

Gauge-gravity duality and aspects of strongly coupled systems

Arnab Kundu
The University of Texas at Austin

Institute of Physics, Bhubaneswar
February 12, 2013

Outline

Holography: gauge-gravity duality

the concept

Where string theory enters

AdS/CFT correspondence: specific realizations and a strong-weak duality

AdS/CFT and strongly coupled physics at RHIC

what we are learning

Applications to other strongly coupled systems

top-down vs bottom-up approaches

Taking stock & Conclusions

where we stand, where to go

Black holes: the “harmonic oscillator” a la mode

Conceived by Laplace a long time back ~ 18th century



Picture taken from Wikipedia image

Solutions of Einstein's equations of motion
They exist!

Characterized by an event-horizon: nothing inside
it can ever come out

Perfect tool to play with various theoretical
concepts

Holography: gauge-gravity duality

Apply quantum mechanics to black holes: the black hole ain't so black!

The event-horizon gives Hawking radiation
black hole has a temperature and an entropy
the entropy goes as the area of the event-horizon

$$S = \frac{A}{4G} \left(\frac{k_B c^3}{\hbar} \right)$$

(Bekenstein, Hawking '70s)

Quantum gravity in (d+1)-dim spacetime = theory living on the d-dim boundary

Quantum field theory in d-spacetime dimensions is described by quantum gravity in (d+1)-
dimensions & vice versa

('t Hooft, Susskind '90s)

Where strings enter: AdS/CFT

Large N gauge theories are secretly string theory

('t Hooft)

Where strings enter: AdS/CFT

Large N gauge theories are secretly string theory

(‘t Hooft)

Concrete examples of the holographic principle can be realized

generally known as the AdS/CFT correspondence

classical gravity in $(d+1)$ -dim anti de-Sitter = strongly coupled conformal field theory in d -dim

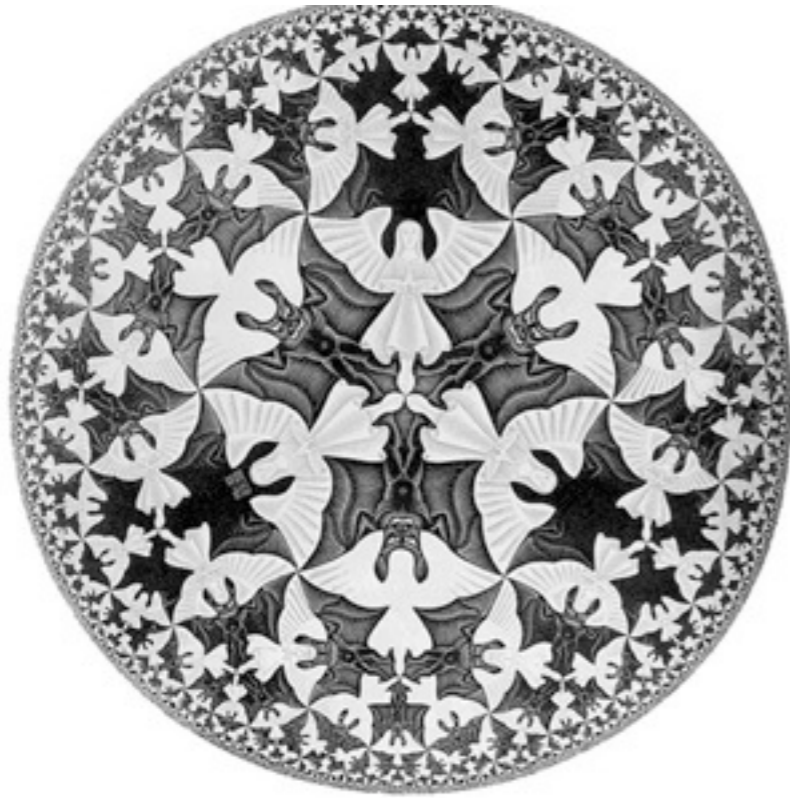
a family of such examples, both conformal and non-conformal

and the list keeps growing ...

(Maldacena '98)

A strong-weak duality, controllable at large N

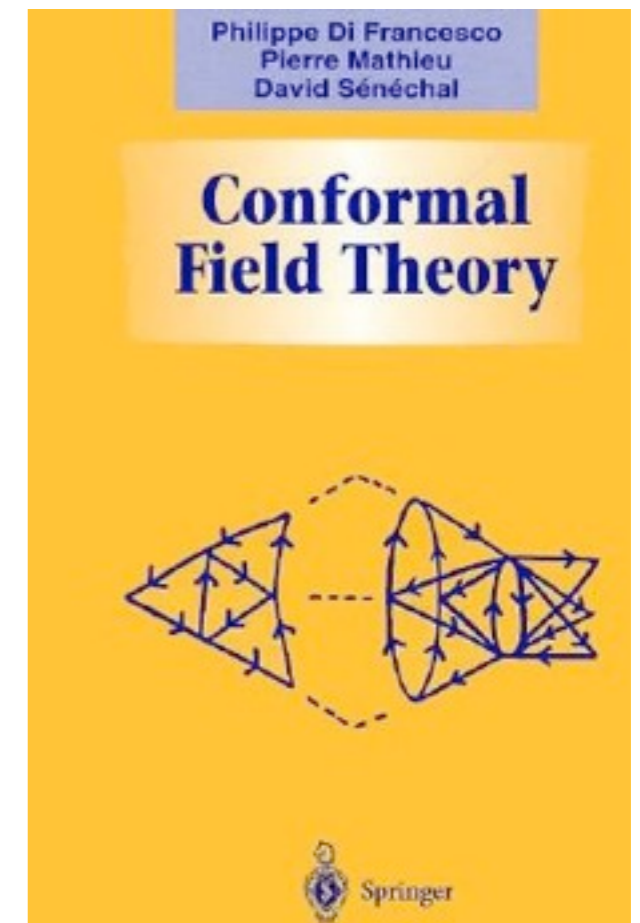
Introducing the characters



Courtesy: M. C. Escher

AdS = solution of Einstein gravity with a -ve cosmological constant

CFT = describes scale-invariant systems

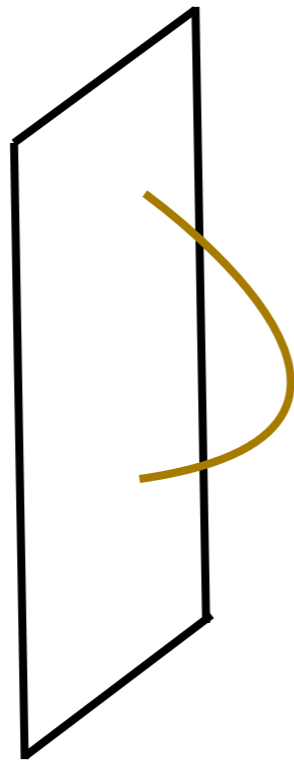


Ingredients from string theory

D_p-branes: (p+1)-dim extended object where a string ends
(e.g., *D3-branes*)

Ingredients from string theory

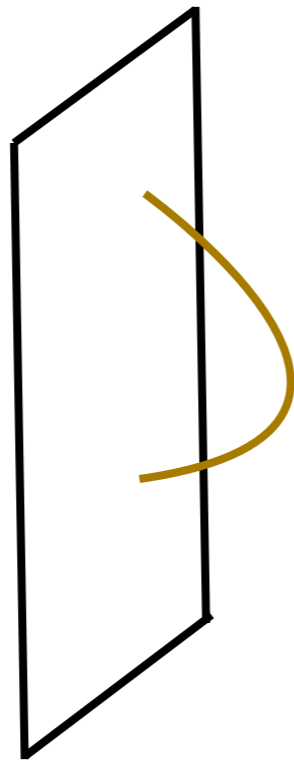
D_p-branes: (p+1)-dim extended object where a string ends
(e.g., D3-branes)



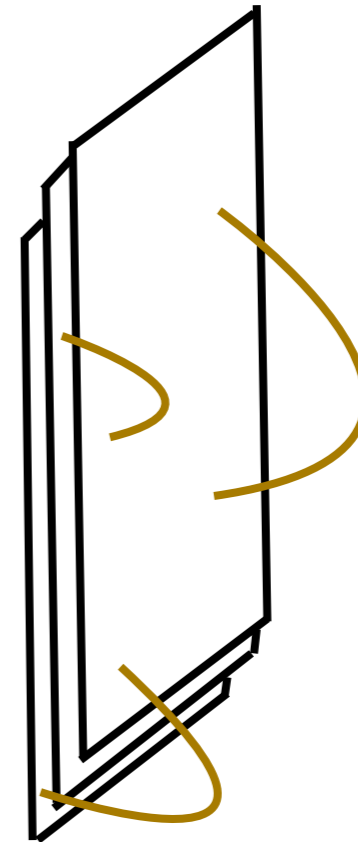
Physics described by U(1)
susy gauge theory

Ingredients from string theory

D_p-branes: (p+1)-dim extended object where a string ends
(e.g., D3-branes)



Physics described by U(1)
susy gauge theory



U(N) for N coincident branes

low energy physics described by a (p+1)-dim gauge theory

(e.g., $\mathcal{N}=4$ SYM)

Ingredients from string theory

(e.g., *D3-branes*)

Branes also have gravitational footprint

Decoupling of the gauge theory from the rest of the “stringy” physics gives a “near-horizon” geometry

$$AdS_5 \times X^5$$

The anti de-Sitter part

some compact manifold

N large

$$\lambda = g_{\text{YM}}^2 N \text{ large}$$

controllable geometry

A summary: an example

Classical (super)-gravity in

$$AdS_5 \times S^5$$

(only closed string modes)

Strongly coupled $\mathcal{N} = 4$

super Yang-Mills (SYM)

(only adjoint d.o.f.)

Isometry group:

$$SO(4, 2) \times SO(6)$$

AdS-part

sphere-part

Global symmetry group:

$$SO(4, 2) \times SO(6)$$

conformal group

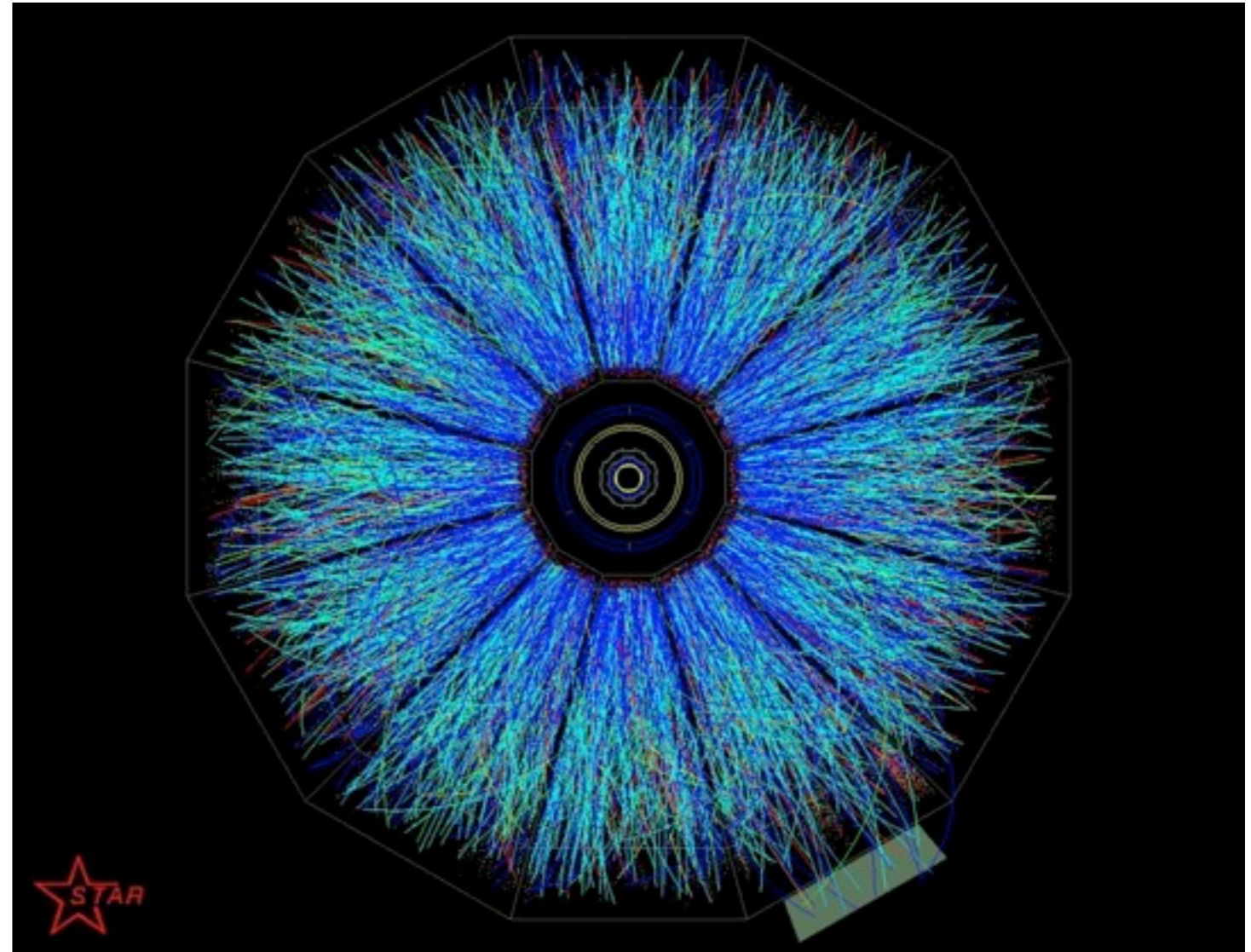
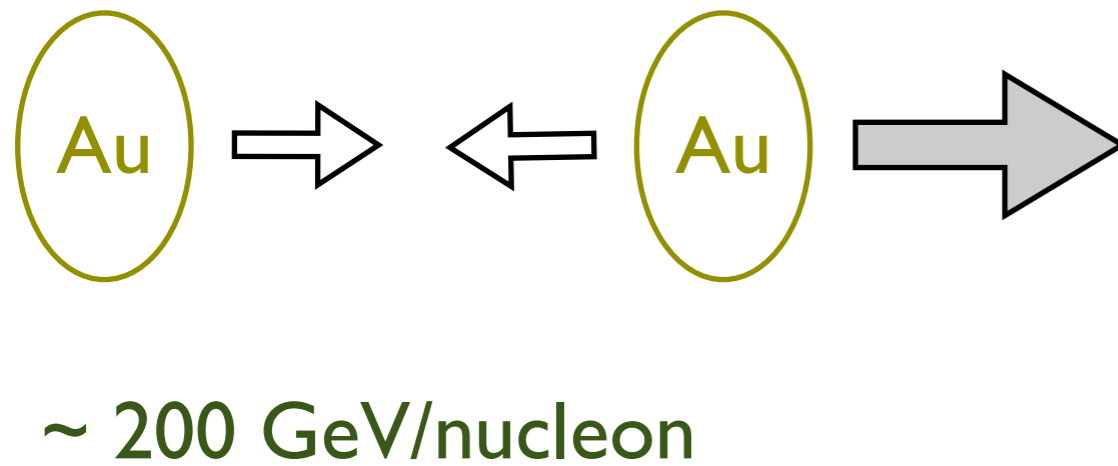
R-symmetry group

AdS-black hole geometry

finite T physics

classical gravity calculations teach us about strongly coupled gauge theory

QGP at RHIC



Courtesy: Wikipedia image

The physics is described by strongly coupled Quantum chromodynamics

This is a hard deal!!

AdS/CFT: can it be useful?

QCD & large N gauge theory: So many differences!!

Perhaps strong coupling and finite T governs the physics ...

QCD

strongly coupled plasma of
gluons and **fundamental**
matter; deconfined, screening,
finite correlation length, ...

SYM

strongly coupled plasma of
gluons and **adjoint** matter;
deconfined, screening, finite
correlation length, ...

⇒ at RHIC energy

May learn qualitatively useful lessons

Towards a smoking gun

RHIC produces a nearly ideal fluid, with a very low viscosity/entropy ratio

$$\frac{\eta}{s} \approx \frac{1 \text{ to } 3}{4\pi}$$

$$\hbar = 1, \quad k_B = 1$$



Courtesy: Wikipedia image

There is no theoretical computation to produce a similar result

The smoking gun

AdS/CFT translates this into a scattering problem
in gravity

Can be performed for a large class of 10-dim
backgrounds:

$$AdS_5 - BH \times X^5$$

some compact manifold

Dual to large N gauge theories with various
amount of susy

(Kovtun, Son, Starinets '05)

Universal result: $\frac{\eta}{s} = \frac{1}{4\pi}$ with $\hbar = 1$
 $k_B = 1$

Taking stock

Some universality indeed exists

The physics is governed by a 5-dim AdS-black hole

Taking stock

Some universality indeed exists

The physics is governed by a 5-dim AdS-black hole

Do the details matter at all: what are the extra dimensions doing?
is it always enough to consider some low dimensional effective gravity theory in AdS?

Taking stock

Some universality indeed exists

The physics is governed by a 5-dim AdS-black hole

Do the details matter at all: what are the extra dimensions doing?

is it always enough to consider some low dimensional effective gravity theory in AdS?

If details do not always matter, can we be more adventurous?

try to capture other strongly coupled systems, without worrying about microscopics

symmetry is the guide

Details matter, at least sometimes

Stringy embedding ensures the duality in a precise sense

Details matter, at least sometimes

Stringy embedding ensures the duality in a precise sense

There is physics where the 10-dimensional details are crucial

e.g. the “QCD” phase diagram

Physics in the flavour sector: “quarks” in AdS/CFT

introduce branes of various dimensions as “test particles” in the 10-dim geometry

these “test particles” are aligned in the 10-dim background in a certain way

Details matter, at least sometimes

Stringy embedding ensures the duality in a precise sense

There is physics where the 10-dimensional details are crucial

e.g. the “QCD” phase diagram

Physics in the flavour sector: “quarks” in AdS/CFT

introduce branes of various dimensions as “test particles” in the 10-dim geometry

these “test particles” are aligned in the 10-dim background in a certain way

stability

what physics we want to engineer

The rough idea: an example

Background geometry is made of N_c D3-branes

Add N_f D7-branes with $N_f \ll N_c$

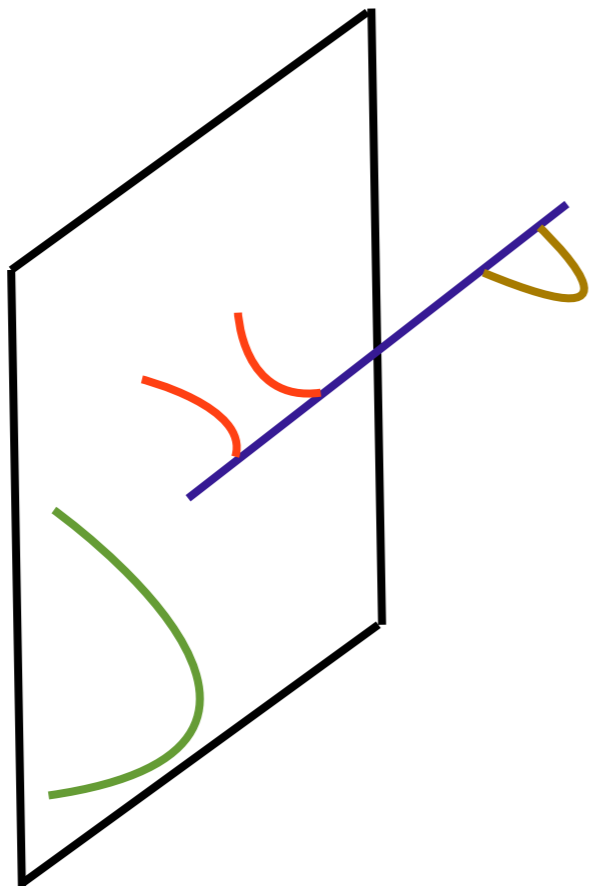
(Karch, Katz '02)

The rough idea: an example

Background geometry is made of N_c D3-branes

Add N_f D7-branes with $N_f \ll N_c$

(Karch, Katz '02)



3-3 strings: adjoint sector

3-7 strings: fundamental matter

7-7 strings: global symmetry

$$U(N_f)$$

D7-branes are simple probes of the geometry

their dynamics determine the physics in the flavour sector

Is there a smoking gun?

A remarkably rich & varied phenomenology is obtained in the flavour sector phase structure

many features are model-dependent

complete QCD phase diagram not well-understood; the results serve as a catalogue, at least

chemical potential is particularly interesting; lattice methods inadequate

Is there a smoking gun?

A remarkably rich & varied phenomenology is obtained in the flavour sector phase structure

many features are model-dependent

complete QCD phase diagram not well-understood; the results serve as a catalogue, at least

chemical potential is particularly interesting; lattice methods inadequate

An intriguing example:

An elegant way to realize spontaneous breaking of chiral symmetry:

$$U(N_f)_L \times U(N_f)_R \rightarrow U(N_f)_{\text{diag}}$$

(Sakai-Sugimoto '04)

Issues that we can address

Limitations apply!

Various phases of flavour matter, in such large N gauge theories + flavours

Phase diagram with various parameters: temperature, chemical potential, electromagnetic fields etc.

Various phase transitions and the order of the transition

Dependence of the phase structure on the number of flavours

Lattice studies are inadequate

Interlude I

When the details don't matter

Physics far from equilibrium

thermalization process of QGP at RHIC or LHC, no good theoretical handle on the physics

beyond linear response theory

When the details don't matter

Physics far from equilibrium

thermalization process of QGP at RHIC or LHC, no good theoretical handle on the physics

beyond linear response theory

The spherical cow approximation

Analogous problem for a strongly coupled
large N gauge theory instead

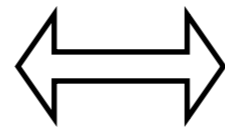
a classical gravity computation



$\mathcal{N} = 4$ SYM

When the black hole is forming

Thermalization process in a
large N gauge theory



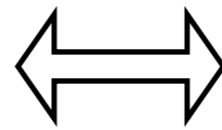
The formation of a black
hole in AdS-space

*Need numerical GR in AdS-space
hard problem*

(Chesler, Lehner, Pretorius etc.)

When the black hole is forming

Thermalization process in a large N gauge theory



The formation of a black hole in AdS-space

*Need numerical GR in AdS-space
hard problem*

(Chesler, Lehner, Pretorius etc.)

A simpler approach: model the black hole formation process *ab initio*

clue: an interpolation between a purely AdS to an AdS-black hole background

strategy: come up with a simple background that does the job

advantage: such a background can be probed for various kinds of physics, *easily*

(Balasubramanian et al '10, etc.)

The AdS-Vaidya background

Gravitational collapse of a shell of null dust

The background is characterized by a mass function:

$$m(t) \rightarrow 0, \quad t \rightarrow -\infty \quad \text{pure AdS}$$

$$m(t) \rightarrow M, \quad t \rightarrow +\infty \quad \text{AdS-black hole}$$

$m(t)$ is a simple interpolating function

We can extract the thermalization time

Study non-local operators

2-pt function, Wilson loop, entanglement entropy = minimal area surface computation in gravity

Is it meaningful?

What one obtains: $\tau_{\text{therm}} \approx 0.3 fm/c$

What one observes: $\tau_{\text{obs}} < 1 fm/c$

(Balasubramanian et al '10, etc.)

Is it meaningful?

What one obtains: $\tau_{\text{therm}} \approx 0.3 fm/c$

What one observes: $\tau_{\text{obs}} < 1 fm/c$

(Balasubramanian et al '10, etc.)

Introducing a chemical potential:

$T\ell \ll 1$ thermalization time decreases with increasing μ/T

$T\ell \gg 1$ thermalization time increases with increasing μ/T

T : temperature

μ : chemical potential

ℓ : length of the operator

(Caceres & Kundu '12)

Interlude II

Top-down vs bottom-up

Top-down

Full 10-dimensional
(super)-gravity embedding

The duality dictionary is
precise

Harder: limited number of
examples exist

Bottom-up

Effective low-dim gravity
model

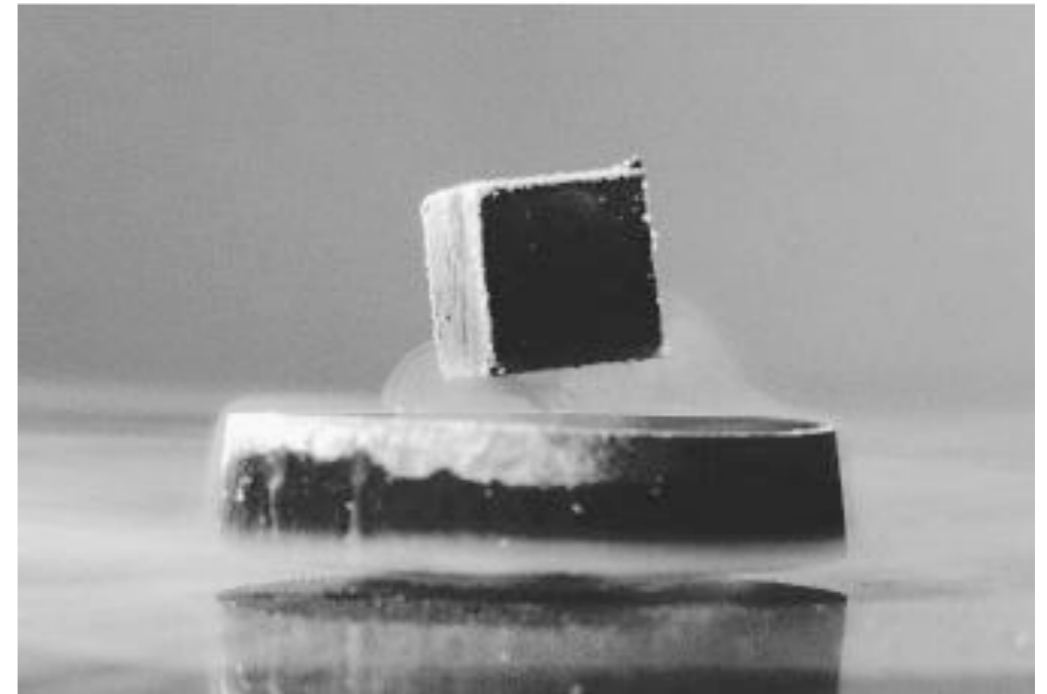
The duality is postulated *ab
initio*

Simpler: more diverse
phenomenology

An "application": holographic superconductors

An interesting state of matter with zero electrical resistivity

BCS theory does explain a class of these

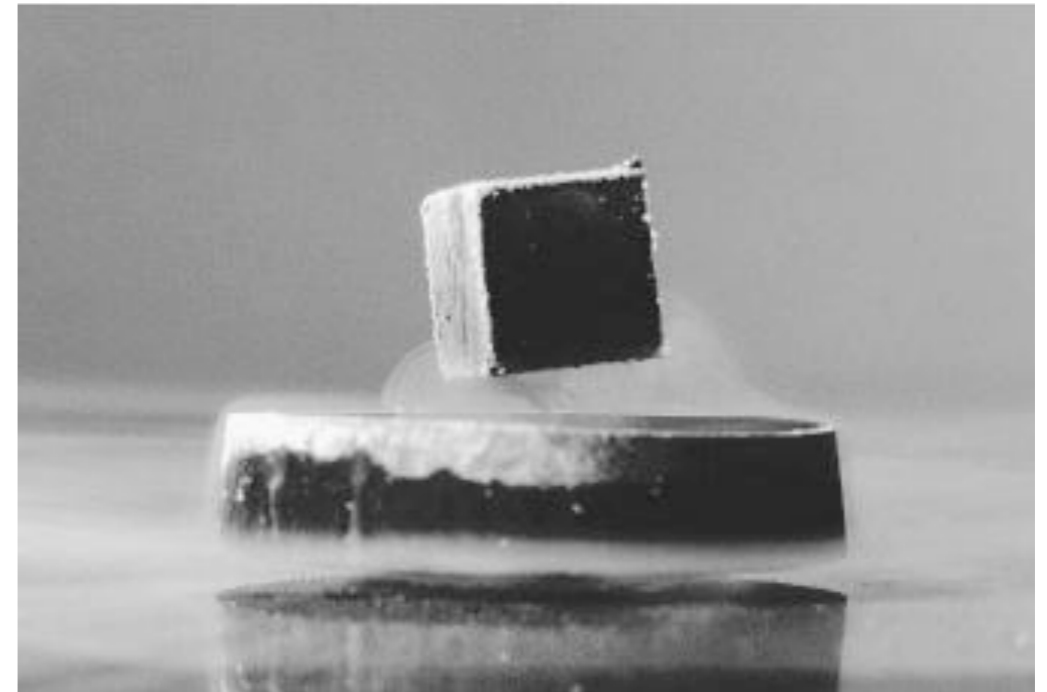


Courtesy: Wiki-image

An "application": holographic superconductors

An interesting state of matter with zero electrical resistivity

BCS theory does explain a class of these



Courtesy: Wiki-image

We will appeal to:

Weinberg: "Superconductivity for particular theorists"

For our purposes, superconductivity = spontaneous breaking of global $U(1)$

Holographic superconductors: ingredients

Abelian Higgs model coupled to Einstein gravity with -ve cosmological constant

(Bottom-up model)

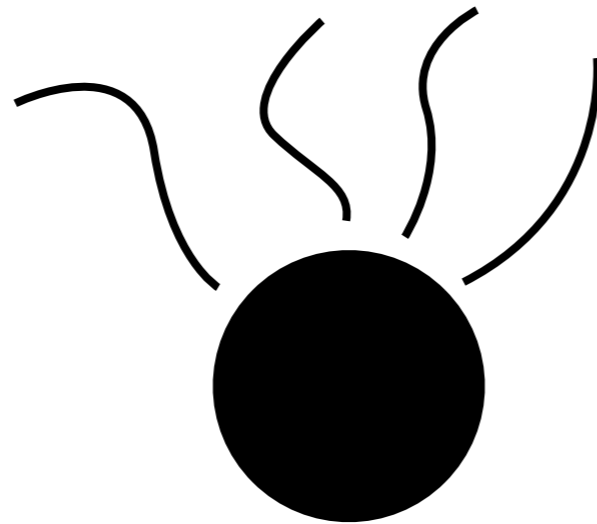
(Gubser; Hartnoll, Herzog, Horowitz '08)

$$\mathcal{L} = \frac{1}{2}R - \frac{1}{4}F^2 - |\partial\phi - iqA\phi|^2 - V(|\phi|)$$

tunable

choose at will

Ground state:



black hole with scalar hair!

hair = condensate, hence symmetry breaking

Holographic SC: stringy embedding

The effective model can be embedded in 10 and 11-dim sugra

(Top-down model, freedom to tune is constrained)

(Gauntlett, Sonner, Wiseman '09)

(Gubser, Herzog, Pufu, Tesileanu '09)

(Most of the) above embeddings are unstable, the fate of the rest is unclear!

this instability can only be seen in the full stringy picture

Holographic SC: stringy embedding

The effective model can be embedded in 10 and 11-dim sugra

(Top-down model, freedom to tune is constrained)

(Gauntlett, Sonner, Wiseman '09)

(Gubser, Herzog, Pufu, Tesileanu '09)

(Most of the) above embeddings are unstable, the fate of the rest is unclear!

this instability can only be seen in the full stringy picture

There is an embedding without any known pathology

it's subtle!

(Bobev, Kundu, Pilch, Warner '11)

Behaves like a superconductor, but has an explicit breaking of the $U(1)$

Universal: hairy black holes are always favoured, a superconducting phase transition with T

can also be a 1st order phase transition

Taking stock: concluding remarks

Stringy embeddings are important

Bottom-up modeling is interesting, but some caution is necessary

Taking stock: concluding remarks

Stringy embeddings are important

Bottom-up modeling is interesting, but some caution is necessary

What are we learning? Why do this?

at least a catalogue of possibilities: phases of matter described by such theories

various phases have been realized: non-Fermi liquids, non-relativistic scale-invariant systems, etc, etc..

Taking stock: concluding remarks

Stringy embeddings are important

Bottom-up modeling is interesting, but some caution is necessary

What are we learning? Why do this?

at least a catalogue of possibilities: phases of matter described by such theories

various phases have been realized: non-Fermi liquids, non-relativistic scale-invariant systems, etc, etc..

if not about condensed-matter physics, we learn about gravity

the existence and constructions of hairy black holes, numerical GR in AdS etc.

Exciting times are ahead!

Thank You

