frac{1}{2}g^{r\sigma}(\partial_\phi g_{\sigma\phi} + \partial_\phi g_{\phi\sigma} - \partial_\sigma g_{\phi\phi}) \\
&= \frac{1}{2}g^{rr}(\partial_\phi g_{r\phi} + \partial_\phi g_{\phi r} - \partial_r g_{\phi\phi}) \\
&= \frac{1}{2}g^{rr}(0 + 0 - \partial_r g_{\phi\phi}) = \frac{1}{2}g^{rr}(\partial_r g_{\phi\phi}) \\
&= \frac{1}{2}(e^{-2\beta(r)})(-\partial_r (r^2 \sin^2 \theta)) \\
&= \frac{1}{2}(e^{-2\beta(r)})(-2r \sin^2 \theta) \\
\Gamma_{\phi\phi}^r &= -r \sin^2 \theta e^{-2\beta(r)}
\end{aligned} \tag{10}$$

$$\begin{aligned}
\Gamma_{\phi\phi}^\theta &= \frac{1}{2}g^{\theta\sigma}(\partial_\phi g_{\sigma\phi} + \partial_\phi g_{\phi\sigma} - \partial_\sigma g_{\phi\phi}) \\
&= \frac{1}{2}g^{\theta\theta}(\partial_\phi g_{\theta\phi} + \partial_\phi g_{\phi\theta} - \partial_\theta g_{\phi\phi}) \\
&= \frac{1}{2}g^{\theta\theta}(0 + 0 - \partial_\theta g_{\phi\phi}) \\
&= \frac{1}{2}(r^2)^{-1}(-\partial_\theta (r^2 \sin^2 \theta)) \\
\Gamma_{\phi\phi}^\theta &= -\sin \theta \cos \theta
\end{aligned} \tag{11}$$

$$\begin{aligned}
\Gamma_{\theta\phi}^\phi &= \frac{1}{2}g^{\phi\sigma}(\partial_\theta g_{\sigma\phi} + \partial_\phi g_{\theta\sigma} - \partial_\sigma g_{\theta\phi}) \\
&= \frac{1}{2}g^{\phi\phi}(\partial_\theta g_{\phi\phi} + \partial_\phi g_{\theta\phi} - \partial_\phi g_{\theta\phi}) \\
&= \frac{1}{2}g^{\phi\phi}(\partial_\theta g_{\phi\phi} + 0 - 0) \\
&= \frac{1}{2}(r^2 \sin^2 \theta)^{-1} \partial_\theta (r^2 \sin^2 \theta) \\
&= \frac{1}{2}(r^2 \sin^2 \theta)^{-1} (2r^2 \sin \theta \cos \theta) \\
\Gamma_{\theta\phi}^\phi &= \frac{\cos \theta}{\sin \theta}
\end{aligned} \tag{12}$$

Summary:

$$\begin{aligned}\Gamma_{tr}^t &= \partial_r \alpha, \Gamma_{tt}^r = e^{2(\alpha-\beta)} \partial_r \alpha, \Gamma_{rr}^r = \partial_r \beta \\ \Gamma_{r\theta}^\theta &= \frac{1}{r}, \Gamma_{\theta\theta}^r = -re^{-2\beta}, \Gamma_{r\phi}^\phi = \frac{1}{r} \\ \Gamma_{\phi\phi}^r &= -re^{-2\beta} \sin^2 \theta, \Gamma_{\phi\phi}^\theta = -\sin \theta \cos \theta, \Gamma_{r\phi}^\phi = \frac{\cos \theta}{\sin \theta}\end{aligned}\quad (13)$$

We know that Riemann tensor

$$R_{\sigma\mu\nu}^\rho = \partial_\mu \Gamma_{\nu\sigma}^\rho - \partial_\nu \Gamma_{\mu\sigma}^\rho + \Gamma_{\mu\lambda}^\rho \Gamma_{\nu\sigma}^\lambda - \Gamma_{\nu\lambda}^\rho \Gamma_{\mu\sigma}^\lambda \quad (14)$$

Now, We calculate Ricci tensor

$$\begin{aligned}R_{\mu\nu} &= R_{\mu\lambda\nu}^\lambda \\ R_{tt} &= R_{t\lambda t}^\lambda = R_{trt}^r + R_{t\theta t}^\theta + R_{t\phi t}^\phi\end{aligned}\quad (15)$$

$$\begin{aligned}R_{trt}^r &= \partial_r \Gamma_{tt}^r - \partial_t \Gamma_{rt}^r + \Gamma_{r\lambda}^r \Gamma_{tt}^\lambda - \Gamma_{t\lambda}^r \Gamma_{rt}^\lambda \\ &= \partial_r \Gamma_{tt}^r - 0 + \Gamma_{rr}^r \Gamma_{tt}^r + \Gamma_{r\theta}^r \Gamma_{tt}^\theta + \Gamma_{r\phi}^r \Gamma_{tt}^\phi - \Gamma_{tt}^r \Gamma_{rt}^t - \Gamma_{tr}^r \Gamma_{rt}^r - \Gamma_{t\theta}^r \Gamma_{rt}^\theta - \Gamma_{t\phi}^r \Gamma_{rt}^\phi \\ &= \partial_r (e^{2(\alpha-\beta)} \partial_r \alpha) + \partial_r \beta e^{2(\alpha-\beta)} \partial_r \alpha + 0 + 0 - e^{2(\alpha-\beta)} \partial_r \alpha \partial_r \alpha - 0 - 0 - 0 \\ &= \partial_r^2 \alpha e^{2(\alpha-\beta)} + \partial_r \alpha \partial_r (e^{2(\alpha-\beta)}) + \partial_r \alpha \partial_r \beta e^{2(\alpha-\beta)} - e^{2(\alpha-\beta)} (\partial_r \alpha)^2 \\ &= \partial_r^2 \alpha e^{2(\alpha-\beta)} + (2e^{2(\alpha-\beta)} \partial_r \alpha - 2e^{2(\alpha-\beta)} \partial_r \beta) \partial_r \alpha + \partial_r \alpha \partial_r \beta e^{2(\alpha-\beta)} - e^{2(\alpha-\beta)} (\partial_r \alpha)^2 \\ &= e^{2(\alpha-\beta)} [\partial_r^2 \alpha + 2(\partial_r \alpha)^2 - 2\partial_r \alpha \partial_r \beta + \partial_r \alpha \partial_r \beta - (\partial_r \alpha)^2] \\ R_{trt}^r &= e^{2(\alpha-\beta)} [\partial_r^2 \alpha + (\partial_r \alpha)^2 - \partial_r \alpha \partial_r \beta]\end{aligned}\quad (16)$$

$$\begin{aligned}R_{t\theta t}^\theta &= \partial_\theta \Gamma_{tt}^\theta - \partial_t \Gamma_{t\theta}^\theta + \Gamma_{\theta\lambda}^\theta \Gamma_{tt}^\lambda - \Gamma_{t\lambda}^\theta \Gamma_{t\theta}^\lambda \\ &= 0 + 0 + \Gamma_{\theta t}^\theta \Gamma_{tt}^t + \Gamma_{\theta r}^\theta \Gamma_{tt}^r + \Gamma_{\theta\theta}^\theta \Gamma_{tt}^\theta + \Gamma_{\theta\phi}^\theta \Gamma_{tt}^\phi - \Gamma_{tt}^\theta \Gamma_{t\theta}^t - \Gamma_{tr}^\theta \Gamma_{t\theta}^r - \Gamma_{t\theta}^\theta \Gamma_{t\theta}^\theta - \Gamma_{t\phi}^\theta \Gamma_{t\theta}^\phi \\ &= 0 + 0 + 0 + \Gamma_{\theta r}^\theta \Gamma_{tt}^r + 0 + 0 - 0 - 0 - 0 - 0 \\ R_{t\theta t}^\theta &= \frac{1}{r} e^{2(\alpha-\beta)} \partial_r \alpha\end{aligned}\quad (17)$$

$$\begin{aligned}R_{t\phi t}^\phi &= \partial_\phi \Gamma_{tt}^\phi - \partial_t \Gamma_{t\phi}^\phi + \Gamma_{\phi\lambda}^\phi \Gamma_{tt}^\lambda - \Gamma_{t\lambda}^\phi \Gamma_{t\phi}^\lambda \\ &= 0 + 0 + \Gamma_{\phi\lambda}^\phi \Gamma_{tt}^\lambda - \Gamma_{t\lambda}^\phi \Gamma_{t\phi}^\lambda \\ &= \Gamma_{\phi t}^\phi \Gamma_{tt}^t + \Gamma_{\phi r}^\phi \Gamma_{tt}^r + \Gamma_{\phi\theta}^\phi \Gamma_{tt}^\theta + \Gamma_{\phi\phi}^\phi \Gamma_{tt}^\phi - \Gamma_{tt}^\phi \Gamma_{t\phi}^t - \Gamma_{tr}^\phi \Gamma_{t\phi}^r - \Gamma_{t\theta}^\phi \Gamma_{t\phi}^\theta - \Gamma_{t\phi}^\phi \Gamma_{t\phi}^\phi \\ &= 0 + \Gamma_{\phi r}^\phi \Gamma_{tt}^r + 0 + 0 - 0 - 0 - 0 - 0 \\ R_{t\phi t}^\phi &= \frac{1}{r} e^{2(\alpha-\beta)} \partial_r \alpha\end{aligned}\quad (18)$$

From eq.(15) can be written as:

$$\begin{aligned}R_{tt} &= R_{t\lambda t}^\lambda = R_{trt}^r + R_{t\theta t}^\theta + R_{t\phi t}^\phi = e^{2(\alpha-\beta)} [\partial_r^2 \alpha + (\partial_r \alpha)^2 - \partial_r \alpha \partial_r \beta] + \frac{2}{r} e^{2(\alpha-\beta)} \partial_r \alpha \\ R_{tt} &= e^{2(\alpha-\beta)} [\partial_r^2 \alpha + (\partial_r \alpha)^2 - \partial_r \alpha \partial_r \beta + \frac{2}{r} \partial_r \alpha]\end{aligned}\quad (19)$$

$$R_{rr} = R_{r\lambda r}^\lambda = R_{rtr}^t + R_{r\theta r}^\theta + R_{r\phi r}^\phi$$

$$\begin{aligned} R_{rtr}^t &= \partial_t \Gamma_{rr}^t - \partial_r \Gamma_{tr}^t + \Gamma_{t\lambda}^t \Gamma_{rr}^\lambda - \Gamma_{r\lambda}^t \Gamma_{tr}^\lambda \\ &= 0 - \partial_r \Gamma_{tr}^t + \Gamma_{tt}^t \Gamma_{rr}^t + \Gamma_{tr}^t \Gamma_{rr}^r + \Gamma_{t\theta}^t \Gamma_{rr}^\theta + \Gamma_{t\phi}^t \Gamma_{rr}^\phi - \Gamma_{rt}^t \Gamma_{tr}^t - \Gamma_{rr}^t \Gamma_{tr}^r - \Gamma_{r\theta}^t \Gamma_{tr}^\theta - \Gamma_{r\phi}^t \Gamma_{tr}^\phi \\ &= -\partial_r(\partial_r \alpha) + 0 + \partial_r \alpha \partial_r \beta + 0 + 0 - \partial_r \alpha \partial_r \alpha - 0 - 0 - 0 \\ R_{rtr}^t &= \partial_r \alpha \partial_r \beta - \partial_r^2 \alpha - (\partial_r \alpha)^2 \quad (20) \end{aligned}$$

$$\begin{aligned} R_{r\theta r}^\theta &= \partial_\theta \Gamma_{rr}^\theta - \partial_r \Gamma_{\theta r}^\theta + \Gamma_{\theta\lambda}^\theta \Gamma_{rr}^\lambda - \Gamma_{r\lambda}^\theta \Gamma_{\theta r}^\lambda \\ &= 0 - \partial_r \left(\frac{1}{r}\right) + \Gamma_{\theta t}^\theta \Gamma_{rr}^t + \Gamma_{\theta r}^\theta \Gamma_{rr}^r + \Gamma_{\theta\theta}^\theta \Gamma_{rr}^\theta + \Gamma_{\theta\phi}^\theta \Gamma_{rr}^\phi - \Gamma_{rt}^\theta \Gamma_{\theta r}^t - \Gamma_{rr}^\theta \Gamma_{\theta r}^r - \Gamma_{r\theta}^\theta \Gamma_{\theta r}^\theta - \Gamma_{r\phi}^\theta \Gamma_{\theta r}^\phi \\ &= \frac{1}{r^2} + 0 + \frac{1}{r} \partial_r \beta + 0 - \frac{1}{r} \frac{1}{r} - 0 - 0 - 0 \\ R_{r\theta r}^\theta &= \frac{1}{r} \partial_r \beta \quad (21) \end{aligned}$$

$$\begin{aligned} R_{r\phi r}^\phi &= \partial_\phi \Gamma_{rr}^\phi - \partial_r \Gamma_{\phi r}^\phi + \Gamma_{\phi\lambda}^\phi \Gamma_{rr}^\lambda - \Gamma_{r\lambda}^\phi \Gamma_{\phi r}^\lambda \\ &= 0 - \partial_r \left(\frac{1}{r}\right) + \Gamma_{\phi t}^\phi \Gamma_{rr}^t + \Gamma_{\phi r}^\phi \Gamma_{rr}^r + \Gamma_{\phi\theta}^\phi \Gamma_{rr}^\theta + \Gamma_{\phi\phi}^\phi \Gamma_{rr}^\phi - \Gamma_{rt}^\phi \Gamma_{\phi r}^t - \Gamma_{rr}^\phi \Gamma_{\phi r}^r - \Gamma_{r\theta}^\phi \Gamma_{\phi r}^\theta - \Gamma_{r\phi}^\phi \Gamma_{\phi r}^\phi \\ &= \frac{1}{r^2} + 0 + \frac{1}{r} \partial_r \beta + 0 - 0 - 0 - 0 - \frac{1}{r} \frac{1}{r} \\ R_{r\phi r}^\phi &= \frac{1}{r} \partial_r \beta \quad (22) \end{aligned}$$

$$\begin{aligned} R_{rr} &= \partial_r \alpha \partial_r \beta - \partial_r^2 \alpha - (\partial_r \alpha)^2 + \frac{1}{r} \partial_r \beta + \frac{1}{r} \partial_r \beta \\ R_{rr} &= \partial_r \alpha \partial_r \beta - \partial_r^2 \alpha - (\partial_r \alpha)^2 + \frac{2}{r} \partial_r \beta \quad (23) \end{aligned}$$

$$R_{\theta\theta} = R_{\theta\lambda\theta}^\lambda = R_{\theta t\theta}^t + R_{\theta r\theta}^r + R_{\theta\phi\theta}^\phi$$

$$\begin{aligned} R_{\theta t\theta}^t &= \partial_t \Gamma_{\theta\theta}^t - \partial_\theta \Gamma_{t\theta}^t + \Gamma_{t\lambda}^t \Gamma_{\theta\theta}^\lambda - \Gamma_{\theta\lambda}^t \Gamma_{t\theta}^\lambda \\ &= 0 - 0 + \Gamma_{tt}^t \Gamma_{\theta\theta}^t + \Gamma_{tr}^t \Gamma_{\theta\theta}^r + \Gamma_{t\theta}^t \Gamma_{\theta\theta}^\theta + \Gamma_{t\phi}^t \Gamma_{\theta\theta}^\phi - \Gamma_{\theta t}^t \Gamma_{t\theta}^t - \Gamma_{\theta r}^t \Gamma_{t\theta}^r - \Gamma_{\theta\theta}^t \Gamma_{t\theta}^\theta - \Gamma_{\theta\phi}^t \Gamma_{t\theta}^\phi \\ &= 0 + \partial_r \alpha (-re^{-2\beta}) + 0 + 0 - 0 - 0 - 0 - 0 \\ R_{\theta t\theta}^t &= -re^{-2\beta} \partial_r \alpha \quad (24) \end{aligned}$$

$$\begin{aligned}
R_{\theta r \theta}^r &= \partial_r \Gamma_{\theta \theta}^r - \partial_\theta \Gamma_{r \theta}^r + \Gamma_{r \lambda}^r \Gamma_{\theta \theta}^\lambda - \Gamma_{\theta \lambda}^r \Gamma_{r \theta}^\lambda \\
&= \partial_r(-re^{-2\beta}) - 0 + \Gamma_{rt}^r \Gamma_{\theta \theta}^t + \Gamma_{rr}^r \Gamma_{\theta \theta}^r + \Gamma_{r\theta}^r \Gamma_{\theta \theta}^\theta + \Gamma_{r\phi}^r \Gamma_{\theta \theta}^\phi - \Gamma_{\theta t}^r \Gamma_{r \theta}^t - \Gamma_{\theta r}^r \Gamma_{r \theta}^r - \Gamma_{\theta \theta}^r \Gamma_{r \theta}^\theta - \Gamma_{\theta \phi}^r \Gamma_{r \theta}^\phi \\
&= 2re^{-2\beta} \partial_r \beta - e^{-2\beta} + 0 + \partial_r \beta (-re^{-2\beta}) + 0 + 0 - 0 - 0 - (-re^{-2\beta}) \frac{1}{r} - 0 \\
&= 2re^{-2\beta} \partial_r \beta - e^{-2\beta} + \partial_r \beta (-re^{-2\beta}) + e^{-2\beta} \\
R_{\theta r \theta}^r &= re^{-2\beta} \partial_r \beta \dots (25)
\end{aligned}$$

$$\begin{aligned}
R_{\theta \phi \theta}^\phi &= \partial_\phi \Gamma_{\theta \theta}^\phi - \partial_\theta \Gamma_{\phi \theta}^\phi + \Gamma_{\phi \lambda}^\phi \Gamma_{\theta \theta}^\lambda - \Gamma_{\theta \lambda}^\phi \Gamma_{\phi \theta}^\lambda \\
&= 0 - \partial_\theta \Gamma_{\phi \theta}^\phi + \Gamma_{\phi t}^\phi \Gamma_{\theta \theta}^t + \Gamma_{\phi r}^\phi \Gamma_{\theta \theta}^r + \Gamma_{\phi \theta}^\phi \Gamma_{\theta \theta}^\theta + \Gamma_{\phi \phi}^\phi \Gamma_{\theta \theta}^\phi - \Gamma_{\theta t}^\phi \Gamma_{\phi \theta}^t - \Gamma_{\theta r}^\phi \Gamma_{\phi \theta}^r - \Gamma_{\theta \theta}^\phi \Gamma_{\phi \theta}^\theta - \Gamma_{\theta \phi}^\phi \Gamma_{\phi \theta}^\phi \\
&= -\partial_\theta (\cot \theta) + 0 + \frac{1}{r} (-re^{-2\beta}) + 0 - 0 - 0 - 0 - 0 - \cot^2 \theta \\
R_{\theta \phi \theta}^\phi &= (1 - e^{-2\beta}) \dots (26)
\end{aligned}$$

$$\begin{aligned}
R_{\theta \theta} &= -re^{-2\beta} \partial_r \alpha + re^{-2\beta} \partial_r \beta + (1 - e^{-2\beta}) \\
R_{\theta \theta} &= e^{-2\beta} [r(\partial_r \beta - \partial_r \alpha) - 1] + 1 \quad (27)
\end{aligned}$$

$$\begin{aligned}
R_{\phi \phi} &= R_{\phi \lambda \phi}^\lambda = R_{\phi t \phi}^t + R_{\phi r \phi}^r + R_{\phi \theta \phi}^\theta \\
&= -re^{-2\beta} \sin^2 \theta \partial_r \alpha + re^{-2\beta} \sin^2 \theta \partial_r \beta + (1 - e^{-2\beta}) \sin^2 \theta \\
&= \{e^{-2\beta} [r(\partial_r \beta - \partial_r \alpha) - 1] + 1\} \sin^2 \theta \\
R_{\phi \phi} &= \sin^2 \theta R_{\theta \theta} \quad (28)
\end{aligned}$$

Note: $\partial_r^2 \alpha = \alpha''$, $\partial_r \alpha = \alpha'$, $\partial_r \beta = \beta'$
Curvature scalar or Ricci scalar is

$$\begin{aligned}
R &= g^{\mu\nu} R_{\mu\nu} = g^{tt} R_{tt} + g^{rr} R_{rr} + g^{\theta\theta} R_{\theta\theta} + g^{\phi\phi} R_{\phi\phi} \\
&= -e^{-2\alpha} e^{2(\alpha-\beta)} [\alpha'' - \alpha'^2 - \alpha' \beta' + \frac{2}{r} \alpha'] + e^{-2\beta} [-\alpha'' - \alpha'^2 + \alpha' \beta' \\
&\quad + \frac{2}{r} \beta'] + \frac{1}{r^2} \{e^{-2\beta} [r(\beta' - \alpha') - 1] + 1\} + \frac{1}{r^2 \sin^2 \theta} \{e^{-2\beta} [r(\beta' - \alpha') - 1] + 1\} \sin^2 \theta \\
&= e^{-2\beta} [-\alpha'' - \alpha'^2 + \alpha' \beta' - \frac{2}{r} \alpha' - \alpha'' - \alpha'^2 + \alpha' \beta' + \frac{2}{r} \beta'] + \frac{2}{r^2} \{e^{-2\beta} [r(\beta' - \alpha') - 1] + 1\} \\
&= 2e^{-2\beta} [-\alpha'' - \alpha'^2 + \alpha' \beta' - \frac{1}{r} (\alpha' - \beta')] + 2e^{-2\beta} [\frac{1}{r} (\beta' - \alpha') - \frac{1}{r^2}] + \frac{2}{r^2} \\
&= -2e^{-2\beta} [-\alpha'' - \alpha'^2 + \alpha' \beta' - \frac{2}{r} (\alpha' - \beta') - \frac{1}{r^2}] + \frac{2}{r^2} \\
&= -2e^{-2\beta} [\alpha'' + \alpha'^2 - \alpha' \beta' + \frac{2}{r} (\alpha' - \beta') + \frac{1}{r^2} (1 - e^{2\beta})] \\
R &= 2e^{-2\beta} [\alpha'' + \alpha'^2 - \alpha' \beta' + \frac{2}{r} (\alpha' - \beta') + \frac{1}{r^2} (1 - e^{2\beta})] \dots (29)
\end{aligned}$$

Now, calculate full Einstien's equation

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R \quad (30)$$

tt-component:

$$\begin{aligned} G_{tt} &= R_{tt} - \frac{1}{2}g_{tt}R \\ &= e^{2(\alpha-\beta)}[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}\alpha'] - \frac{1}{2}(-2e^{2\alpha})(-2e^{-2\beta})[\alpha'' + \alpha' - \alpha'\beta' + \frac{2}{r}(\alpha' - \beta') + \frac{1}{r^2}(1 - e^{2\beta})] \\ &= e^{2(\alpha-\beta)}[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}\alpha' - \alpha'' - \alpha'^2 + \alpha'\beta' - \frac{2}{r}(\alpha' - \beta') - \frac{1}{r^2}(1 - e^{2\beta})] \\ &= e^{2(\alpha-\beta)}[\frac{2}{r}\beta' - \frac{1}{r^2}(1 - e^{2\beta})] \\ &= \frac{1}{r^2}e^{2(\alpha-\beta)}[2\beta'r - (1 - e^{2\beta})] \\ G_{tt} &= \frac{1}{r^2}e^{2(\alpha-\beta)}[2r\beta' - 1 + e^{2\beta}].....(31) \end{aligned}$$

rr-component:

$$\begin{aligned} G_{rr} &= R_{rr} - \frac{1}{2}g_{rr}R \\ &= -\alpha'' - \alpha'^2 + \alpha'\beta' + \frac{2}{r}\beta' - \frac{1}{2}e^{2\beta}(-2e^{-2\beta})[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}(\alpha' - \beta') + \frac{1}{r^2}(1 - e^{2\beta})] \\ &= \frac{2}{r}\alpha' + \frac{1}{r^2}(1 - e^{2\beta}) \\ G_{rr} &= \frac{1}{r^2}[2r\alpha' + 1 - e^{2\beta}]..... (32) \end{aligned}$$

$$\begin{aligned} G_{\theta\theta} &= R_{\theta\theta} - \frac{1}{2}g_{\theta\theta}R \\ &= e^{-2\beta}[r(\beta' - \alpha') - 1] + 1 - \frac{1}{r^2}(-2e^{-2\beta})[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}(\alpha' - \beta') + \frac{1}{r^2}(1 - e^{2\beta})] \\ G_{\theta\theta} &= r^2e^{-2\beta}[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{1}{r}(\alpha' - \beta')]..... (33) \end{aligned}$$

$$\begin{aligned} G_{\phi\phi} &= R_{\phi\phi} - \frac{1}{2}g_{\phi\phi}R \\ &= \sin^2\theta\{e^{-2\beta}[r(\beta' - \alpha') - 1] + 1\} - \frac{1}{2}r^2\sin^2\theta(-2e^{-2\beta})[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}(\alpha' - \beta') + \frac{1}{r^2}(1 - e^{2\beta})] \\ &= e^{-2\beta}\sin^2\theta[r(\beta' - \alpha') - 1 + e^{2\beta}] + r^2\sin^2\theta e^{-2\beta}[\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{2}{r}(\alpha' - \beta') + \frac{1}{r^2}(1 - e^{2\beta})] \\ &= e^{-2\beta}\sin^2\theta[r\beta' - r\alpha' - 1 + e^{2\beta} + r^2\alpha'' + r^2\alpha'^2 - r^2\alpha'\beta' + 2r\alpha' - 2r\beta' + 1 - e^{2\beta}] \\ &= e^{-2\beta}\sin^2\theta[r\alpha' - r\beta' + r^2\alpha'' + r^2\alpha'^2 - r^2\alpha'\beta'] \\ G_{\phi\phi} &= \sin^2\theta G_{\theta\theta}....(34) \end{aligned}$$

Summary:

$$\begin{aligned}
G_{tt} &= \frac{1}{r^2} e^{2(\alpha-\beta)} (2r\beta' - 1 + e^{2\beta}) \\
G_{rr} &= \frac{1}{r^2} (2r\alpha' + 1 - e^{2\beta}) \\
G_{\theta\theta} &= r^2 e^{-2\beta} [\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{1}{r}(\alpha' - \beta')] \\
G_{\phi\phi} &= \sin^2\theta G_{\theta\theta}
\end{aligned} \tag{35}$$

The energy-momentum tensor of the star itself as a perfect fluid is given by

$$T_{\mu\nu} = (\epsilon + p)u_\mu u_\nu + pg_{\mu\nu} \tag{36}$$

Here, the energy density (ϵ) and pressure density (p) is a function of r alone. Since, we pursue static solution, we can take the four velocity to be pointing in timelike direction. Normalized to $u_\mu u_\nu = -1$, it becomes

$$u_\mu = (e^\alpha, 0, 0, 0) \tag{37}$$

So, the component of energy momentum tensor will be

$$T = \begin{bmatrix} e^{2\alpha(r)}\epsilon & 0 & 0 & 0 \\ 0 & e^{2\beta(r)}p & 0 & 0 \\ 0 & 0 & r^2p & 0 \\ 0 & 0 & 0 & r^2p\sin^2\theta \end{bmatrix}$$

Einstein tensor:

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi GT_{\mu\nu} \tag{38}$$

Using all previous results we can find that the components of the Einstein tensor are:
The tt -component:

$$\begin{aligned}
\frac{1}{r^2} e^{2(\alpha-\beta)} (2r\beta' - 1 + e^{2\beta}) &= 8\pi G e^{2\alpha} \epsilon \\
\frac{1}{r^2} e^{-2\beta} (2r\beta' - 1 + e^{2\beta}) &= 8\pi G \epsilon
\end{aligned} \tag{39}$$

The rr -component:

$$\begin{aligned}
\frac{1}{r^2} (2r\alpha' + 1 - e^{2\beta}) &= 8\pi G e^{2\beta} p \\
\frac{1}{r^2} e^{-2\beta} (2r\alpha' + 1 - e^{2\beta}) &= 8\pi G p
\end{aligned} \tag{40}$$

The $\theta\theta$ -component:

$$\begin{aligned}
r^2 e^{-2\beta} [\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{1}{r}(\alpha' - \beta')] &= 8\pi G r^2 p \\
e^{-2\beta} [\alpha'' + \alpha'^2 - \alpha'\beta' + \frac{1}{r}(\alpha' - \beta')] &= 8\pi G p
\end{aligned} \tag{41}$$

The $\phi\phi$ -component is proportional to the $\theta\theta$ -equation, so there is no need to consider separately. So, the tt-component of the Einstein equation gives:

$$\begin{aligned}
\frac{1}{r^2}e^{-2\beta}(2r\beta' - 1 + e^{2\beta}) &= 8\pi G\epsilon \\
e^{-2\beta}(2r\beta' - 1) + 1 &= 8\pi G\epsilon r^2 \\
-\frac{d}{dr}\{r(e^{-2\beta} - 1)\} &= 8\pi G\epsilon r^2 \\
\frac{d}{dr}\{r(e^{-2\beta} - 1)\} &= -k\epsilon r^2 \\
d\{r(e^{-2\beta} - 1)\} &= -k\epsilon r^2 dr
\end{aligned} \tag{42}$$

Where, $k = -8\pi G$ and previous equation can be integrated

$$e^{-2\beta} = 1 - \frac{k}{r} \int_0^R \epsilon(r) r^2 dr \tag{43}$$

Let us define

$$\begin{aligned}
m(r) &= 4\pi \int_0^R \epsilon(r) r^2 dr \\
e^{-2\beta} &= 1 - \frac{8\pi G}{4\pi r} m(r) \\
e^{-2\beta} &= \left(1 - \frac{2Gm(r)}{r}\right) \\
e^{2\beta} &= \left(1 - \frac{2Gm(r)}{r}\right)^{-1}
\end{aligned} \tag{44}$$

So that, the metric will be

$$ds^2 = -e^{2\alpha(r)} dt^2 + \left(1 - \frac{2Gm(r)}{r}\right)^{-1} dr^2 + r^2 d\theta^2 + r^2 \sin^2\theta d\phi^2 \tag{45}$$

M_G = graviational mass

$$\begin{aligned}
M_G &= 4\pi \int_0^R \epsilon(r) r^2 e^{2\beta} dr \\
M_G &= 4\pi \int_0^R \frac{\epsilon(r) r^2}{\left(1 - \frac{2Gm(r)}{r}\right)} dr
\end{aligned} \tag{46}$$

The binding energy due to the internal gravitational attraction of the fluid elements in the star, which is given by

$$E_B = M_G - M > 0 \tag{47}$$

The binding energy is the amount of energy that would be required to disperse the matter in

the star to be infinity. The rr-component can be written as

$$\begin{aligned}
\frac{1}{r^2}e^{-2\beta}(2r\frac{d\alpha}{dr} + 1 - e^{-2\beta}) &= 8\pi Gp \\
2r\frac{d\alpha}{dr} + 1 - e^{2\beta} &= 8\pi Gpr^2e^{2\beta} \\
2r\frac{d\alpha}{dr} &= 8\pi Gpr^2e^{2\beta} + e^{2\beta} - 1 \\
&= \frac{(8\pi Gpr^2 + 1)}{(1 - \frac{2Gm(r)}{r})} - 1 \\
&= \frac{(8\pi Gpr^3 + r)}{(r - 2Gm(r))} - 1 \\
&= \frac{8\pi Gpr^3 + r - r + 2Gm(r)}{(r - 2Gm(r))} \\
&= \frac{8\pi Gpr^3 + 2Gm(r)}{(r - 2Gm(r))} \\
\frac{d\alpha}{dr} &= \frac{(4\pi Gpr^3 + Gm(r))}{(r - 2Gm(r))} \tag{48}
\end{aligned}$$

Now, we have to calculate $\frac{dp}{dr} = ?$. Again, solve the tt-component equation

$$\begin{aligned}
(2r\beta' - 1 + e^{2\beta}) &= 8\pi G\epsilon r^2e^{2\beta} \\
2r\beta' &= (8\pi G\epsilon r^2e^{2\beta} - 1)e^{2\beta} + 1 \tag{49}
\end{aligned}$$

$$\begin{aligned}
(2r\alpha' + 1 - e^{2\beta}) &= 8\pi Gpr^2e^{2\beta} \\
2r\alpha' &= (1 + 8\pi Gpr^2e^{2\beta})e^{2\beta} - 1 \tag{50}
\end{aligned}$$

Take the derivative of the previous equation then

$$\begin{aligned}
2\alpha' + 2r\alpha'' &= (8\pi Gp'r^2 + 16\pi Gpr)e^{2\beta} + 2\beta'(1 + 8\pi Gpr^2)e^{2\beta} \\
2r\alpha' + 2r^2\alpha'' &= [2r\beta'(1 + 8\pi Gr^2p) + (16\pi Gr^2p + 8\pi Gr^3p')]e^{2\beta} \\
2r^2\alpha'' &= [2r\beta'(1 + 8\pi Gr^2p) + (16\pi Gr^2p + 8\pi Gr^3p')]e^{2\beta} - 2r\alpha' \tag{51}
\end{aligned}$$

Put the values of $2r\alpha'$ and $2r\beta'$ from previous equations then

$$2r^2\alpha'' = 1 + (16\pi Gr^2p + 8\pi Gr^3p')e^{2\beta} - (1 + 8\pi r^2p)(1 - 8\pi Gr^2\epsilon)e^{4\beta} \tag{52}$$

Square eq.(50) to obtain the result

$$2r^2\alpha'^2 = \frac{1}{2}(1 + 8\pi Gr^2p)^2e^{4\beta} - (1 + 8\pi Gr^2p)e^{2\beta} - \frac{1}{2} \tag{53}$$

Now, we have the expressions for α' , α'' , α'^2 , and β' in terms of p , p' , ϵ , and $e^{2\beta}$. Hence, equation (52) can be written as:

$$\begin{aligned}
(\epsilon + p)\frac{d\alpha}{dr} &= -\frac{dp}{dr} \\
\frac{dp}{dr} &= -\frac{(\epsilon + p)(Gm(r) + 4\pi Gr^3p)}{r(r - 2Gm(r))} \tag{54}
\end{aligned}$$

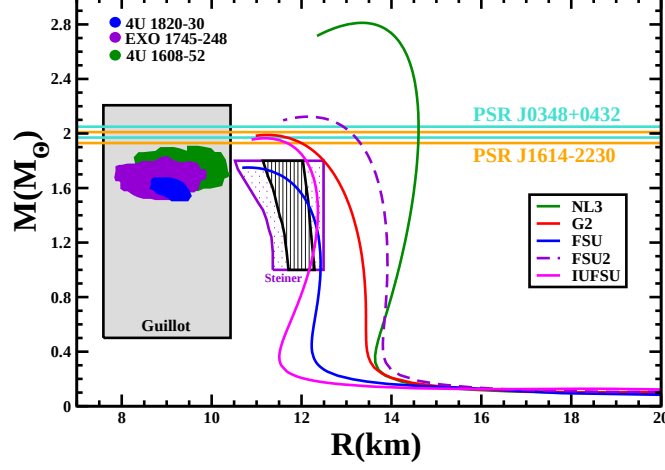


FIG. 1: The mass-radius profile for the force parameters.

Summary: For $G=c=1$, the TOV equations are written as [2]

$$\frac{dp}{dr} = -\frac{(\epsilon + p)(m(r) + 4\pi r^3 p)}{r(r - 2m(r))} \quad (55)$$

$$\frac{dm}{dr} = 4\pi r^2 \epsilon(r) \quad (56)$$

For a given EOS [3], the Tolmann-Oppenheimer-Volkov (TOV) equations must be integrated from the boundary conditions $P(0) = P_c$, and $M(0) = 0$, where P_c and $M(0)$ are the pressure and mass of the star at $r = 0$ and the value of $r (= R)$, where the pressure vanish defines the surface of the star. Thus, at each central density we can uniquely determine a mass M and a radius R of the static neutron and hyperon stars using the four chosen EOSs. The calculated results shown in Fig.1.

References

- [1] Sean M. Carroll, Spacetime and Geometry, An introduction to general relativity.
- [2] Compact stars by Glendenning, 1996.
- [3] Bharat Kumar, S. K. Biswal, and S. K. Patra, Phys. Rev. C **95** (2017) 015801.